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Microscopic Approaches for Macroscopic Phenomena of Neutron Stars

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Brief personal history

1988: Born in Tokyo, grew up in Satte City in Saitama



Cherry blossoms
in Satte City

2003-2006: Inagakuen Public High School

2006-2010: Tokyo University of Science (BSc)



2010-2015: University of Tsukuba (MSc-PhD)

Apr. 2015-Aug. 2017: Warsaw University of Technology, Poland (Postdoc)

Sep. 2017-Dec. 2017: University of Washington, USA (Postdoc)

Jan. 2018-Mar. 2021: Niigata University (Assistant Professor, tenure track)

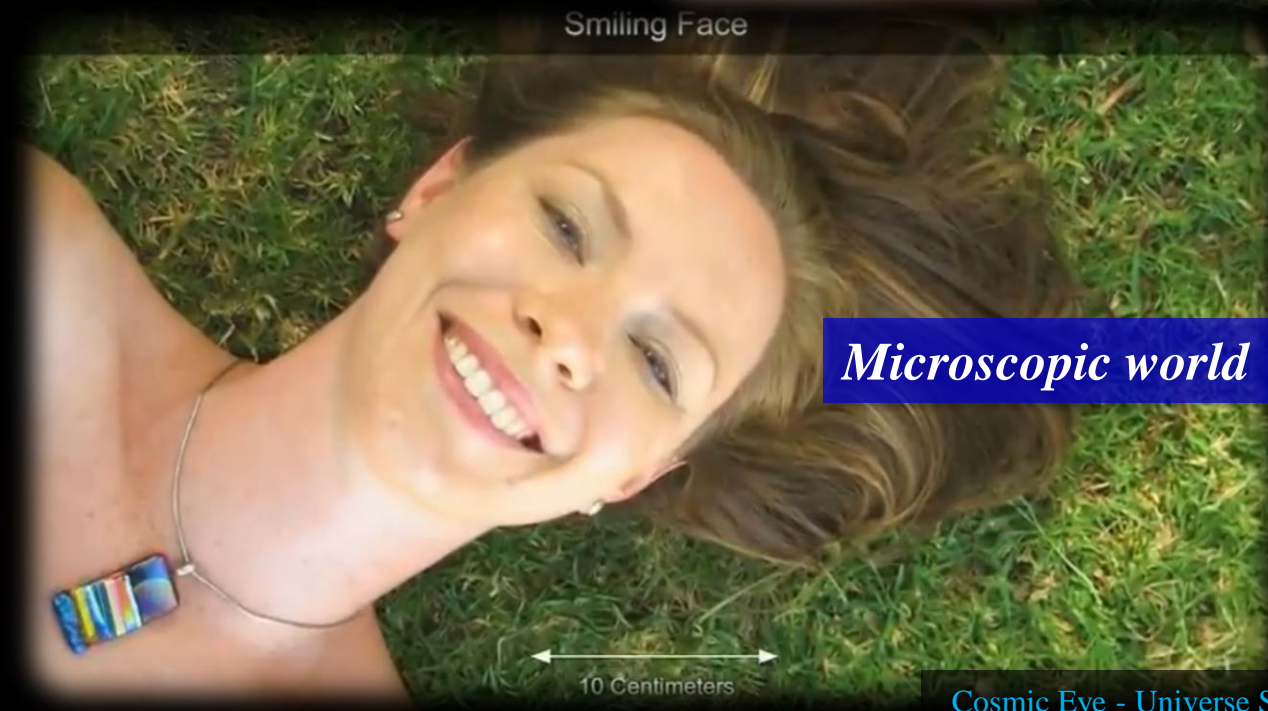


Apr. 2021-Present: Tokyo Institute of Technology (Associate Professor)

Main research field: Nuclear Theory



Hierarchy of Scales in the Universe



Hierarchy of Scales in the Universe

Neutron stars, NS merger, nucleosynthesis, GW, ...



Macroscopic

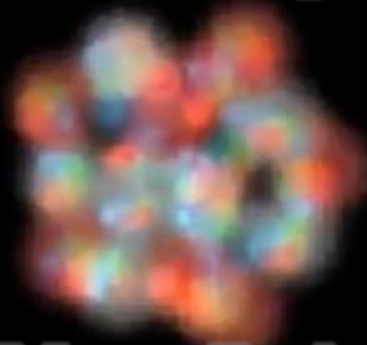


(Nuclear)Astrophysics

Nuclear structure, Equation of State (EoS)
Superfluidity & Superconductivity
Reaction rates, Fission fragments, ...

Neutron-star structure, Star quakes, GW
Pulsar glitches, Cooling
Stellar evolution, Nucleosynthesis, ...

Microscopic



Nuclear Many-Body Problem



10 fm

Hierarchy of Scales in the Universe

Neutron stars, NS merger, nucleosynthesis, GW, ...



Macroscopic



(Nuclear)Astrophysics

Our Mission:

Nuclear structure, Fermi
Superfluid
Reaction rates

To establish a concrete microscopic foundation of macroscopic models

Equation of state, Star quakes, GW
Cooling
Nucleosynthesis, ...



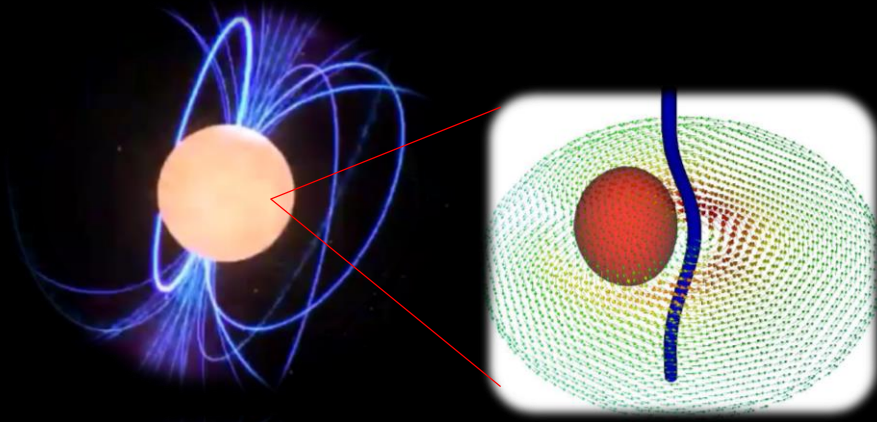
Nuclear Many-Body Problem



10 fm

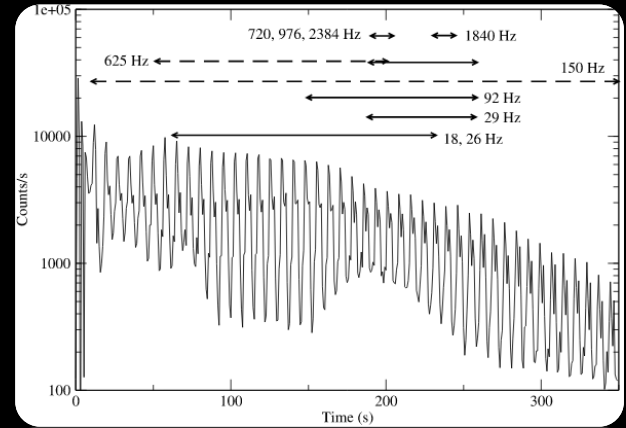
Today, I will mainly talk about:

**Dynamics of quantum vortices
of superfluid neutrons**



↔ Neutron-star glitches

**Time-dependent band theory for
the inner crust of neutron stars**



↔ Quasi-periodic oscillations

..and some related topics of nuclear physics

From **quarks** to atomic nuclei













Standard model of the elementary particles

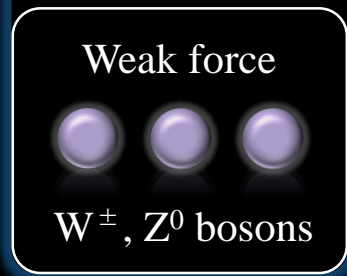
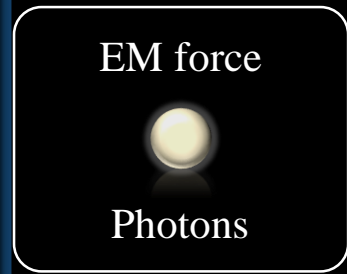
- ✓ Elementary particles: fundamental particles **without structure**
- ✓ Four forces: **strong, weak, electromagnetic**, and gravitational forces
- ✓ Particle physics explores an **ultimate theory** of the Universe



Fermions

Gauge bosons

	1st gen.	2nd gen.	3rd gen.	Charge
Quarks	 <i>u</i> up	 <i>c</i> charm	 <i>t</i> top	$+\frac{2}{3}$
	 <i>d</i> down	 <i>s</i> strange	 <i>b</i> bottom	
Leptons	 <i>e</i> electron	 μ muon	 τ tau	-1
	 ν_e electron neutrino	 ν_μ muon neutrino	 ν_τ tau neutrino	



From quarks to hadrons

Temperature

The QCD phase diagram

Early universe

LHC

✓ Exploring the evolution of the Universe through high-energy nuclear experiments

Quark Gluon Plasma (QGP)

RHIC

Hadrons

✓ Quark matter in NS core?
→ color superconductivity

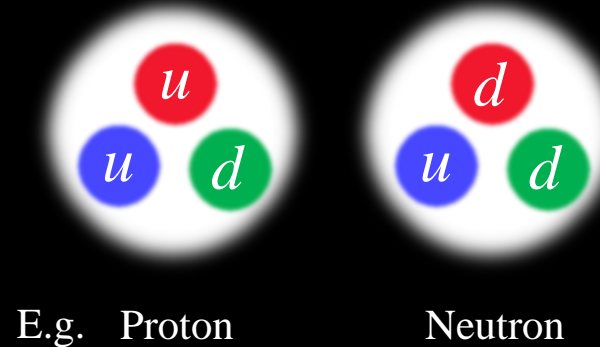
Atomic Nuclei

Neutron Stars

Density

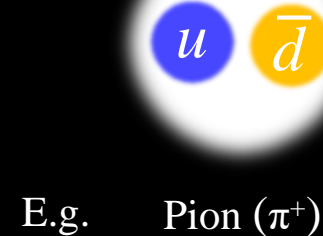
✓ **Hadrons:** composite particles of quarks

Baryons



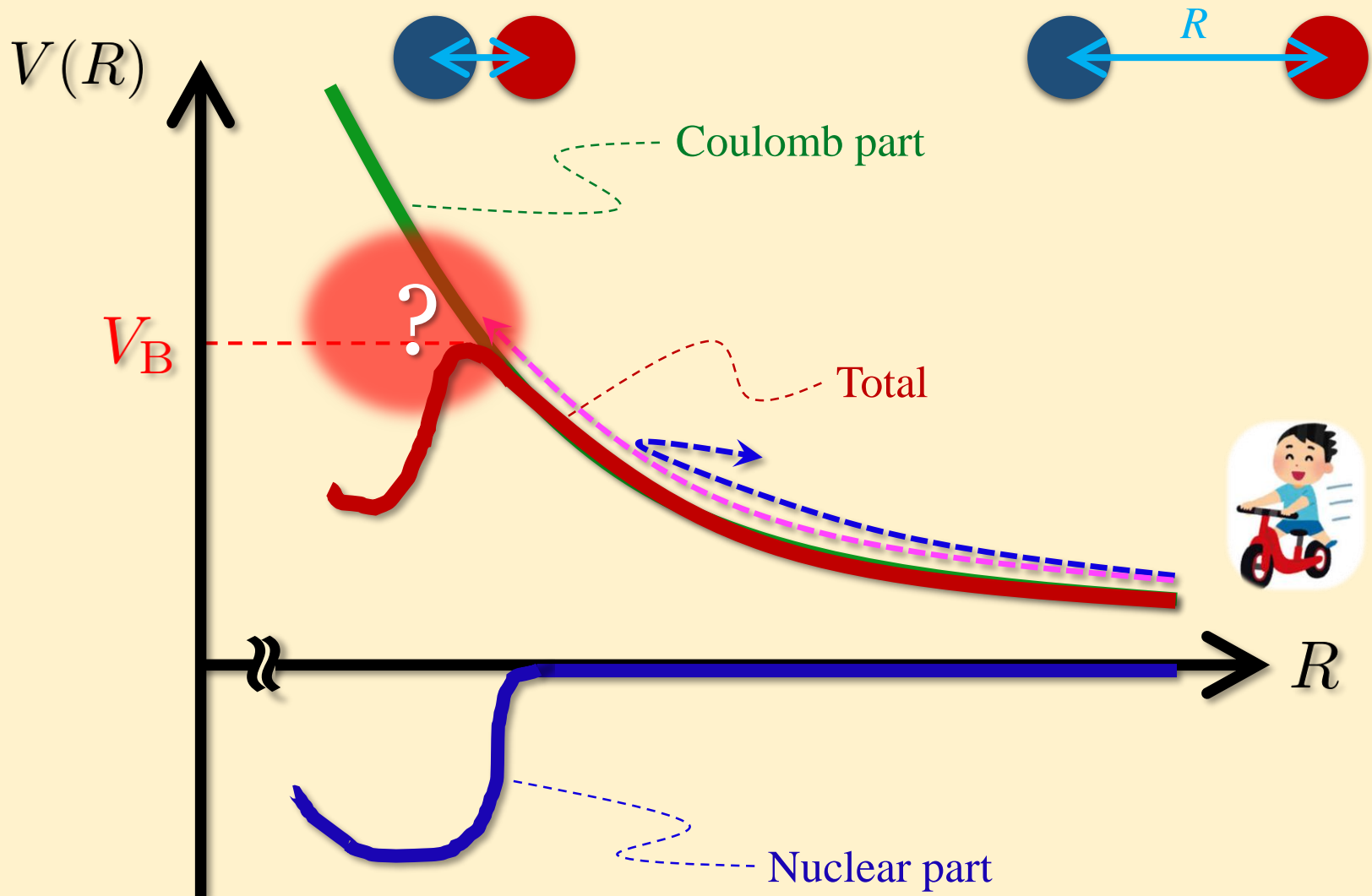
➤ A nucleon (proton or neutron) is a baryon.

Mesons



We collide two nuclei “gently”

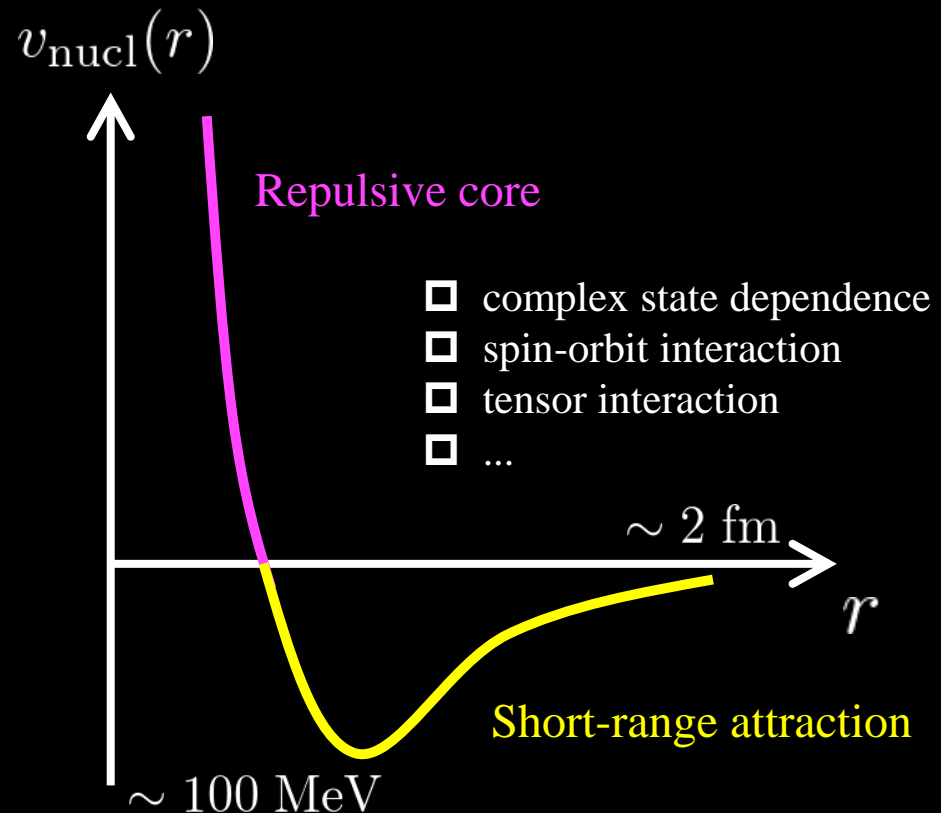
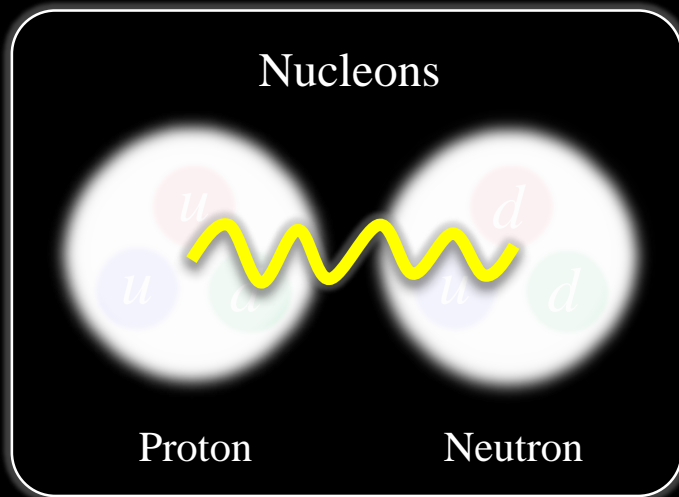
and study quantum many-body dynamics of neutrons and protons



In “low-energy” nuclear physics,
we treat neutrons and protons as building blocks

What we study is:

A quantum many-body problem of fermions interacting through the nuclear force



■ (TD)DFT in a tiny nutshell



A theory which gives us access to the *exact* solution

Equivalent!
(for a special EDF)

$$\hat{H}\Psi(\mathbf{r}_1, \dots, \mathbf{r}_N) = E\Psi(\mathbf{r}_1, \dots, \mathbf{r}_N)$$

Kohn-Sham equation

$$\left[-\frac{\hbar^2}{2m}\nabla^2 + v_{\text{KS}}[\rho(\mathbf{r})] \right] \phi_i(\mathbf{r}) = \varepsilon_i \phi_i(\mathbf{r})$$

$$v_{\text{KS}}[\rho(\mathbf{r})] = \frac{\delta \mathcal{E}[\rho]}{\delta \rho} \quad \rho(\mathbf{r}) = \sum_{i=1}^N |\phi_i(\mathbf{r})|^2$$

EDF

This is the key!

Quantum Many-Body Problem



Energy can also be written as a functional of density

$$E[\rho] = \langle \Psi[\rho] | \hat{H} | \Psi[\rho] \rangle$$

w.f. is a functional of density

P. Hohenberg and W. Kohn, Phys. Rev. B **136**, 864 (1964)

CAUTION!

The existence was proven, but its shape is unknown..

“Inverse Kohn-Sham”

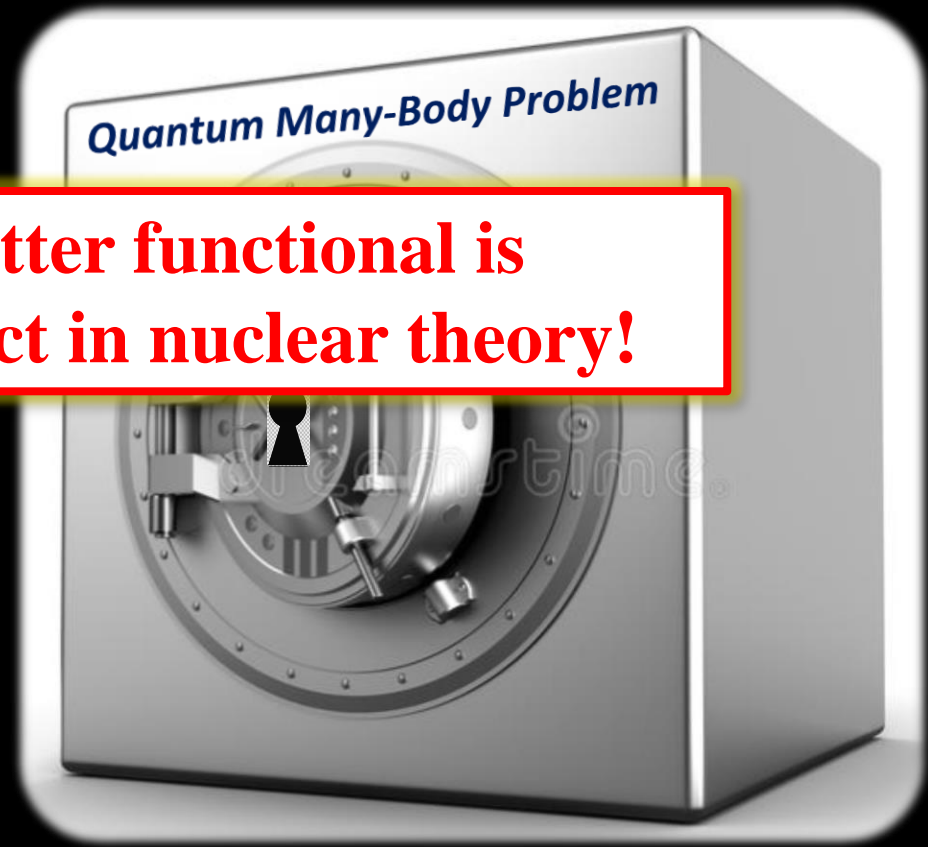


CAUTION!

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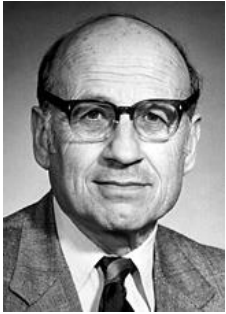
“Inverse Kohn-Sham”

**Developing a better functional is
an important subject in nuclear theory!**

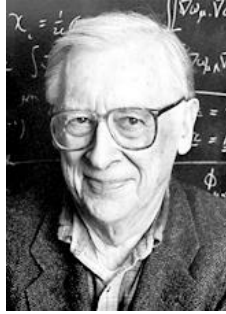


Great Success of the Density Functional Theory

The Nobel Prize in Chemistry 1998



Walter Kohn

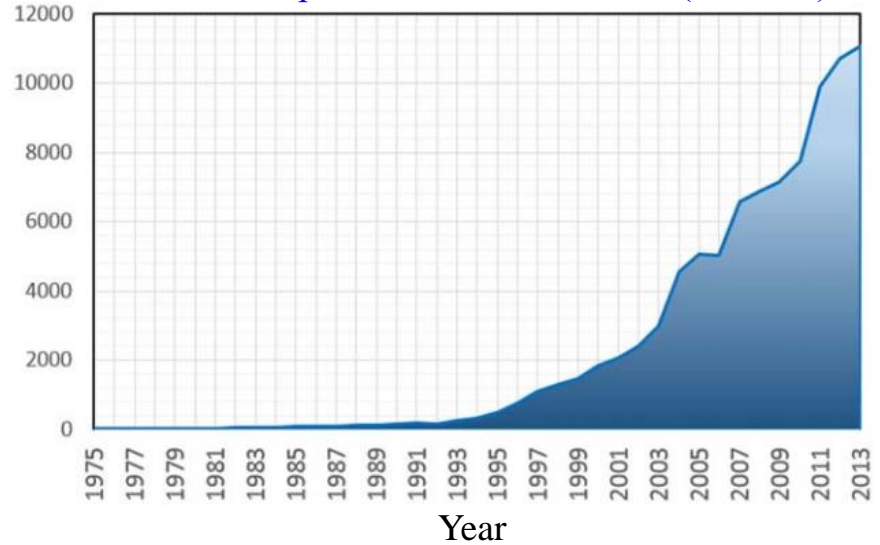


John Pople



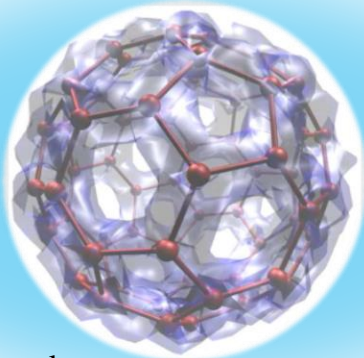
©<https://www.nobelprize.org>

Number of publications with “DFT” (till 2013)



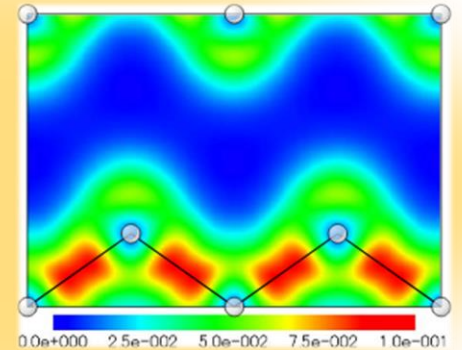
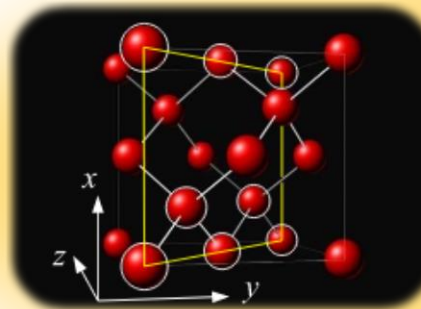
A. Galano and J.R. Alvarez-Idadoy, *J. Compt. Chem.* **35**, 2019 (2014)

Fullerene: C_{60}



C-Z. Gao et al.,
J. Phys. B: At. Mol. Opt. Phys. **48**, 105102 (2015)

Si crystal

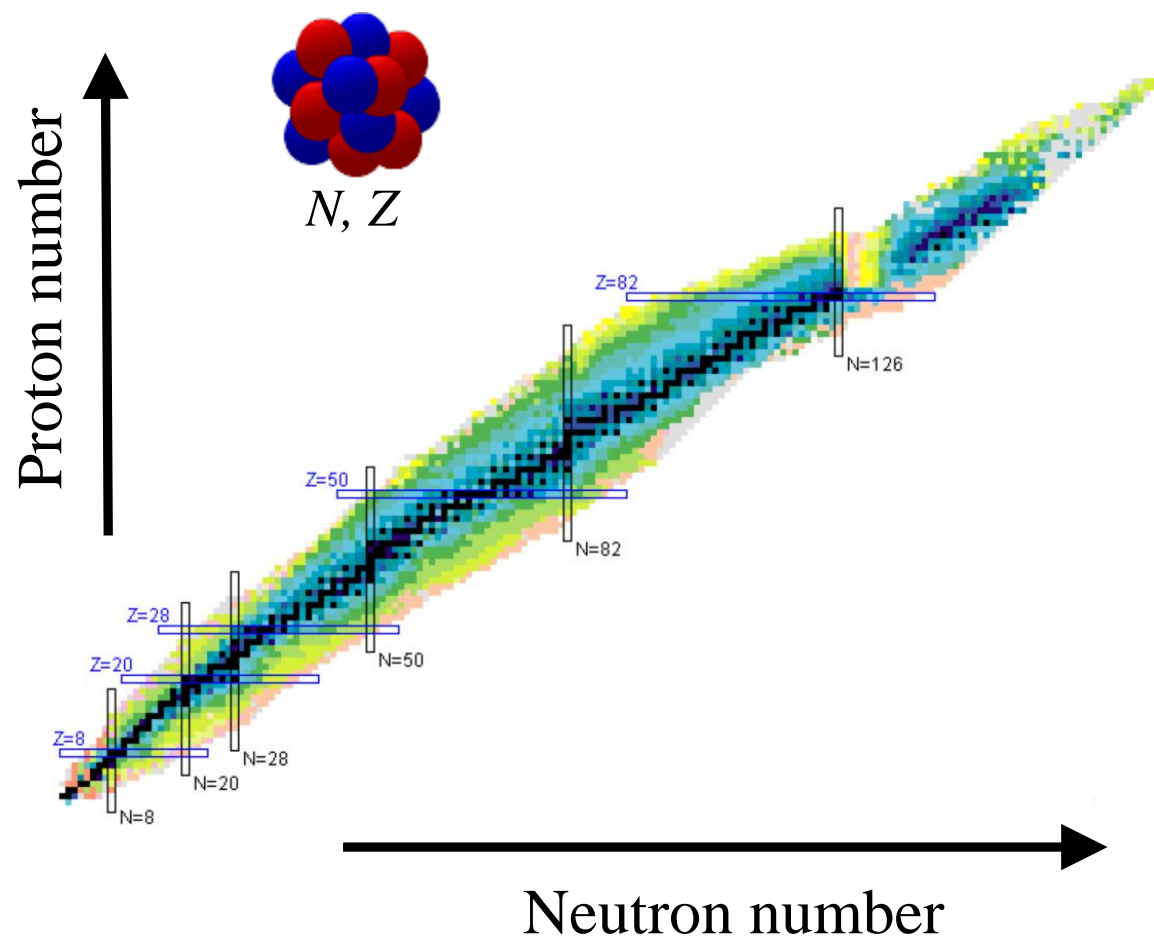


Y. Shinohara, K. Yabana, Y. Kawashita, J.-I. Iwata, T. Otobe, and G. F. Bertsch,
Phys. Rev. B **82**, 155110 (2010)

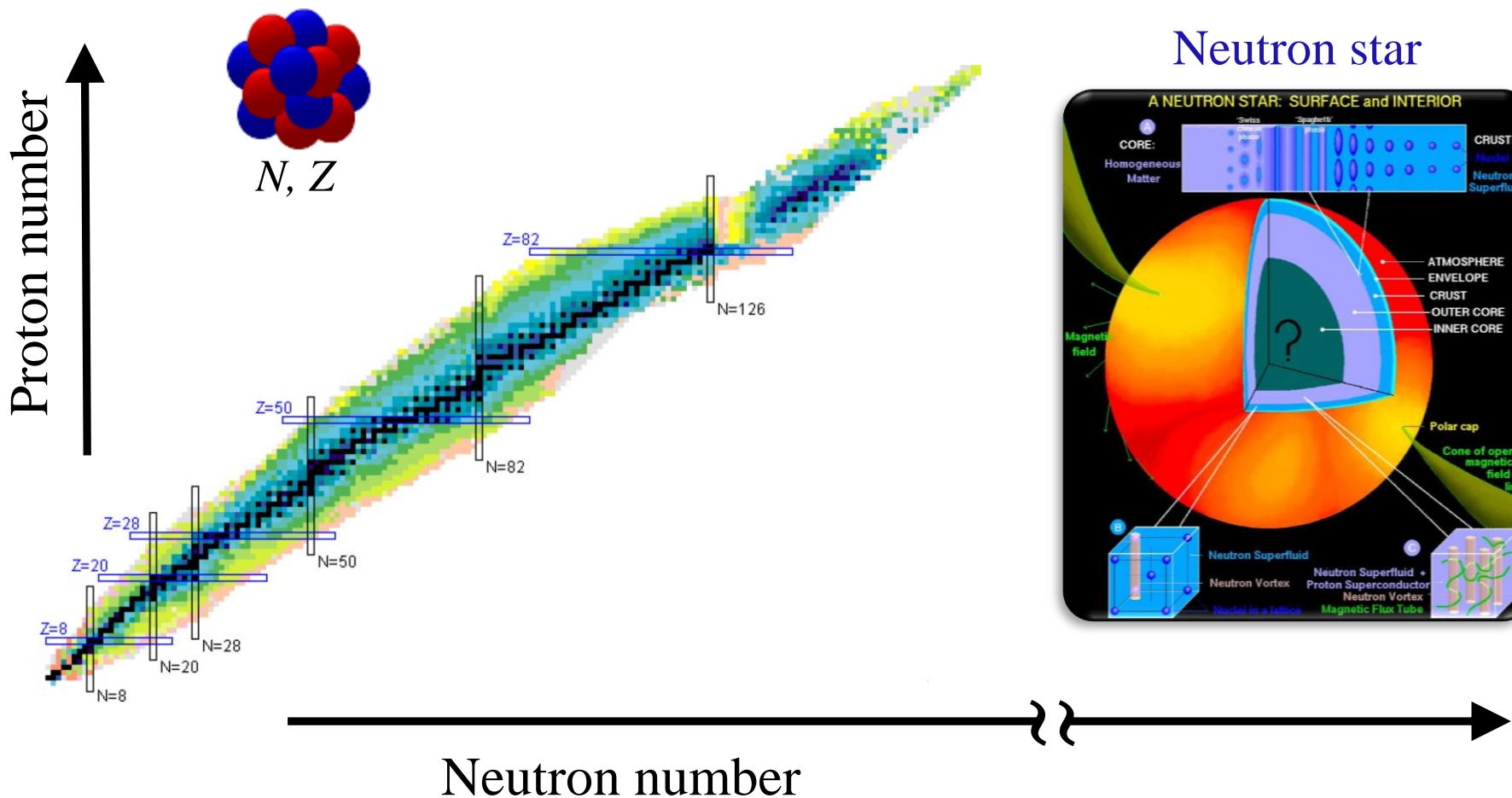
The seminal papers on DFT

- P. Hohenberg and W. Kohn, *Phys. Rev.* **136**, B864 (1964) ➔ **19,015 citations!**
- W. Kohn and L.J. Sham, *Phys. Rev.* **140**, A1133 (1965) ➔ **24,384 citations!**

All nuclei can be described with a *single* EDF

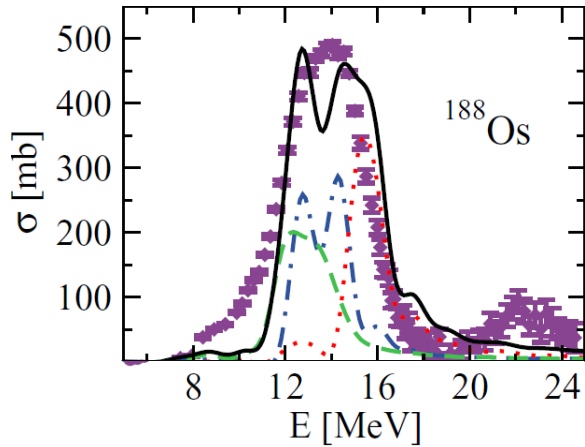


All nuclei can be described with a *single* EDF

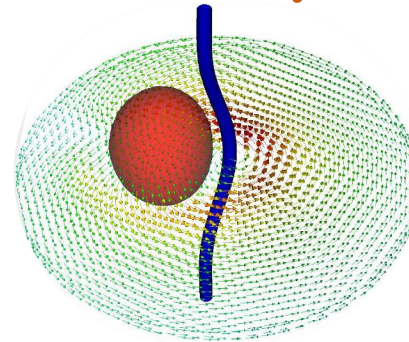


TDDFT is a versatile tool!!

IVGDR

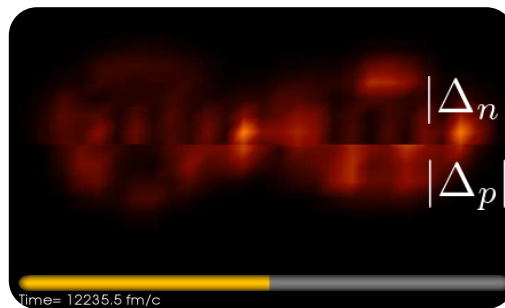


Vortex-nucleus dynamics

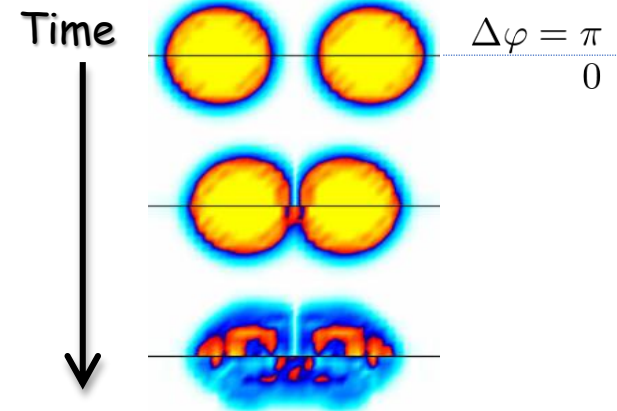


Phys. Rev. C **84**, 051309(R) (2011)
I. Stetcu, A. Bulgac, P. Magierski, and K.J. Roche

Induced fission of ^{240}Pu



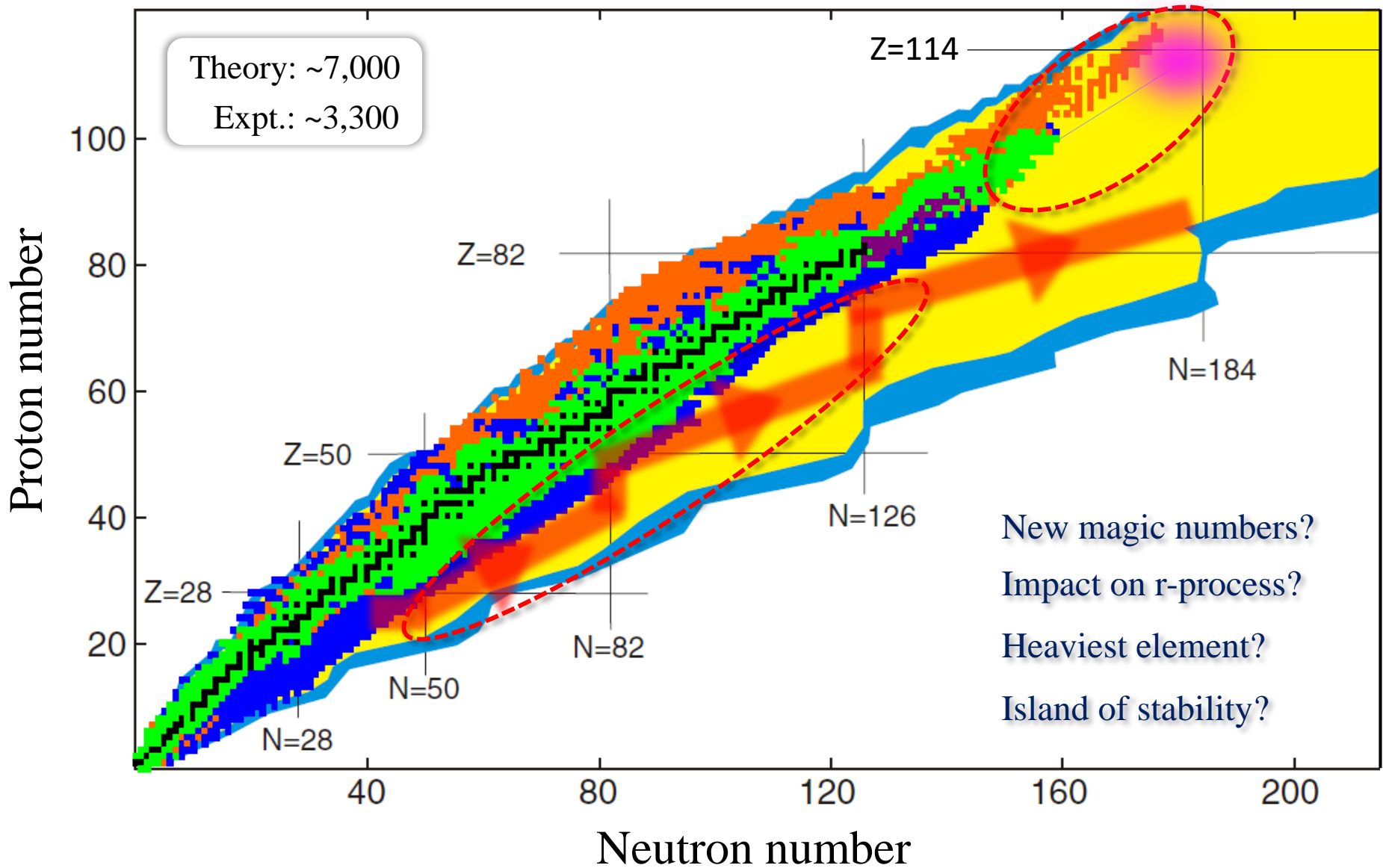
Low-energy heavy-ion reactions



At the frontiers in nuclear physics I:

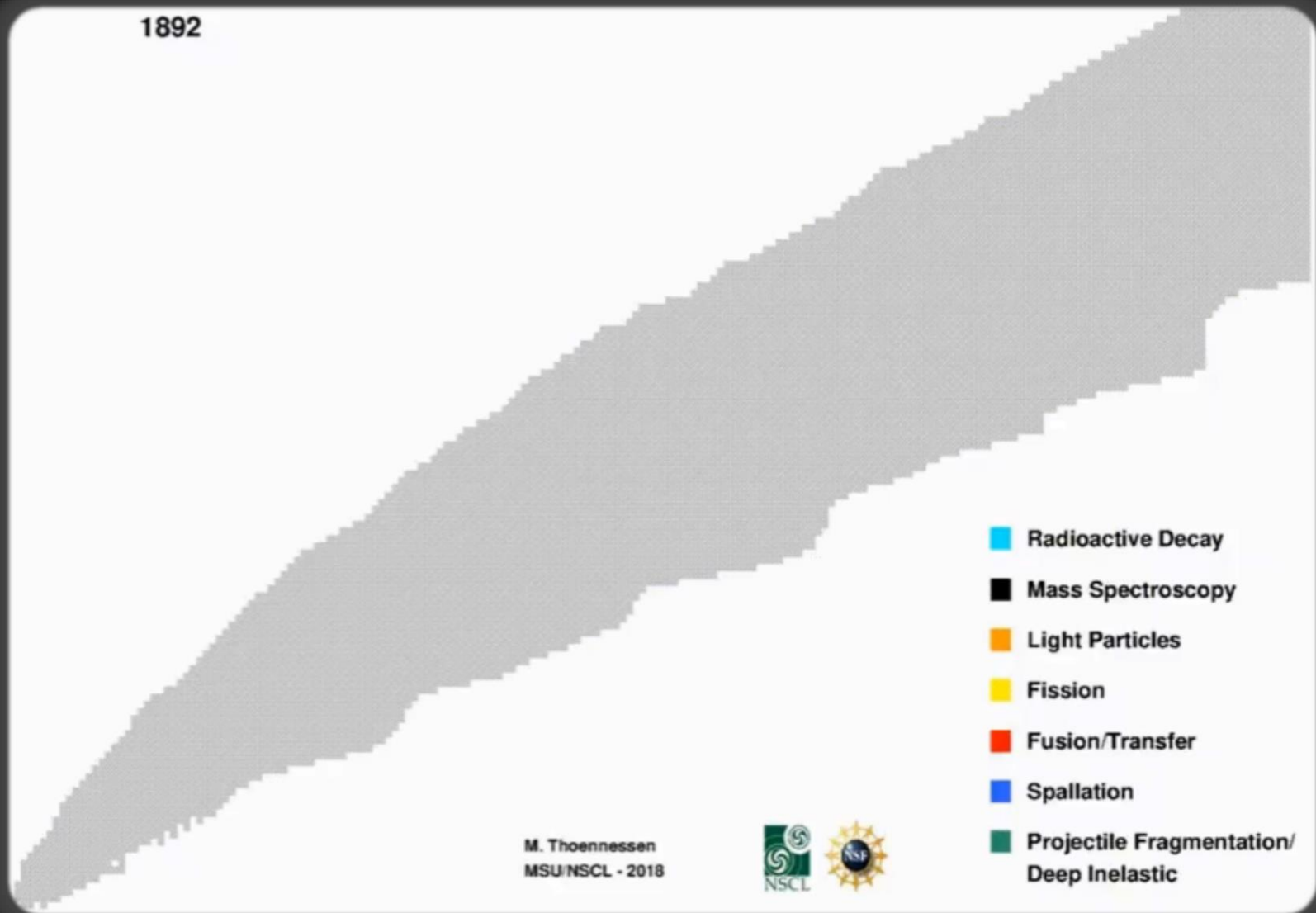
Voyage towards the limit of nuclear existence

The “map” of atomic nuclei: the nuclear chart (a.k.a. the Segrè chart)



Movie from “Discovery of Nuclides Project” by Michael Thoennesen

<https://people.nsl.msui.edu/~thoennes/isotopes/>



The continent of stability has been well explored..



Now we are sailing towards the edge of the nuclear landscape..



Stable nuclei: 288

Experiment: ~3300

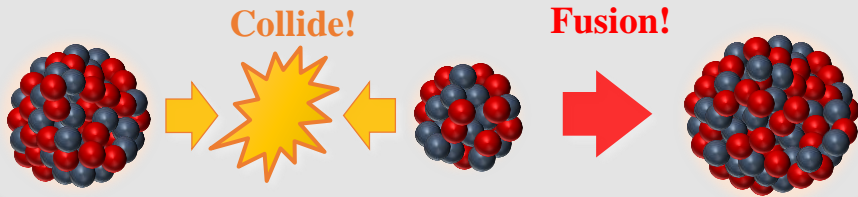
Theory: ~7000-10000

- ❑ drip lines
- ❑ shell structure
- ❑ deformation
- ❑ skin, halo
- ❑ nuclear matter properties
- ❑ nucleosynthesis
- ❑ ...

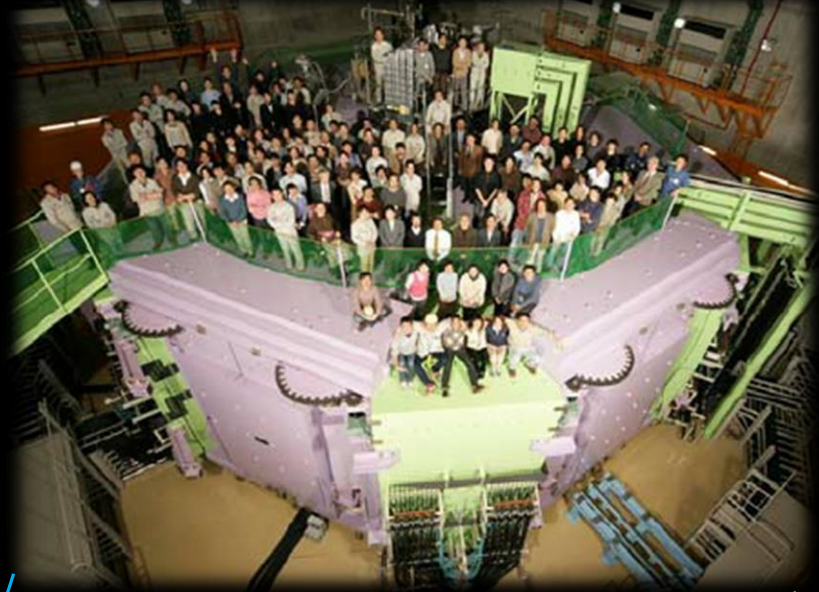


Nucleosynthesis at accelerator facilities

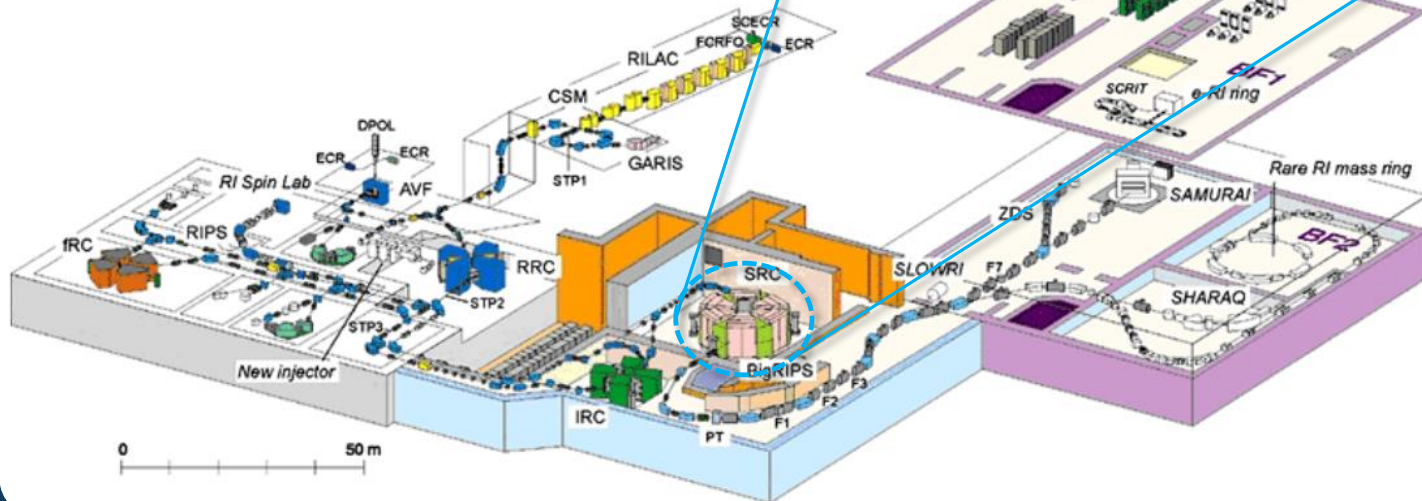
Superconducting Ring Cyclotron (SRC)



The world-leading factory of unstable nuclei!



RI Beam Factory (RIBF) @RIKEN



At the frontiers in nuclear physics II:

Physics of “Neutron Stars”

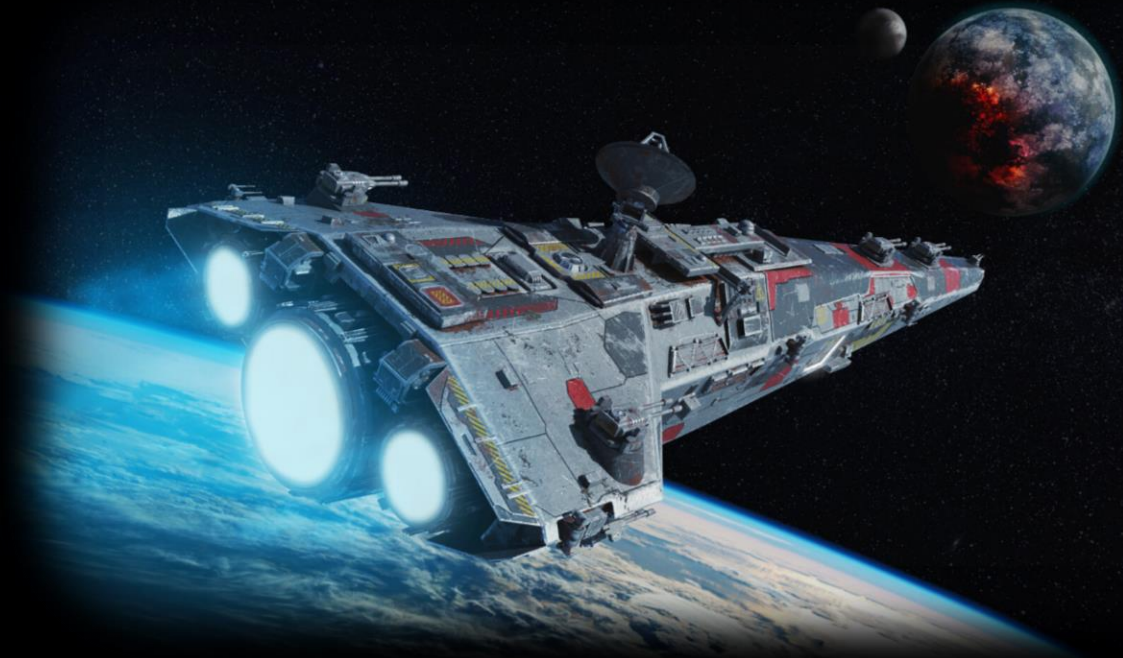
→ also relevant to:

**condensed matter physics, solid-state physics,
as well as astrophysics**

Now we are sailing towards the edge of the nuclear landscape..



*Let's leave the planet
of finite nuclei!*

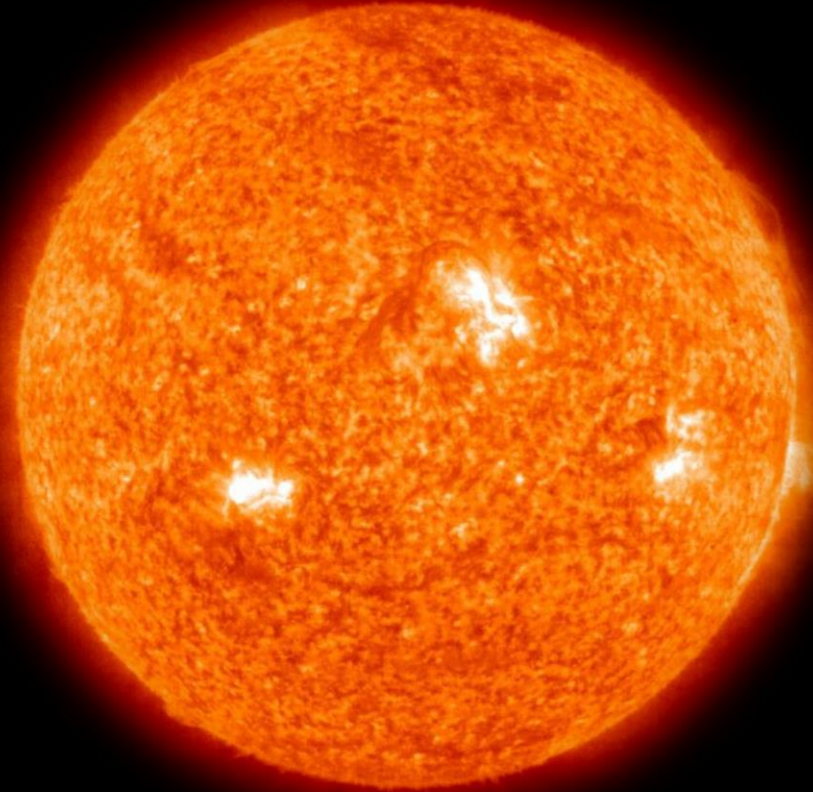




From nuclei to **neutron stars**

The Sun

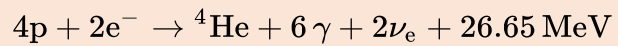
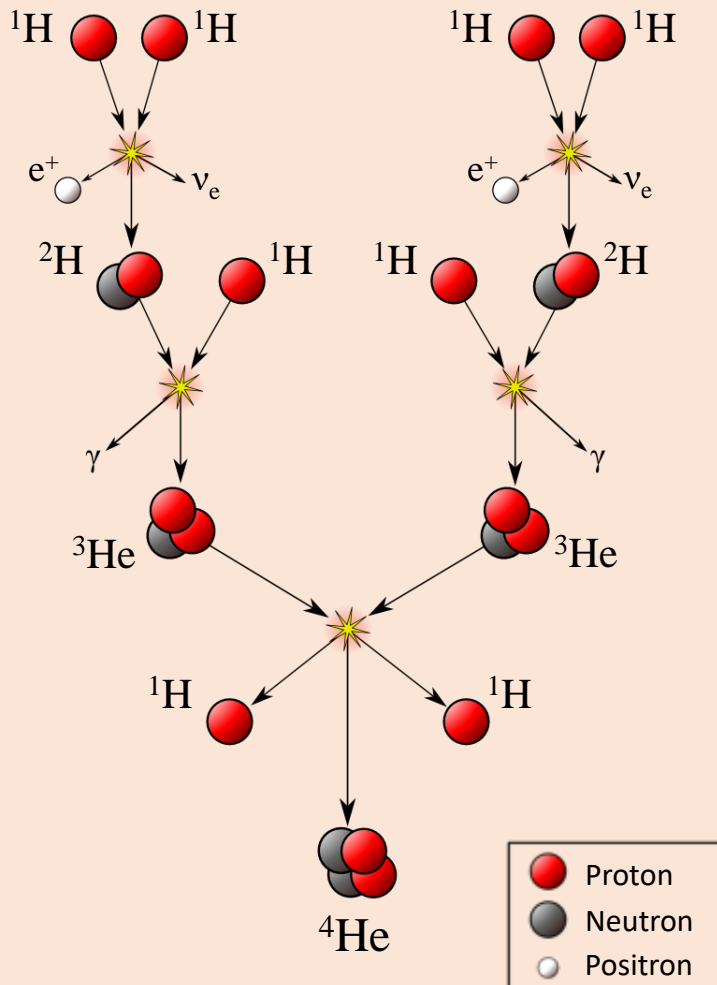
- Radius: $\sim 7 \times 10^8 \text{m}$ (~ 109 times bigger than Earth)
- Mass: $\sim 2 \times 10^{30} \text{kg}$ (~ 330 thousands times heavier than Earth)
- Central temp.: ~ 10 million $^{\circ}\text{C}$
- Surface temp.: ~ 5000 $^{\circ}\text{C}$



Stars shine due to nuclear fusion reactions

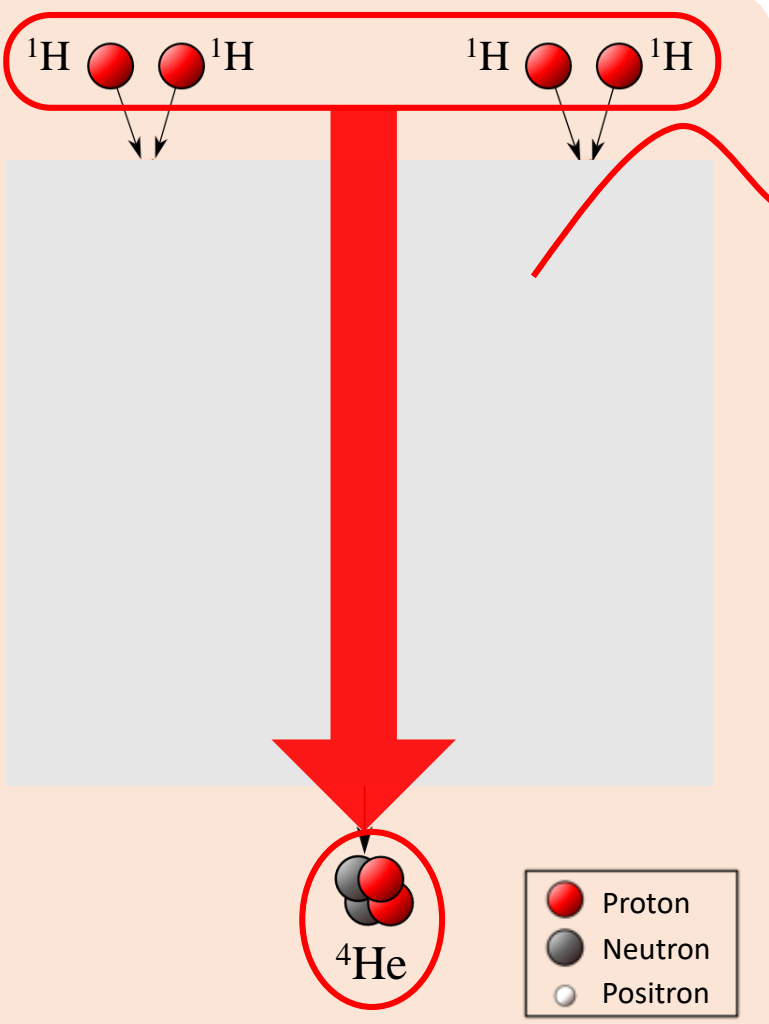
Energy source of the Sun: Nuclear fusion

Proton-proton (p-p) chain

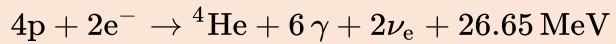
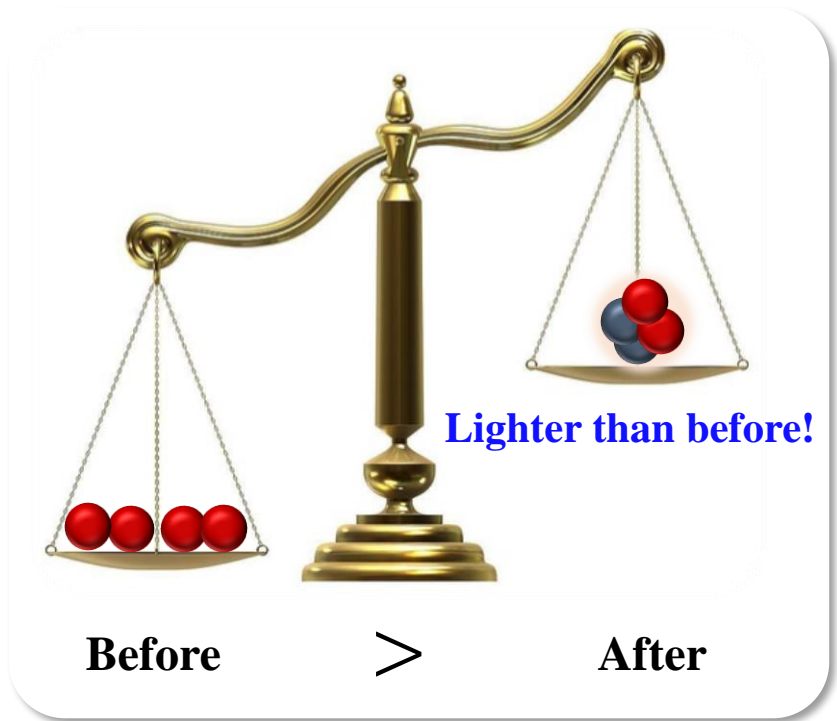


Energy source of the Sun: Nuclear fusion

Proton-proton (p-p) chain

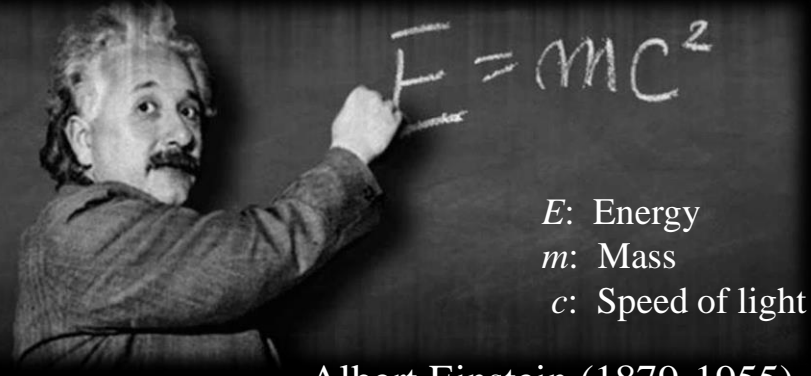


What happened:
Four protons became a helium nucleus!



Energy source of the Sun: Nuclear fusion

Equivalence between mass and energy

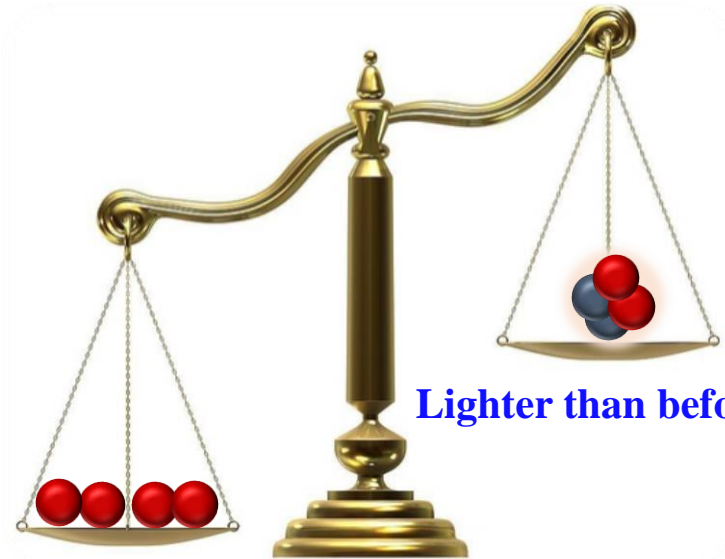


Albert Einstein (1879-1955)

E : Energy
 m : Mass
 c : Speed of light

What happened:

Four protons became a helium nucleus!



Lighter than before!

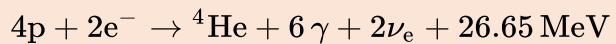
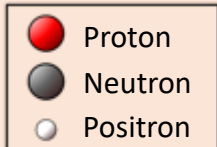
Before

>

After

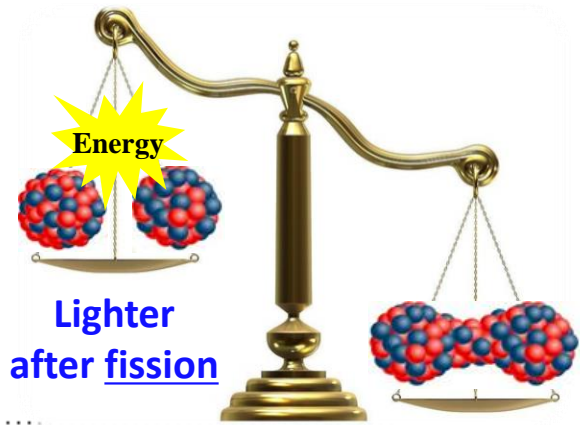
This reaction generates energy
(exothermic)

Energy

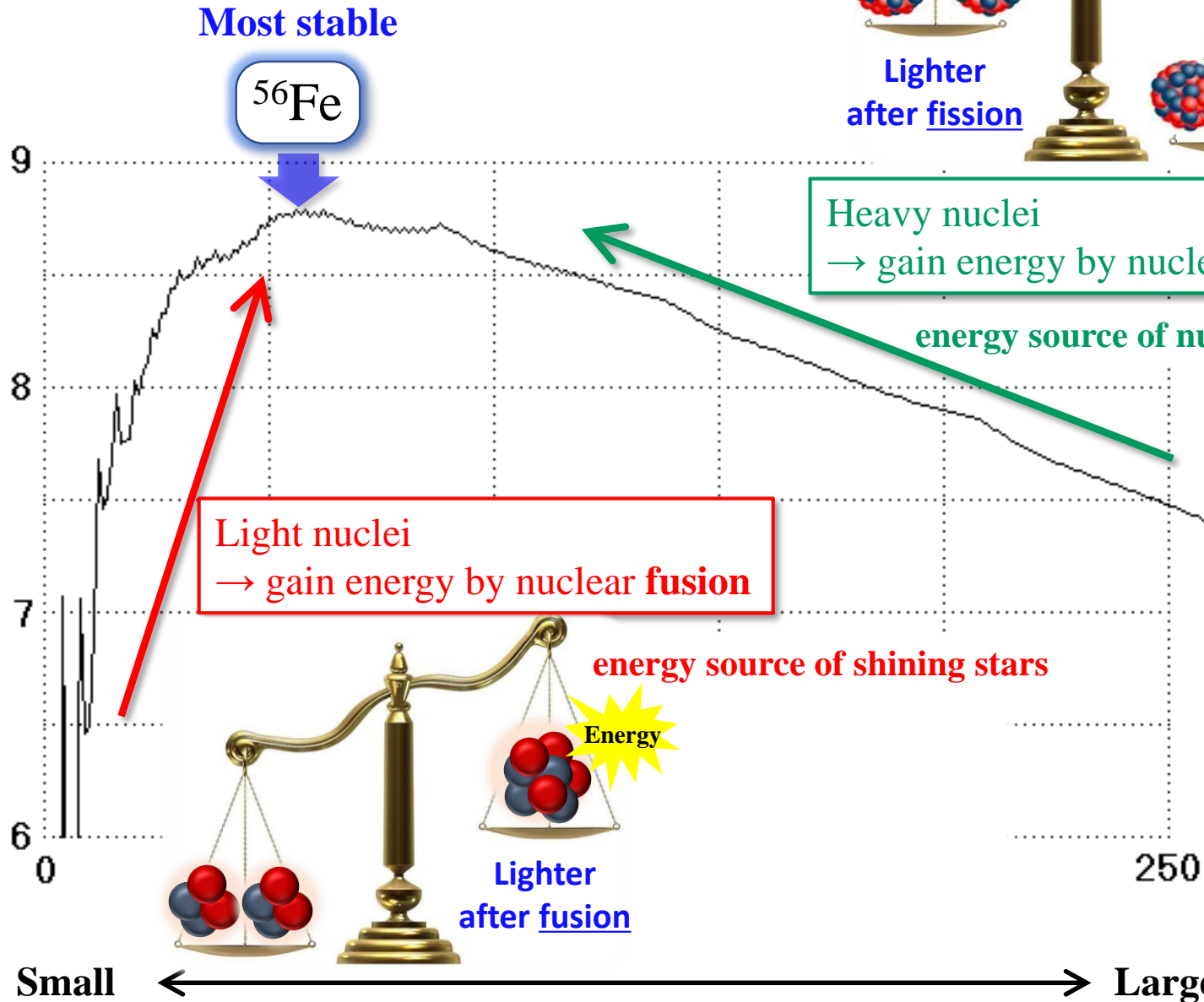


Quantum tunneling allows for
overcoming the Coulomb barrier!

Stability of Nuclei



Binding energy per nucleon (MeV)



Most stable

^{56}Fe

Heavy nuclei
→ gain energy by nuclear fission

Light nuclei
→ gain energy by nuclear fusion

energy source of nuclear plants

energy source of shining stars

Lighter after fusion

Lighter after fission

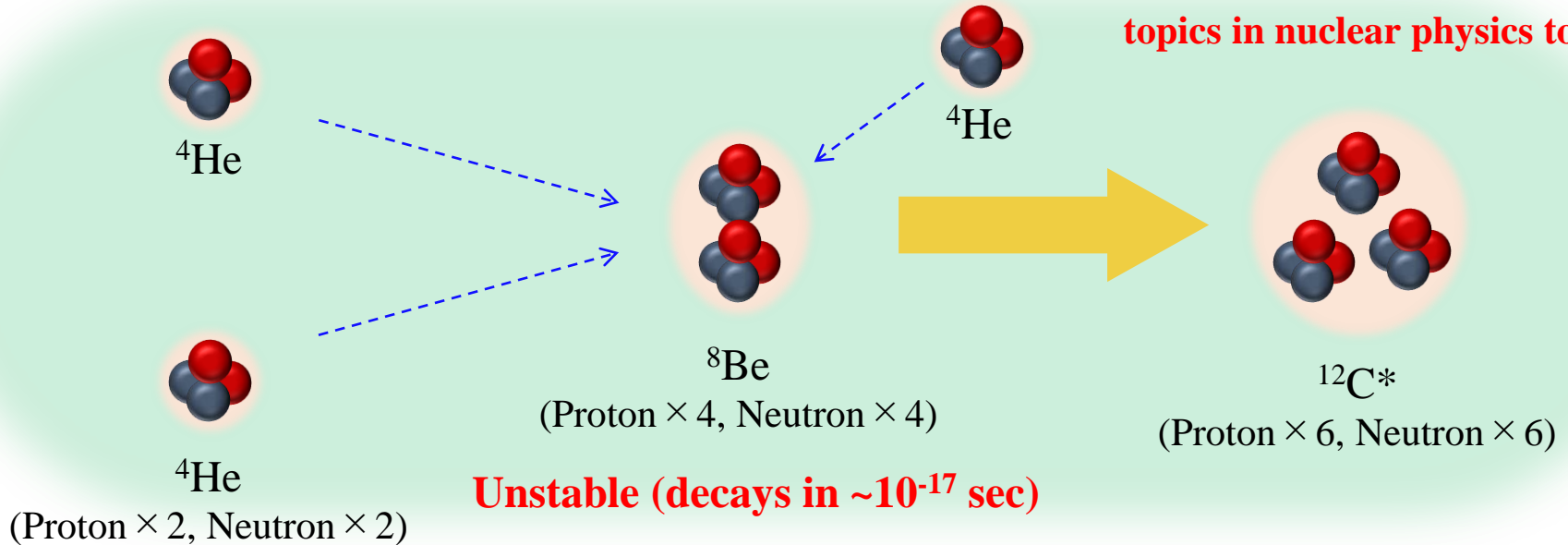
Small

Large

In massive stars, reactions should proceed further, but..

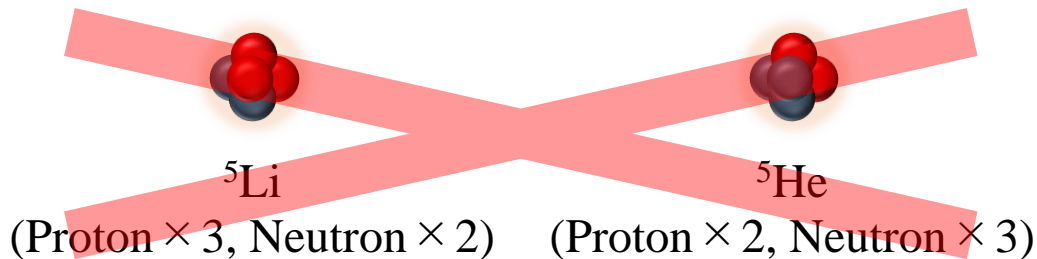
“Triple-alpha reaction”

➤ Cluster structure is one of hot topics in nuclear physics today!



Hoyle predicted an excited state with 3α cluster structure near the 3α threshold!

Neither ${}^5\text{Li}$ nor ${}^5\text{Be}$ can be an alternative



Unstable (decays in $\sim 10^{-22}$ sec)

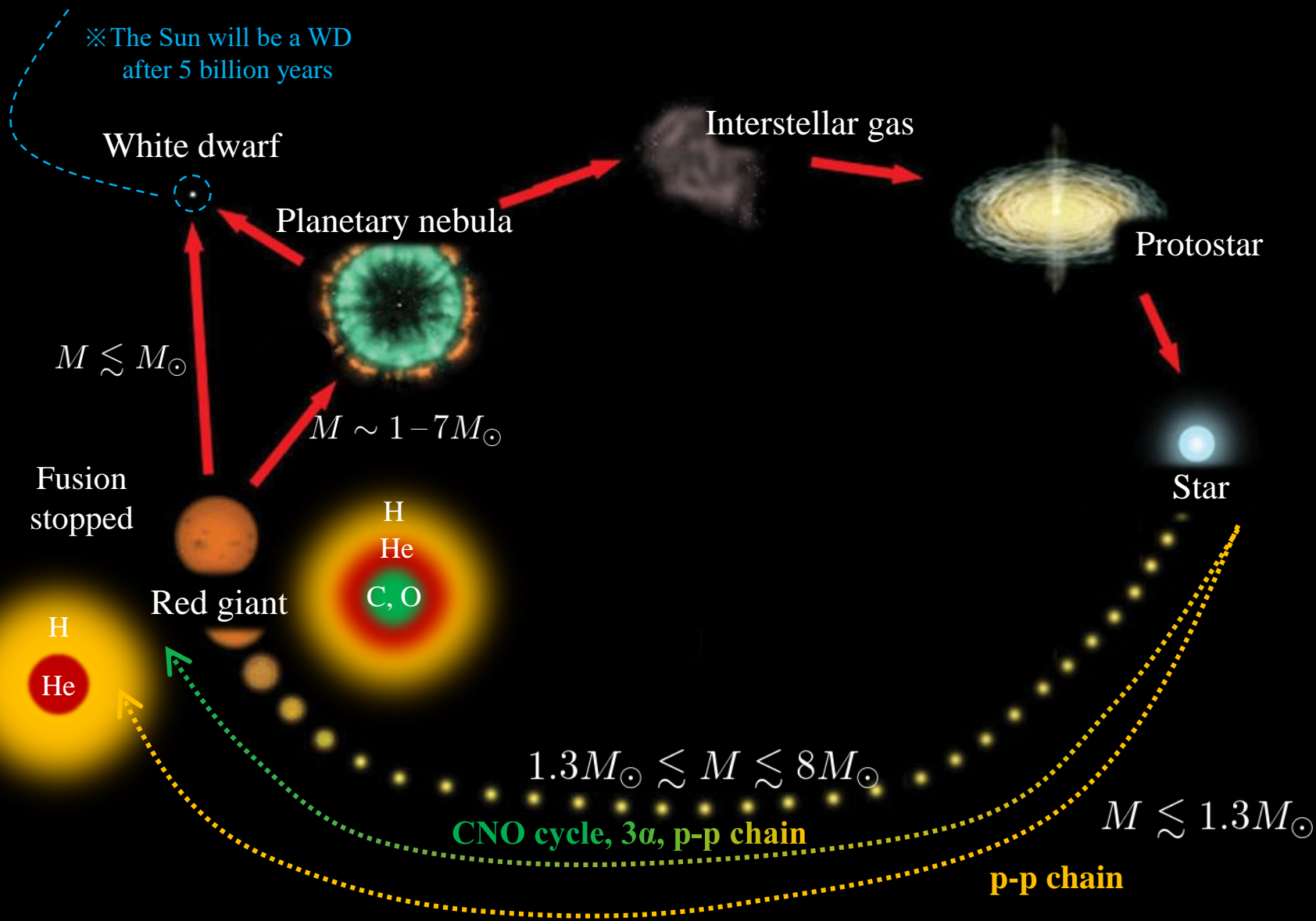


F. Hoyle (1915-2001)

Life cycle of a **light** star

The size is like Earth,
but with the solar mass

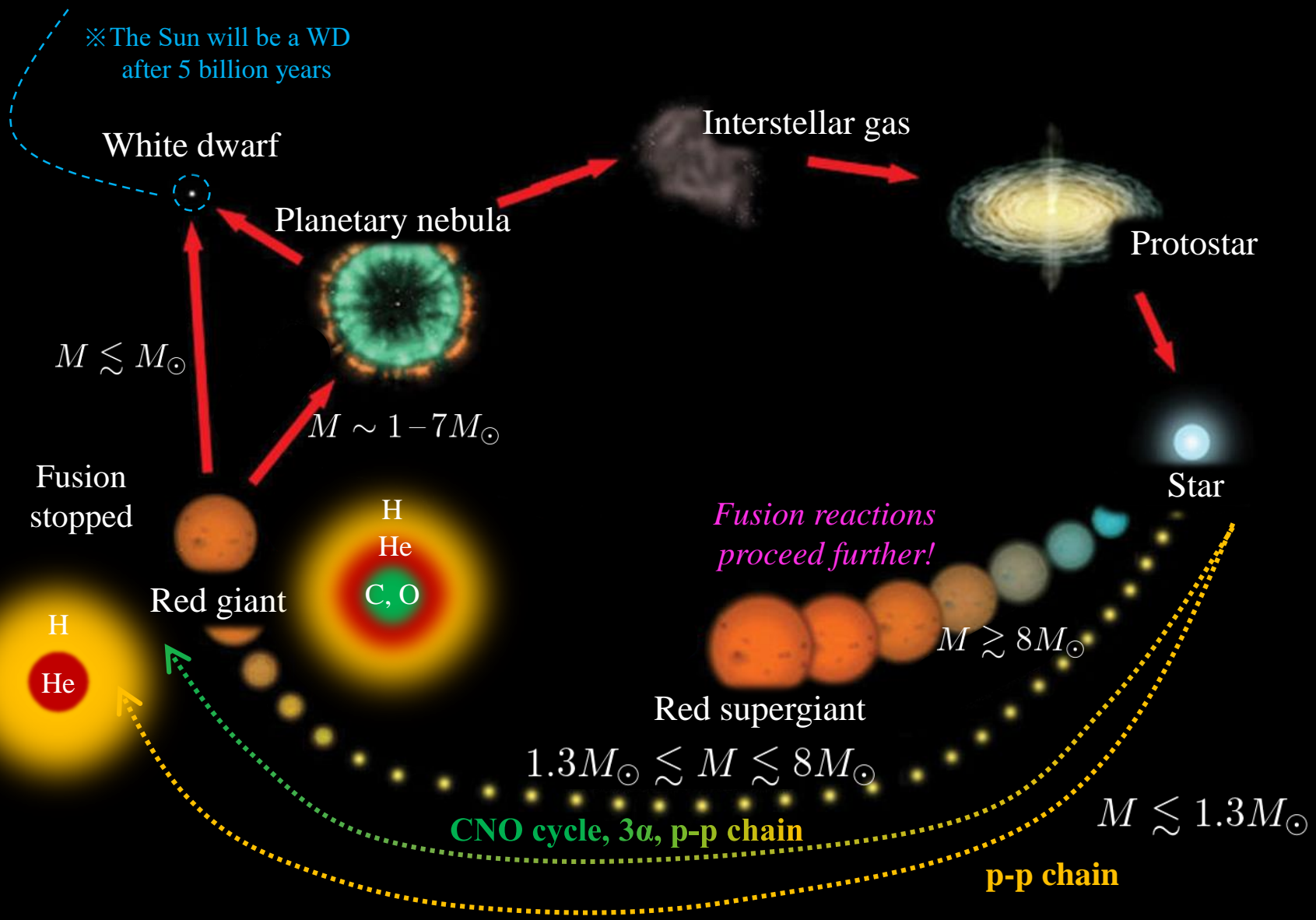
※ The Sun will be a WD
after 5 billion years



Life cycle of a massive star

The size is like Earth, but with the solar mass

※ The Sun will be a WD after 5 billion years

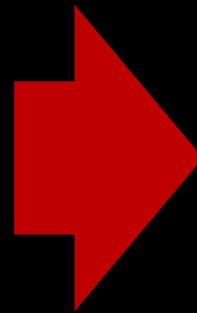
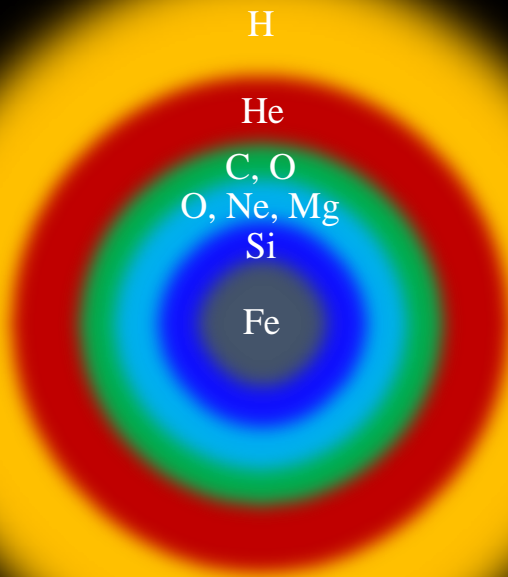


The fate of a massive star

Nuclear reactions:

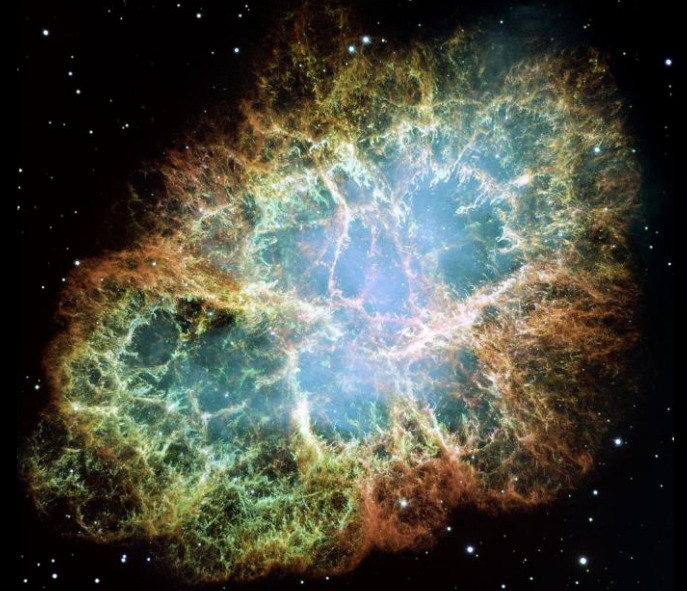


“Onion structure”



After forming the iron core...

- no more fuel
- gravitational collapse
- supernova explosion

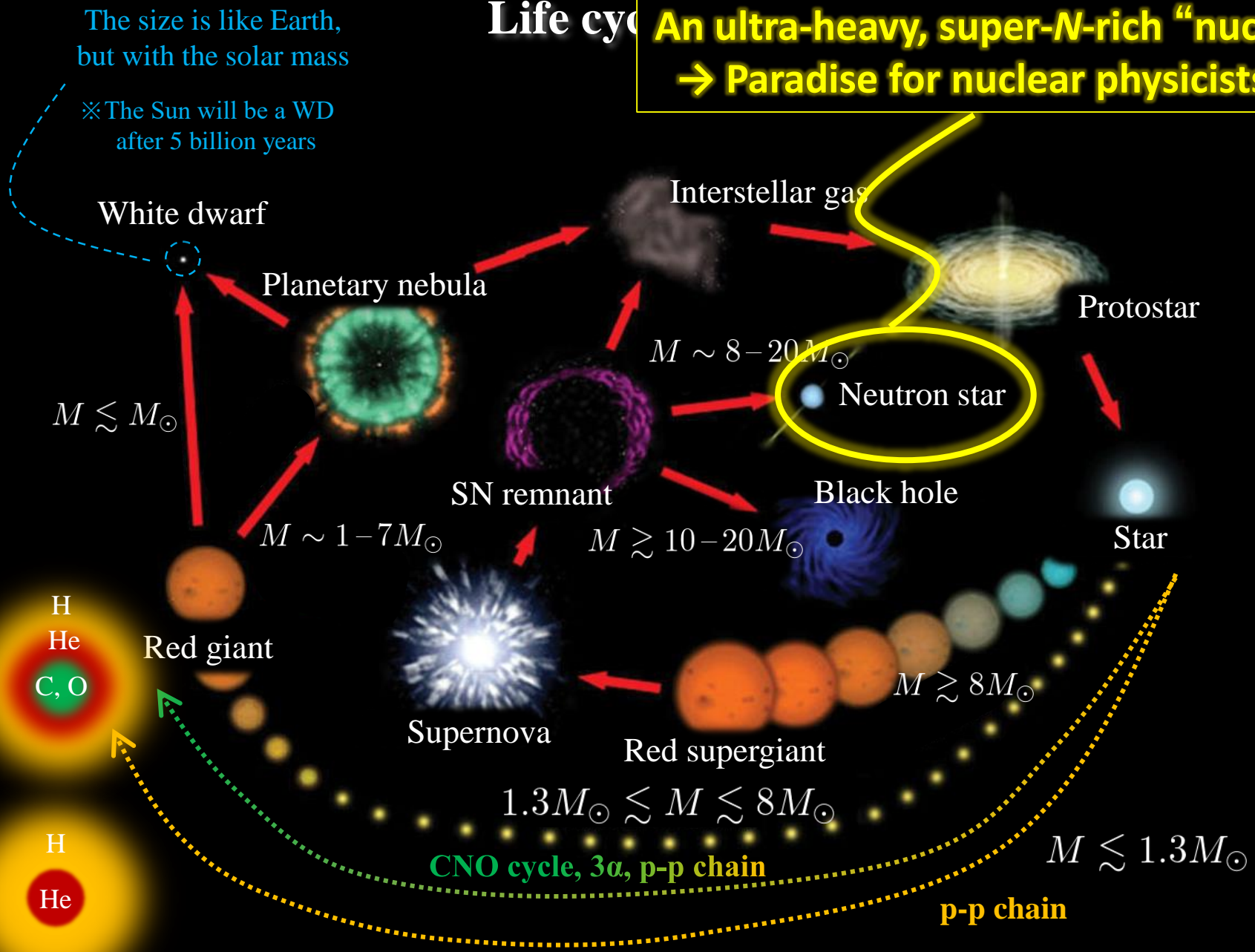


The Crab Nebula
Remnant of the SN in 1054



Life cycle

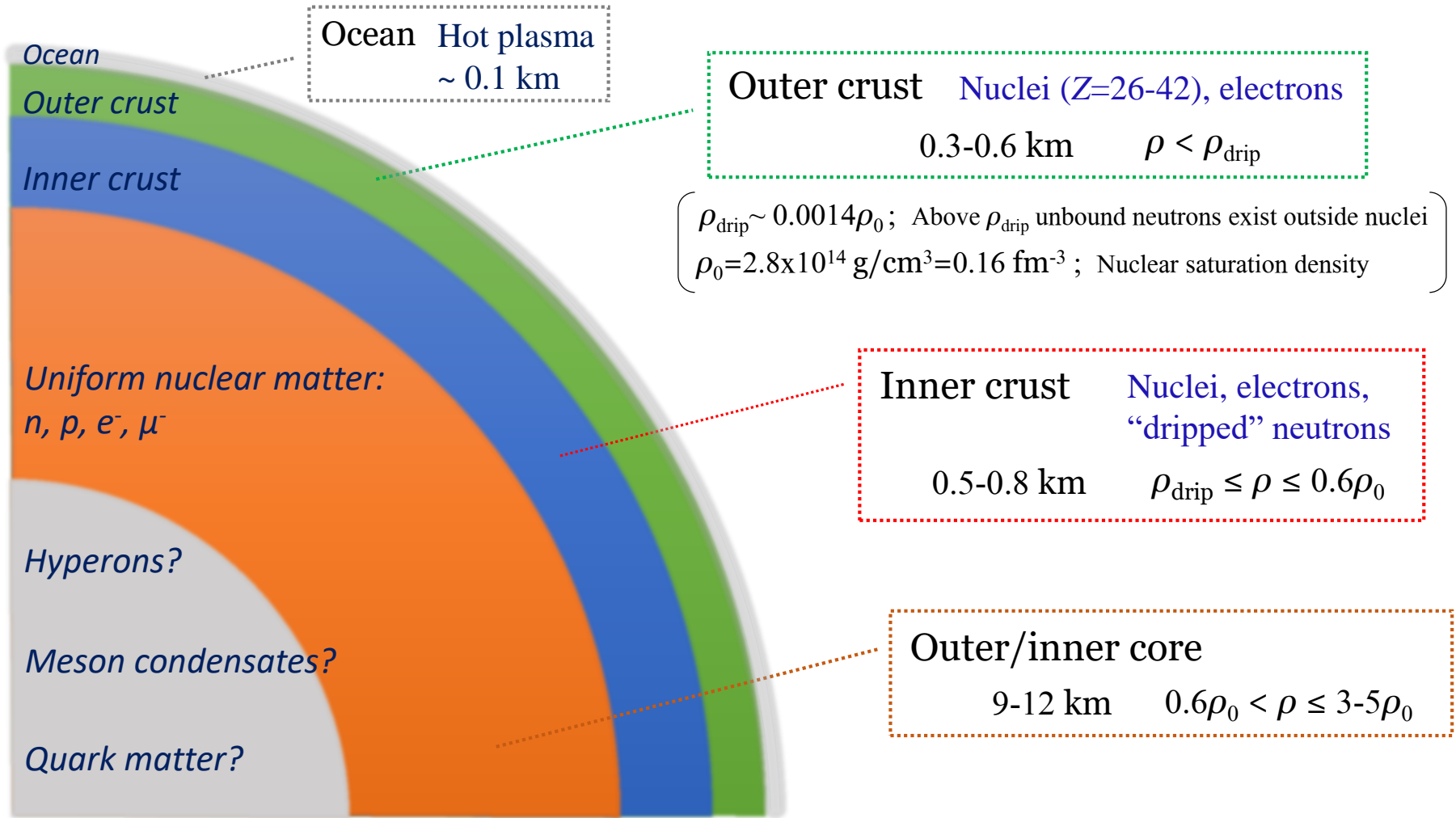
An ultra-heavy, super-N-rich "nucleus"
 → Paradise for nuclear physicists! :D



What's inside a neutron star?

Neutron star is a great playground for nuclear physicists

- ✓ It offers extreme situations which can not be realized in terrestrial experiments



In an outer (low-density) region of neutron stars,
nuclear matter is not actually homogeneous

The nuclear interaction “clusterize” neutrons and protons,
akin to finite nuclei, which form a Coulomb lattice
(*i.e.* a crystal, like a solid)

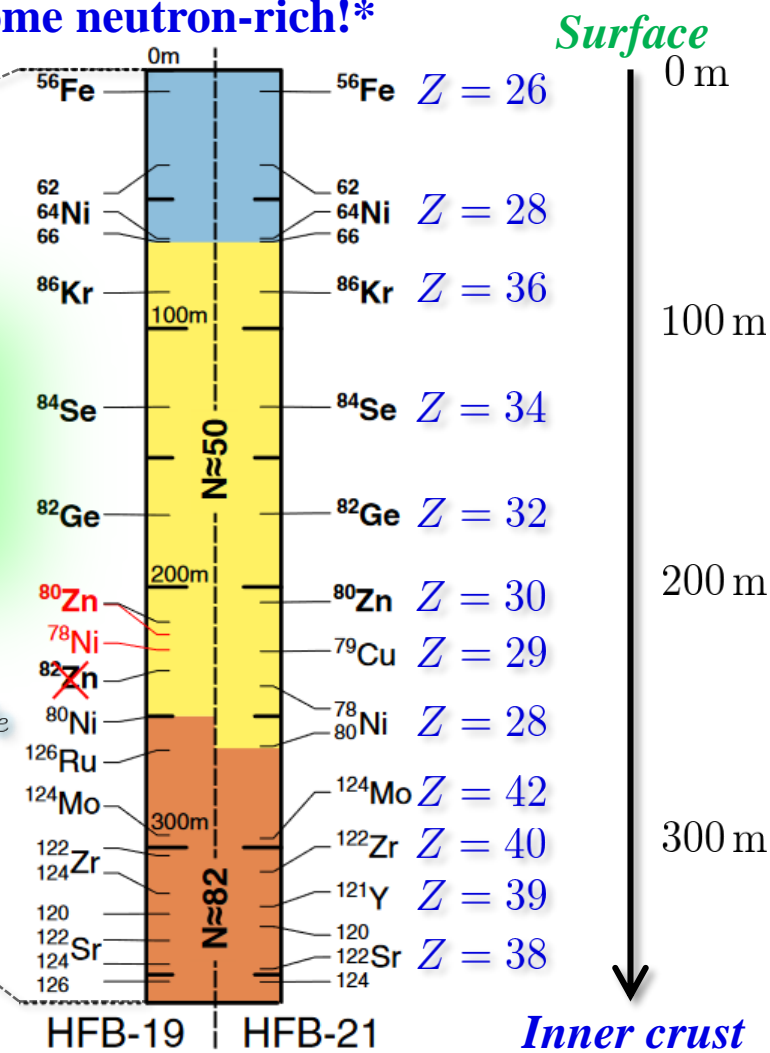
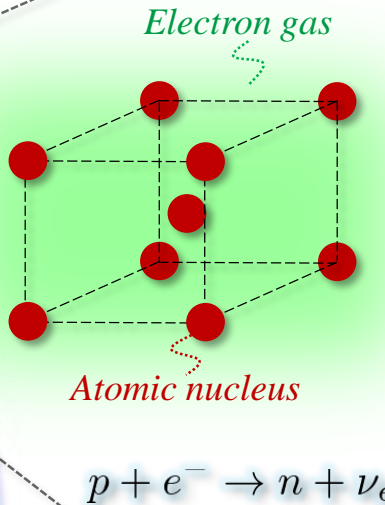
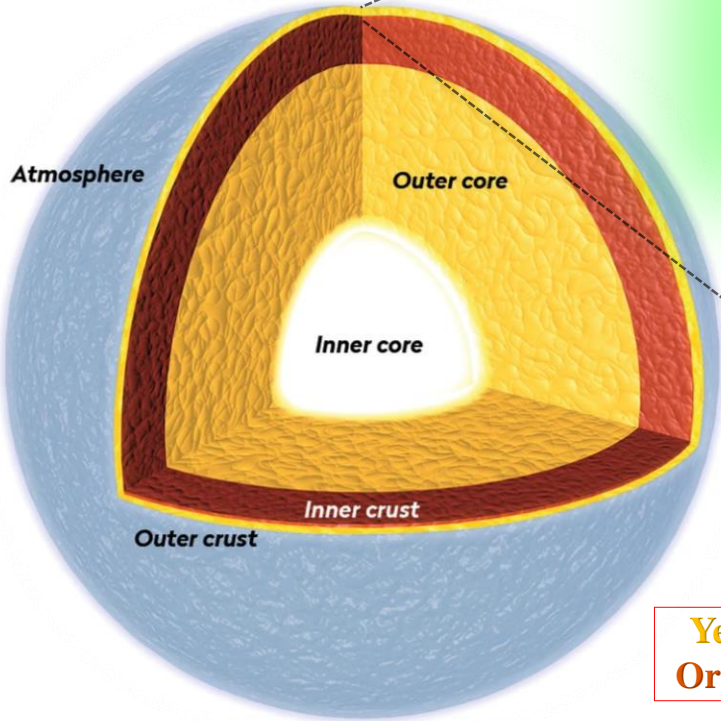
Let’s see: from the outer crust to the inner crust

Structure of the outer crust is “similar” to that of a white dwarf

but, nuclei are different and become neutron-rich!

Composition of the outer crust:

- ✓ Coulomb lattice of nuclei
- + Electron gas



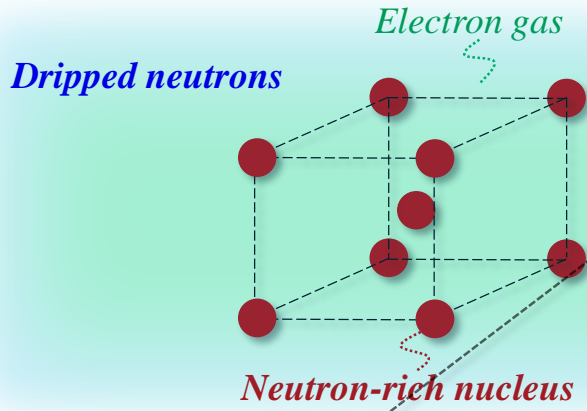
Yellow region: $N \approx 50$
Orange region: $N \approx 82$

Figure: R.N. Wolf et al., Phys. Rev. Lett. **110**, 041101 (2013)

Neutron star: <https://www.skyatnightmagazine.com/space-science/neutron-star/>

Inner crust

In the inner crust, a sea of “dripped neutrons” permeates the Coulomb lattice



✓ Distance between nuclei decreases with increasing density **Outer crust**

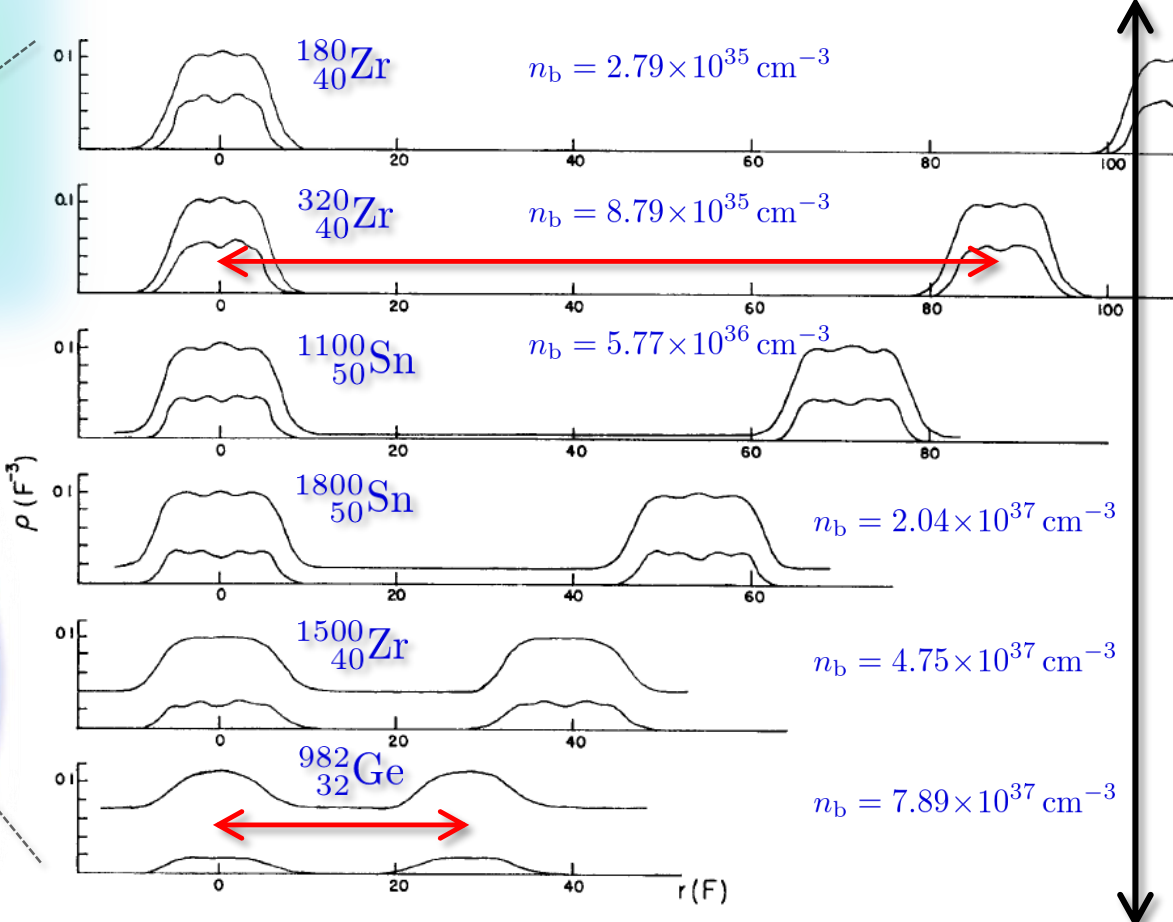
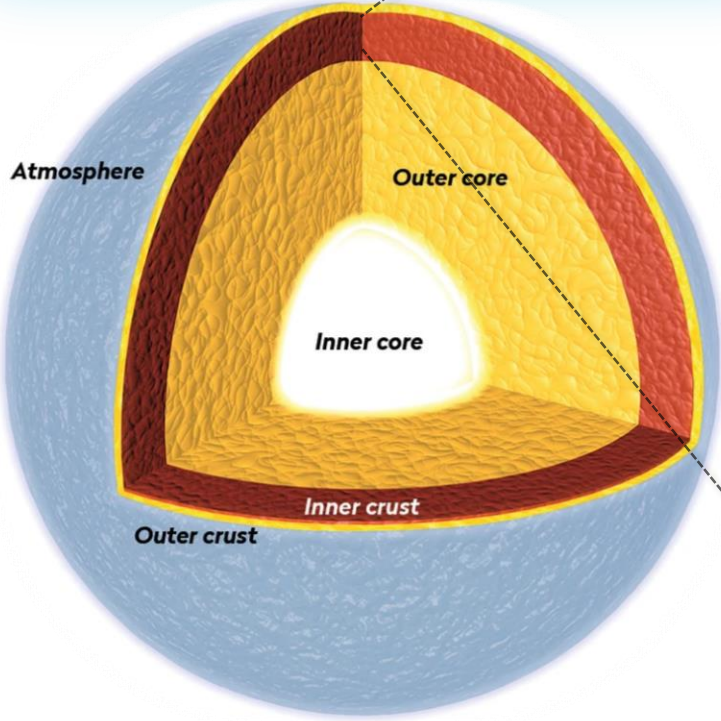


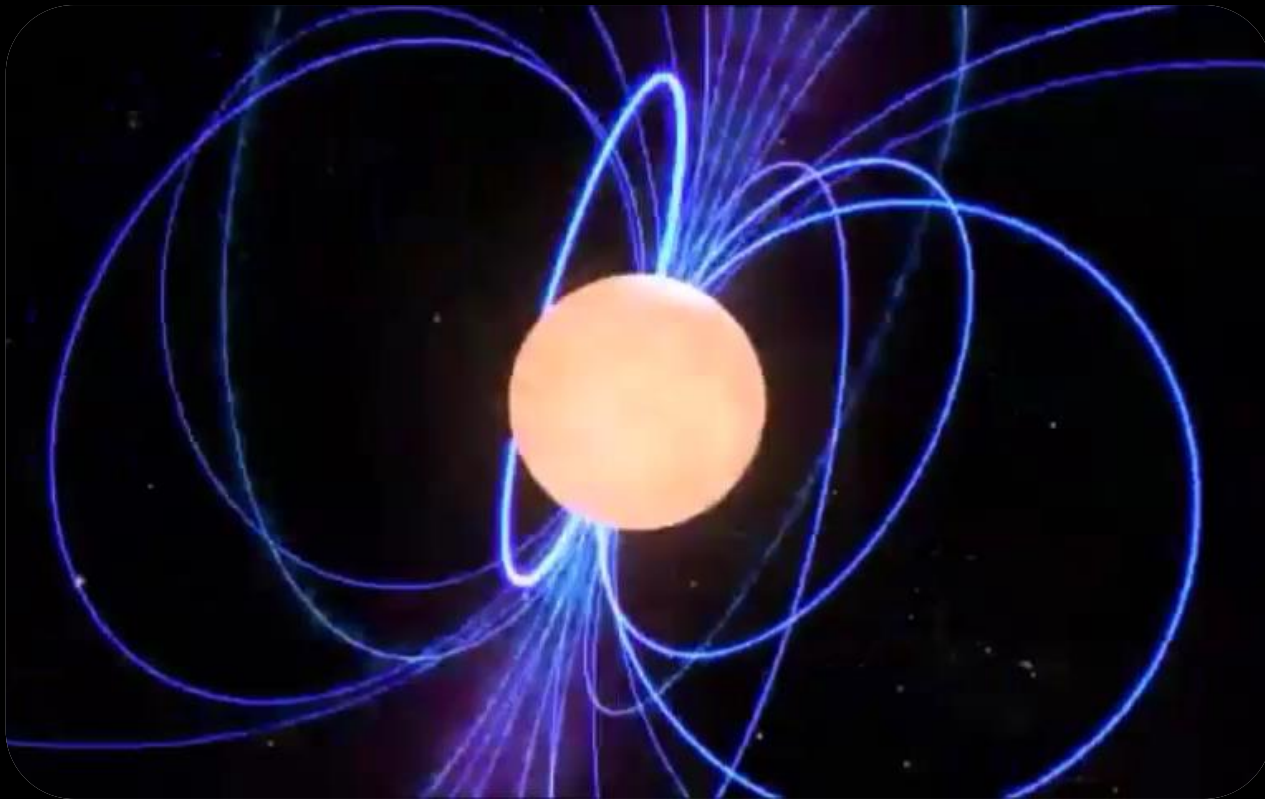
Figure: J.W. Negele and D. Vautherin, Nucl. Phys. **A207**, 298 (1978)
 Neutron star: <https://www.skyatnightmagazine.com/space-science/neutron-star/>

Neutron-star “glitch”



Pulsar - a rotating neutron star

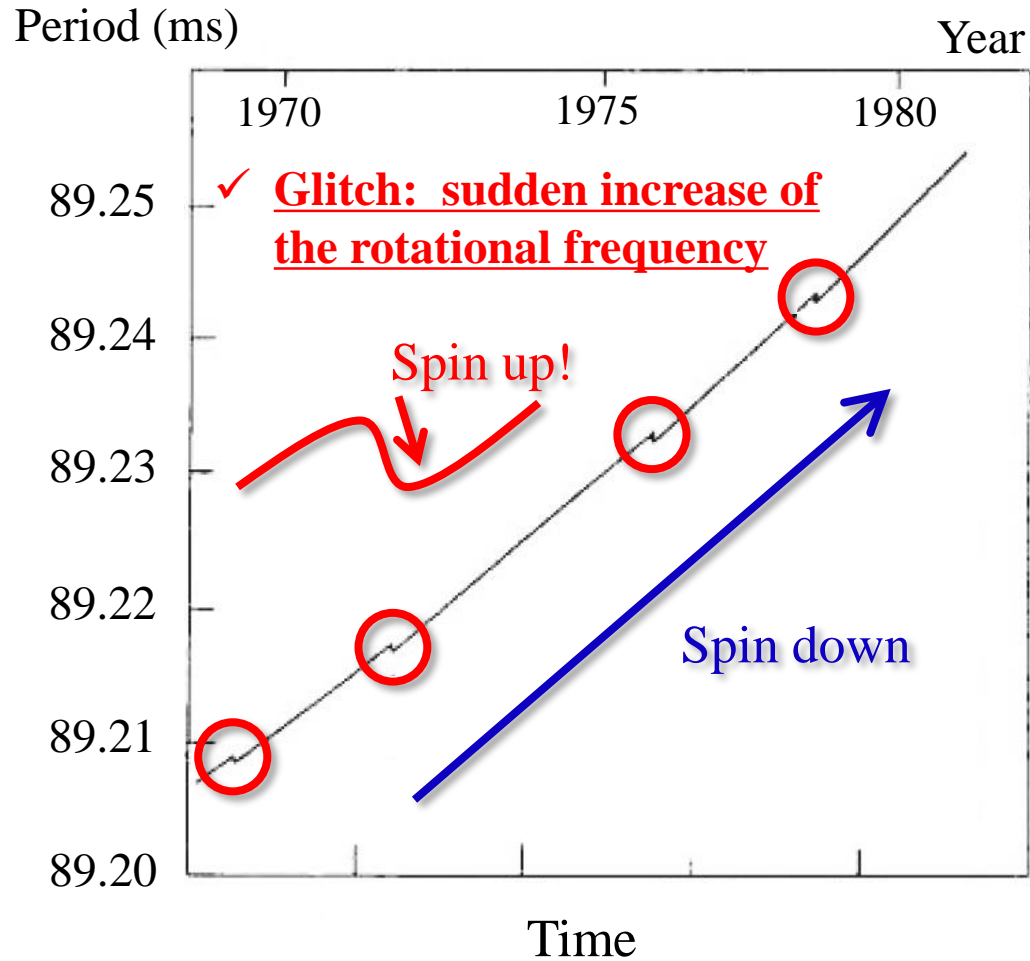
- ✓ First discovery in August 1967 → “Little Green Man” LGM-1 → PSR B1919+21
- ✓ Since then, more than 2650 pulsars have been observed
- ✓ It gradually spins down due to the EM radiation



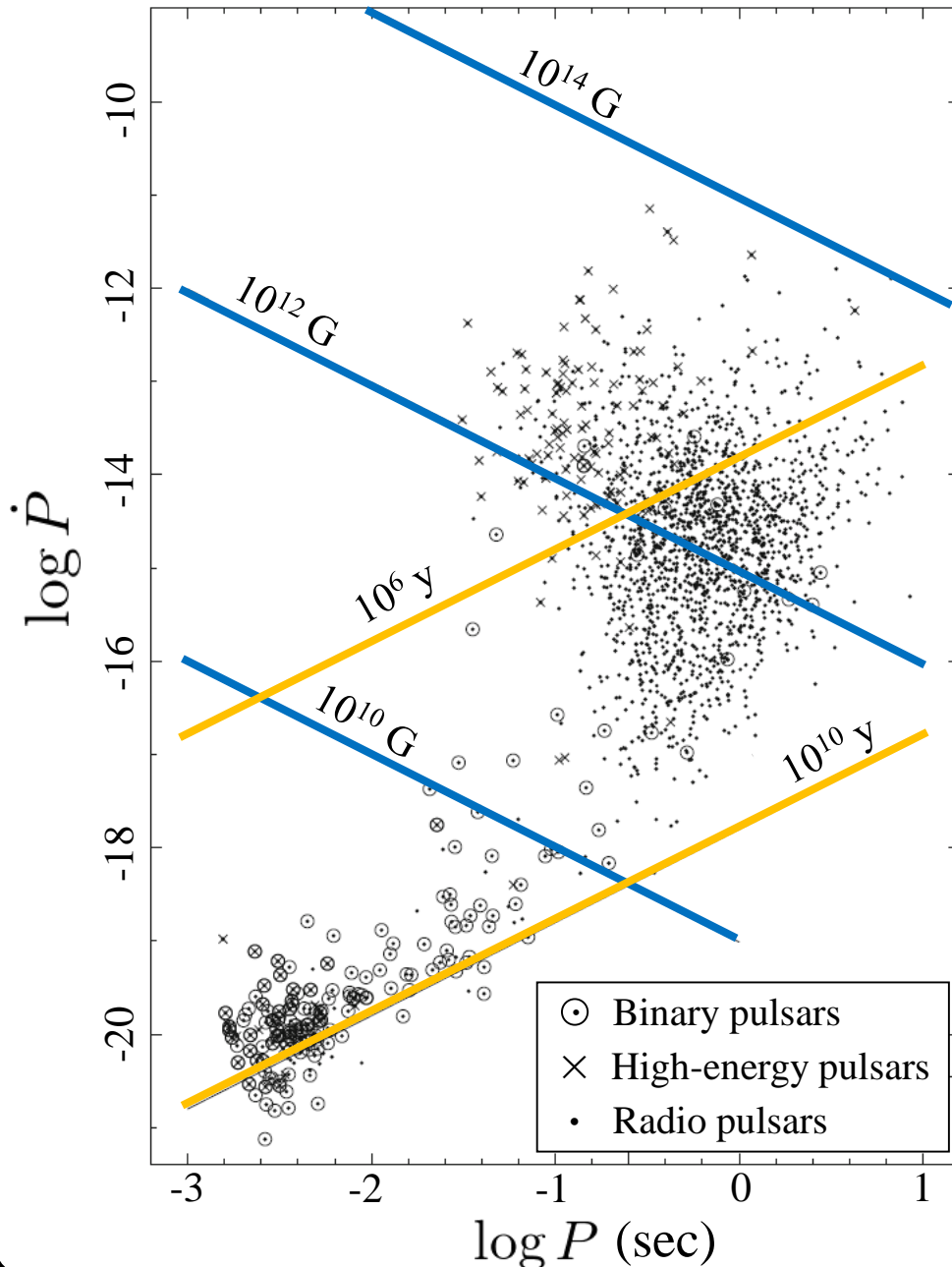
What is the glitch?

Typical example: the Vela pulsar

- *Irregularity* has been observed from continuous monitoring of the pulsation period



$P-\dot{P}$ diagram for pulsars (*not necessarily glitchers)



- ✓ Period (P): milliseconds to seconds
- ✓ Gradually spins down ($\dot{P} > 0$) due to the EM radiations
- ✓ Very stable “clock”, especially for millisecond pulsars, i.e. $\dot{P} \sim 10^{-20}$

➤ Characteristic age:

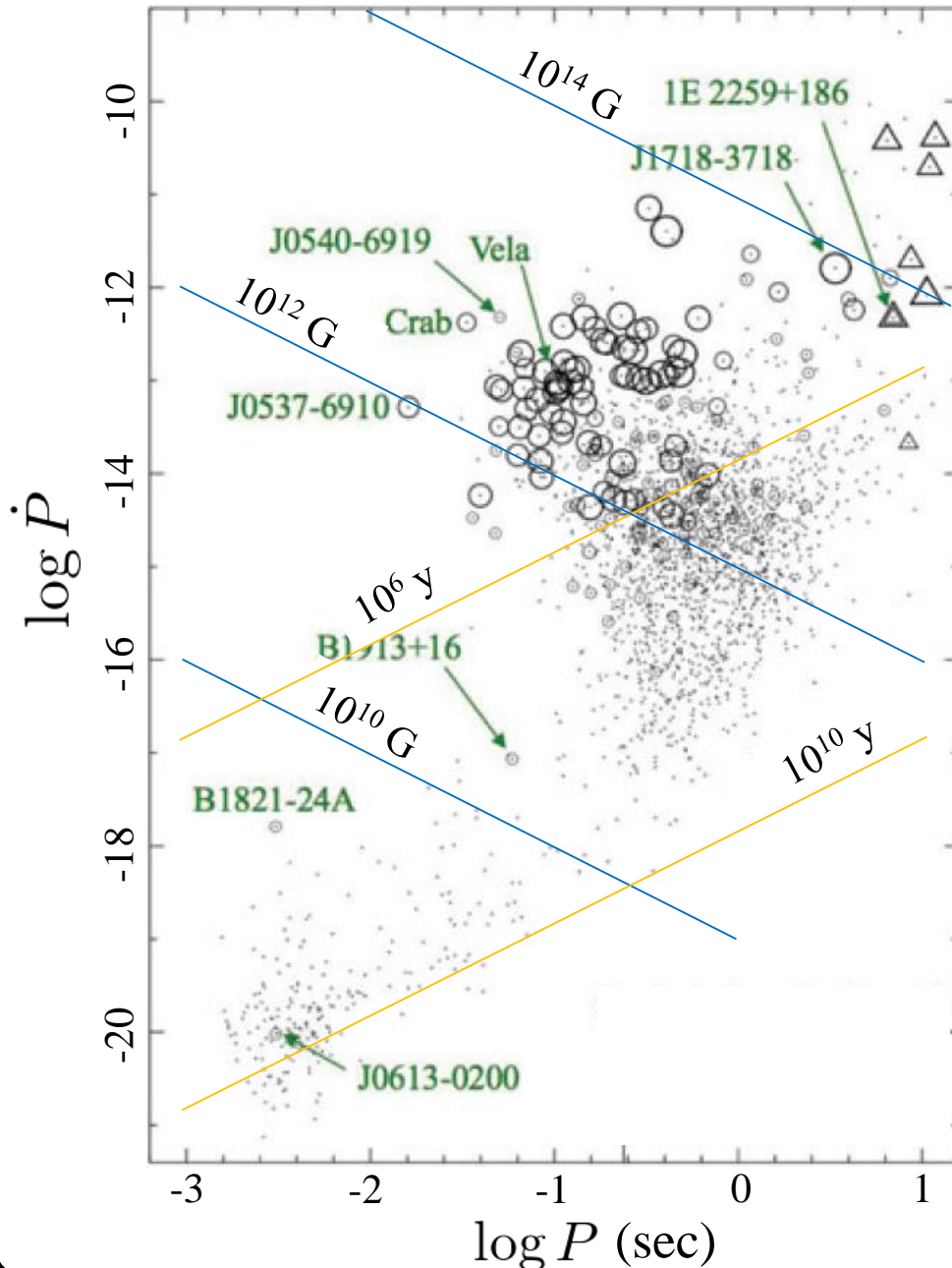
$$\tau_c = P/(2\dot{P})$$

➤ Surface dipole magnetic field strength:

$$B = 3.2 \times 10^{19} (P\dot{P}) \text{ G}$$

Figure taken from:

R.N. Manchester, *J. Astrophys. Astr.* **38**, 42 (2017)

$P-\dot{P}$ diagram for pulsars (*not necessarily glitchers)

- ✓ **More than 548 glitches** have been observed **in more than 180 pulsars**
- ✓ **Symbol size: glitch size** (typically, $\log(\Delta v/v) \sim 10^{-10}$ - 10^{-5})
- ✓ Young pulsars (including magnetars) exhibit larger glitches than older ones

➤ Characteristic age:

$$\tau_c = P/(2\dot{P})$$

➤ Surface dipole magnetic field strength:

$$B = 3.2 \times 10^{19} (P\dot{P}) \text{ G}$$

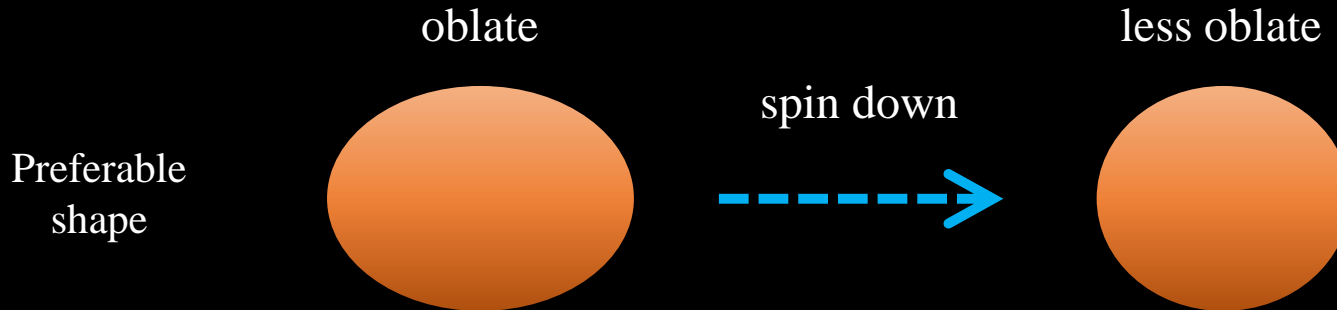
Figure taken from:

R.N. Manchester, Proc. IAU Symp. **337**, 197 (2017)

What happened?

Something must happen inside the neutron star!

- ✓ “Starquake” model [G. Baym et al., Nature **224**, 872 (1969)]



- ✓ Sudden change of moment of inertia → Glitch?

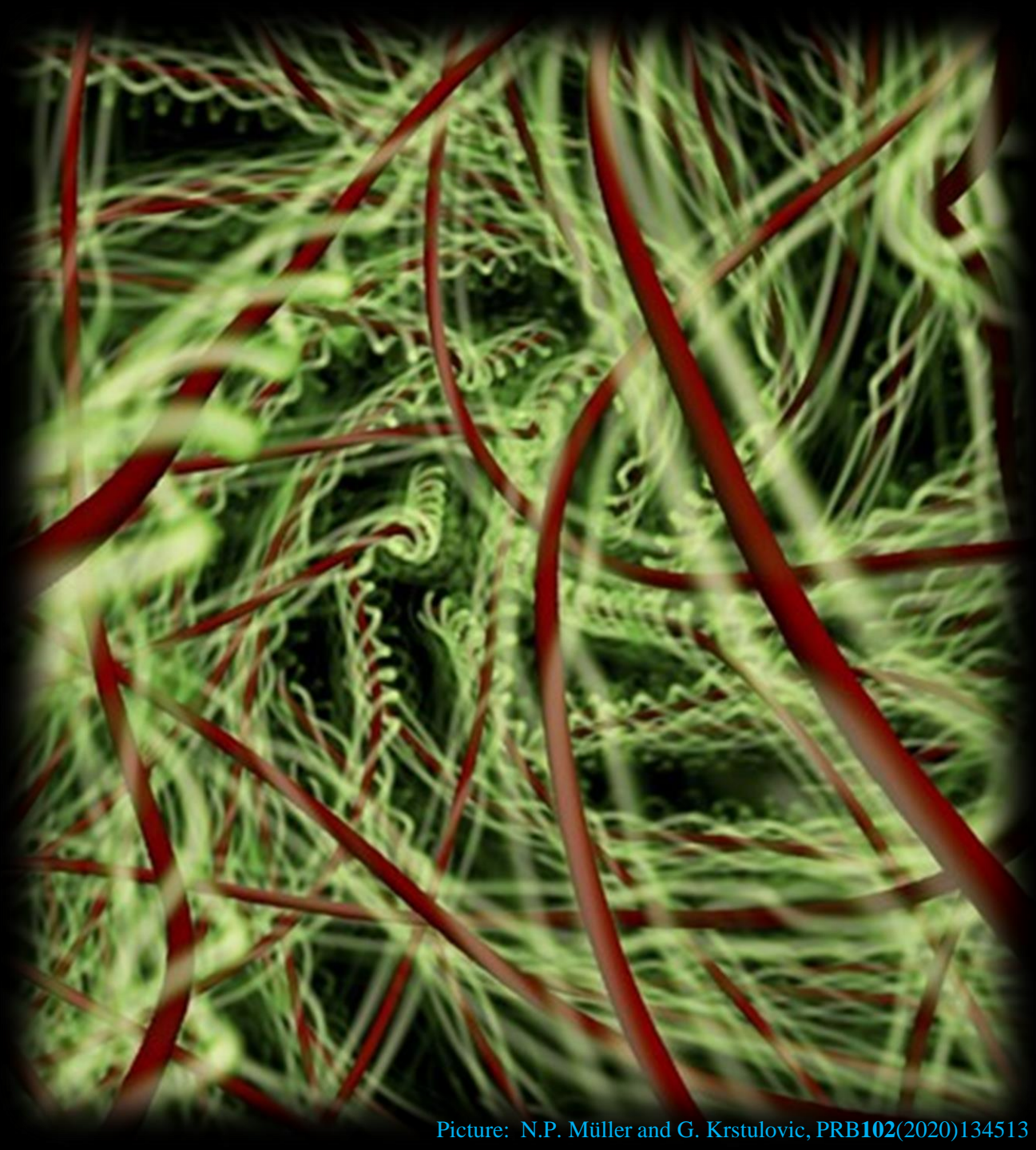
It requires hundreds of years to have a next glitch..

- ✓ Vortex mediated glitch [P.W. Anderson and N. Itoh, Nature **256**, 25 (1975)]

➤ Dynamics of superfluid “quantized vortices” play a key role!



Quantum vortices



In superfluid, vortices are quantized!

Superfluid order parameter:

$$\Delta(\mathbf{r}, t) = |\Delta(\mathbf{r}, t)|e^{i\phi(\mathbf{r}, t)}$$

Superfluid velocity:

$$\mathbf{v}_s(\mathbf{r}, t) = \frac{\hbar}{m}\nabla\phi(\mathbf{r}, t)$$



Vorticity:

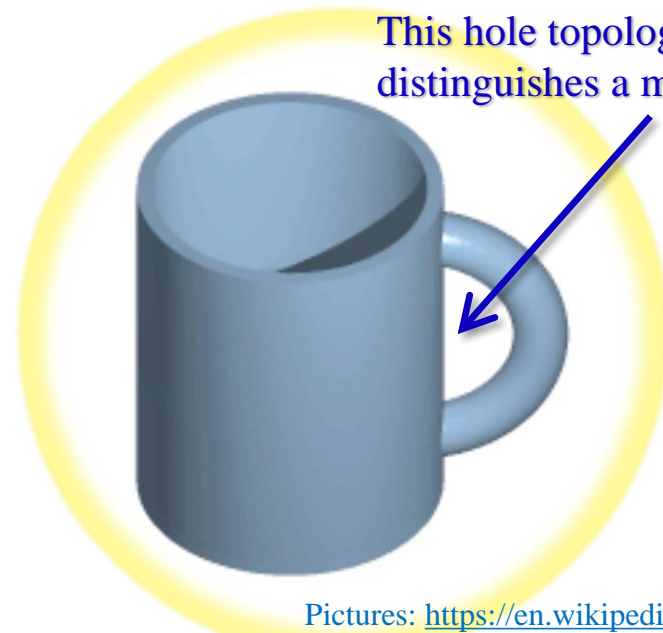
$$\underline{\omega = \nabla \times \mathbf{v}_s = 0}$$

superfluid is irrotational

Circulation:

$$\kappa = \int_S (\nabla \times \mathbf{v}_s) \cdot d\mathbf{S} = 0$$

*Unless, there is no topological defect



This hole topologically distinguishes a mug from a cow

Pictures: <https://en.wikipedia.org/wiki/Topology>

In superfluid, vortices are quantized!

Superfluid order parameter:

$$\Delta(\mathbf{r}, t) = |\Delta(\mathbf{r}, t)|e^{i\phi(\mathbf{r}, t)}$$

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superfluid is irrotational

Circulation:

$$\kappa = \int_S (\nabla \times \mathbf{v}_s) \cdot d\mathbf{S} = 0$$

*Unless, there is no topological defect

If there is a defect:



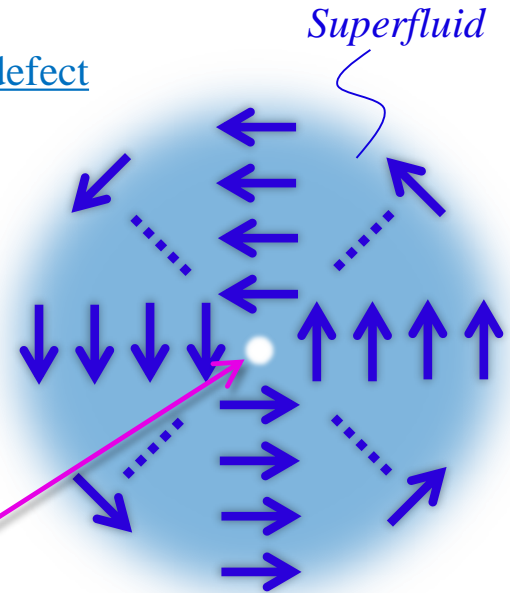
$$\kappa = \oint_C \mathbf{v}_s \cdot d\mathbf{l} = \frac{\hbar}{m} \oint_C \nabla \phi \cdot d\mathbf{l} = \frac{\hbar}{m} (2\pi n)$$

Quantization of circulation

*the phase $\phi(\mathbf{r})$ at the same point must be equivalent!

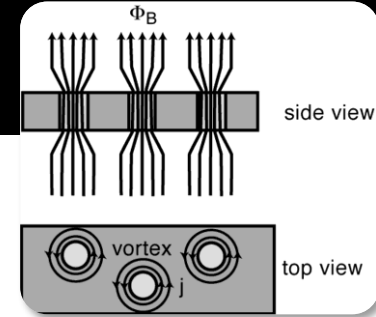
Flow velocity of rotation shall be quantized!

A hole at which superfluidity is lost



Quantum vortex

What is a quantum vortex?



In superconductor, magnetic flux is quantized!

Magnetic flux:

$$\Phi = \int_S \mathbf{B} \cdot d\mathbf{S} = \int_S (\nabla \times \mathbf{A}) \cdot d\mathbf{S} = 0$$

Meissner effect

$$\mathbf{j}_s = -\frac{n_s e_s^2}{m_s} \mathbf{A} + \frac{n_s e_s \hbar}{m_s} \nabla \phi : \text{the London equation}$$

If there is a defect:

$$\Phi = \oint_C \mathbf{A} \cdot d\mathbf{l} \approx \frac{\hbar}{e_s} \oint_C \nabla \phi \cdot d\mathbf{l} = \frac{\hbar}{e_s} (2\pi n)$$

n_s, m_s, e_s : density, mass, and charge of a carrier (Cooper pair)

Quantization of magnetic flux (fluxtube, fluxoid, or fluxon)

Circulation:

$$\kappa = \int_S (\nabla \times \mathbf{v}_s) \cdot d\mathbf{S} = 0$$

*Unless, there is no topological defect

If there is a defect:

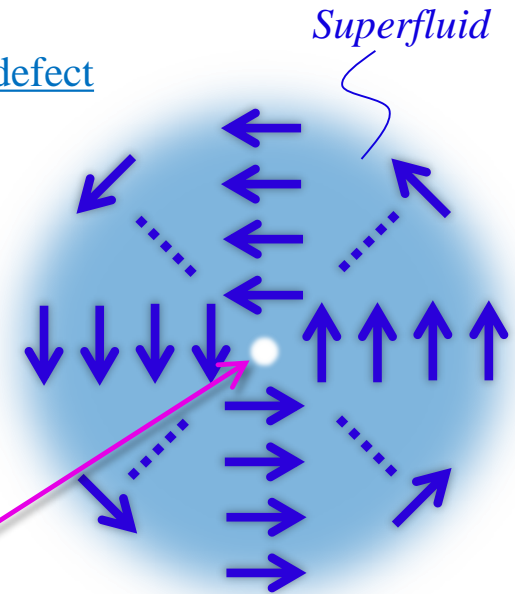
$$\kappa = \oint_C \mathbf{v}_s \cdot d\mathbf{l} = \frac{\hbar}{m} \oint_C \nabla \phi \cdot d\mathbf{l} = \frac{\hbar}{m} (2\pi n)$$

Quantization of circulation

*the phase $\phi(\mathbf{r})$ at the same point must be equivalent!

Flow velocity of rotation shall be quantized!

A hole at which superfluidity is lost

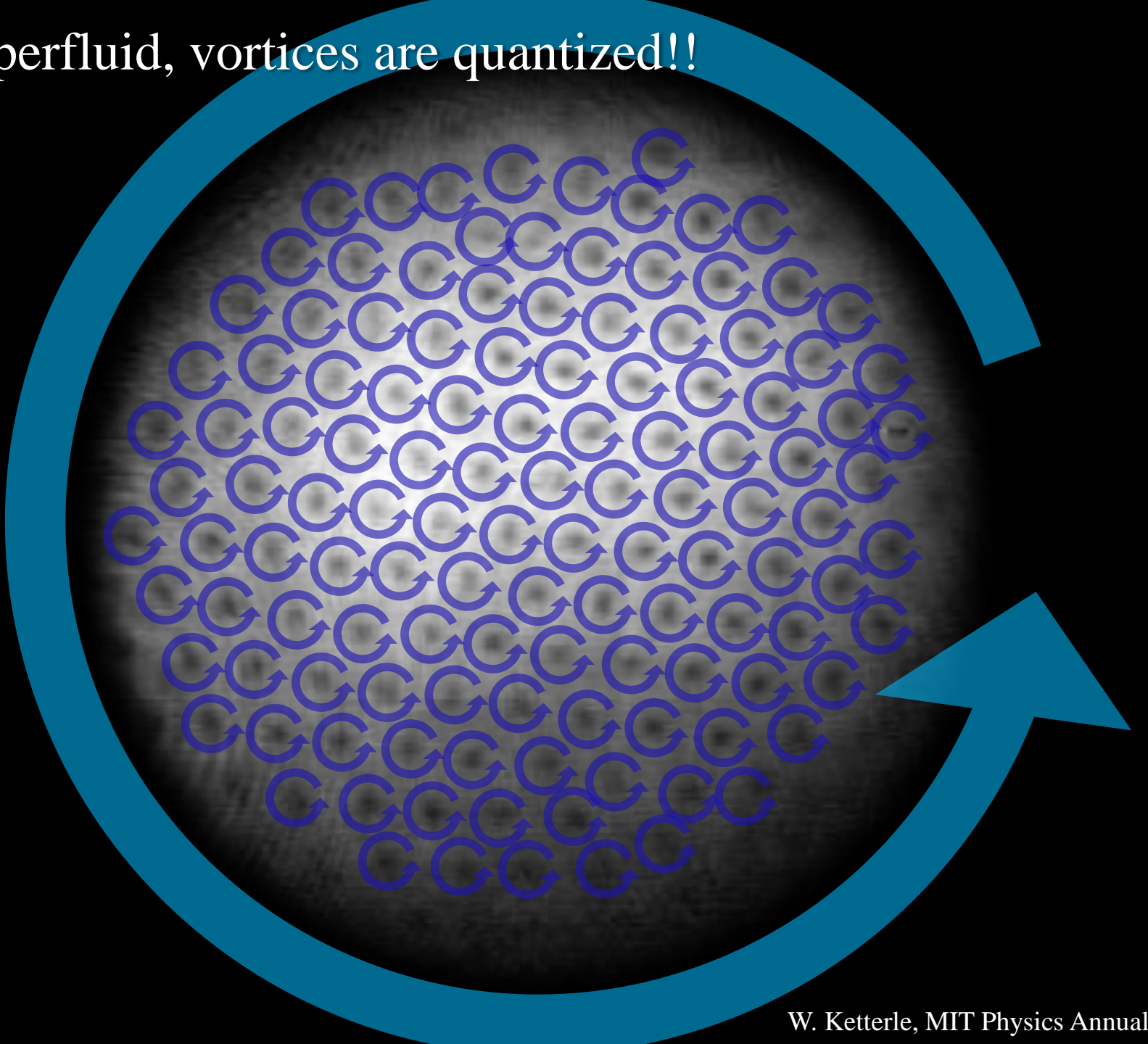


Quantum vortex

In daily life, a vortex is continuous..



In superfluid, vortices are quantized!!



A movie from a talk by W. Guo (available from <https://youtu.be/P2ckefSAN20>) at
INT Program 19-1a “Quantum Turbulence: Cold Atoms, Heavy Ions, and Neutron Stars”
March 18 - April 19, 2019

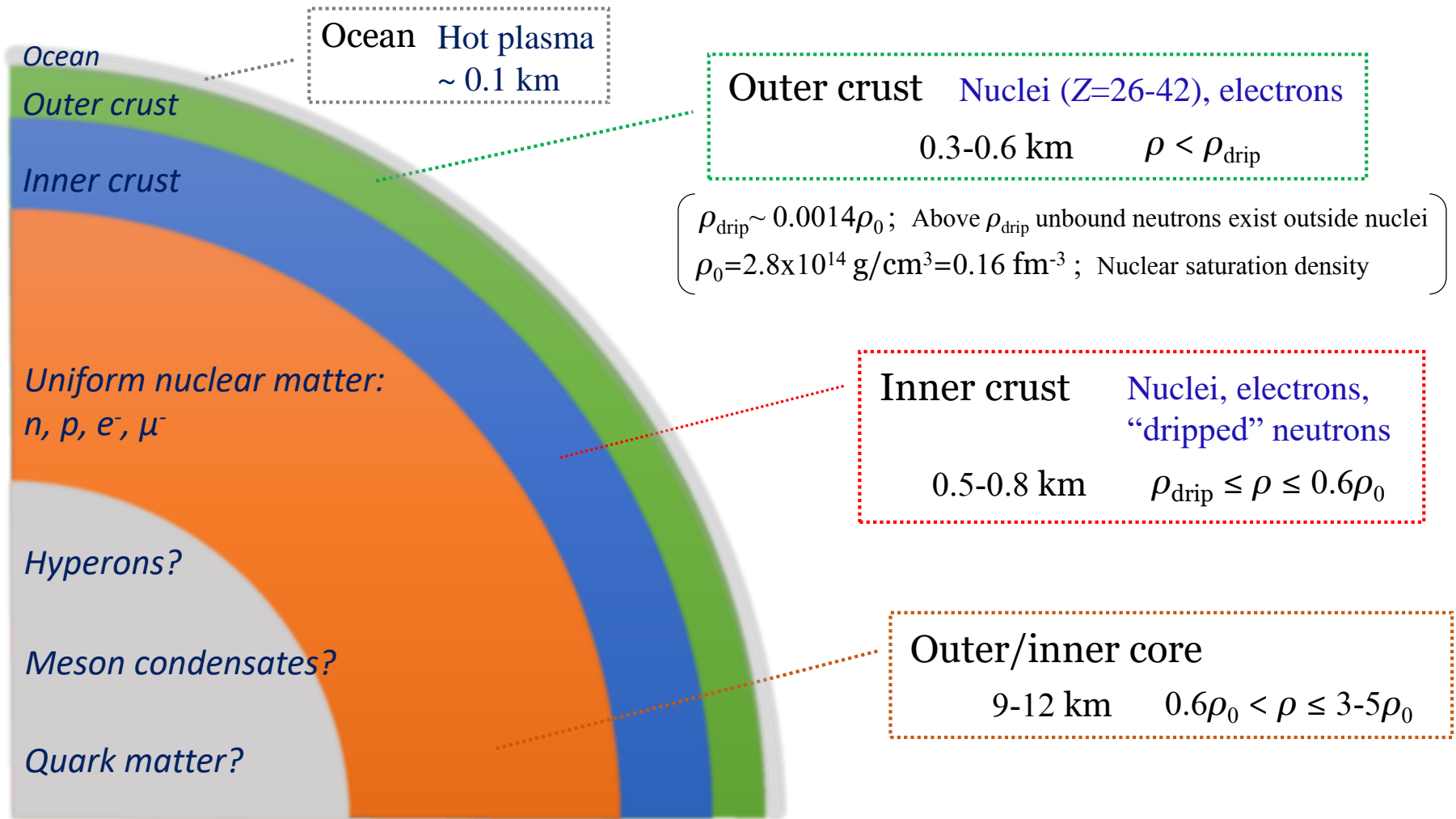
Direct visualization of quantized vortices



Hydrogen particles were trapped in the vortex core, then worked as a tracer

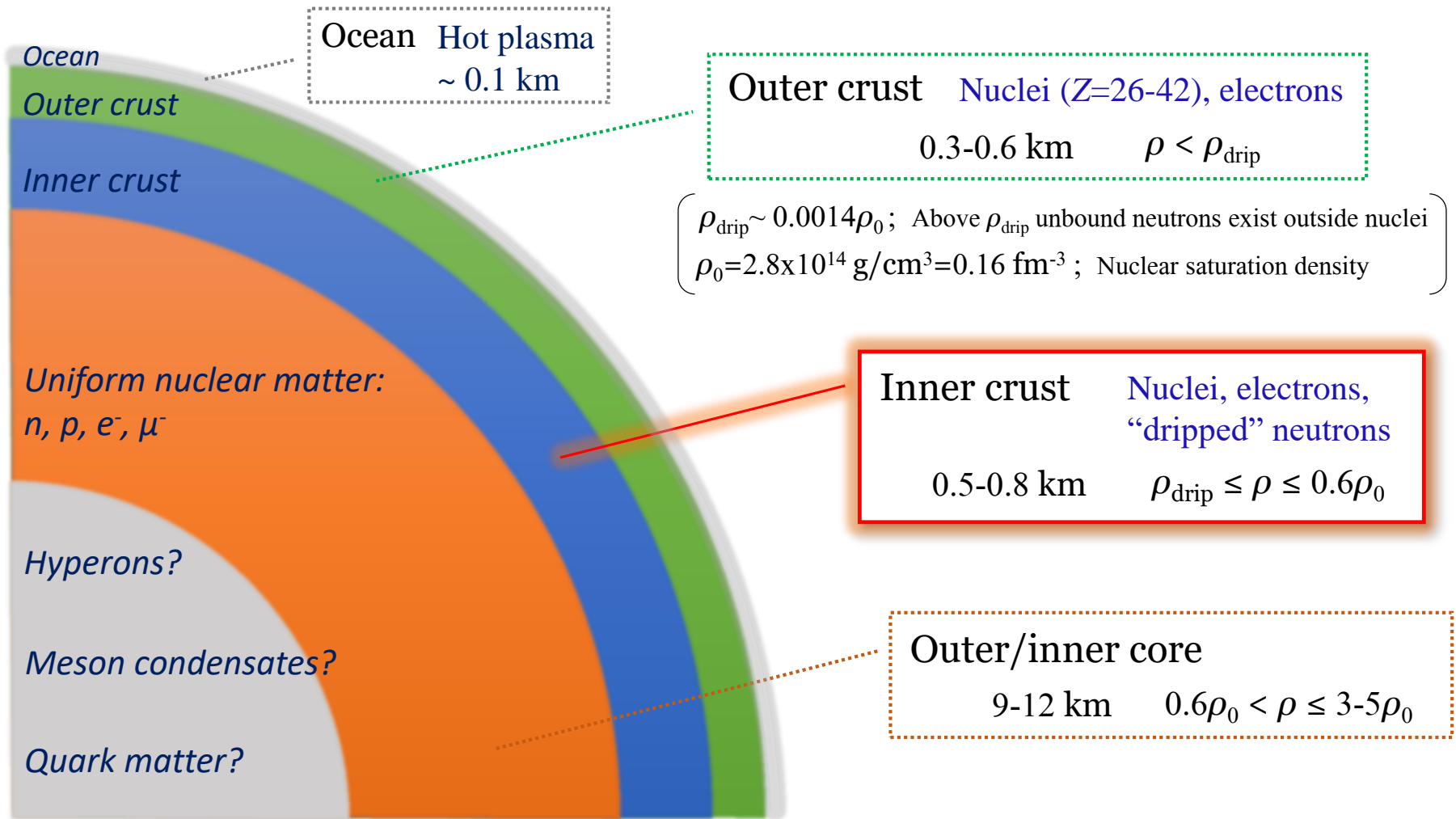
Neutron star is a great playground for nuclear physicists

- ✓ It offers extreme situations which can not be realized in terrestrial experiments



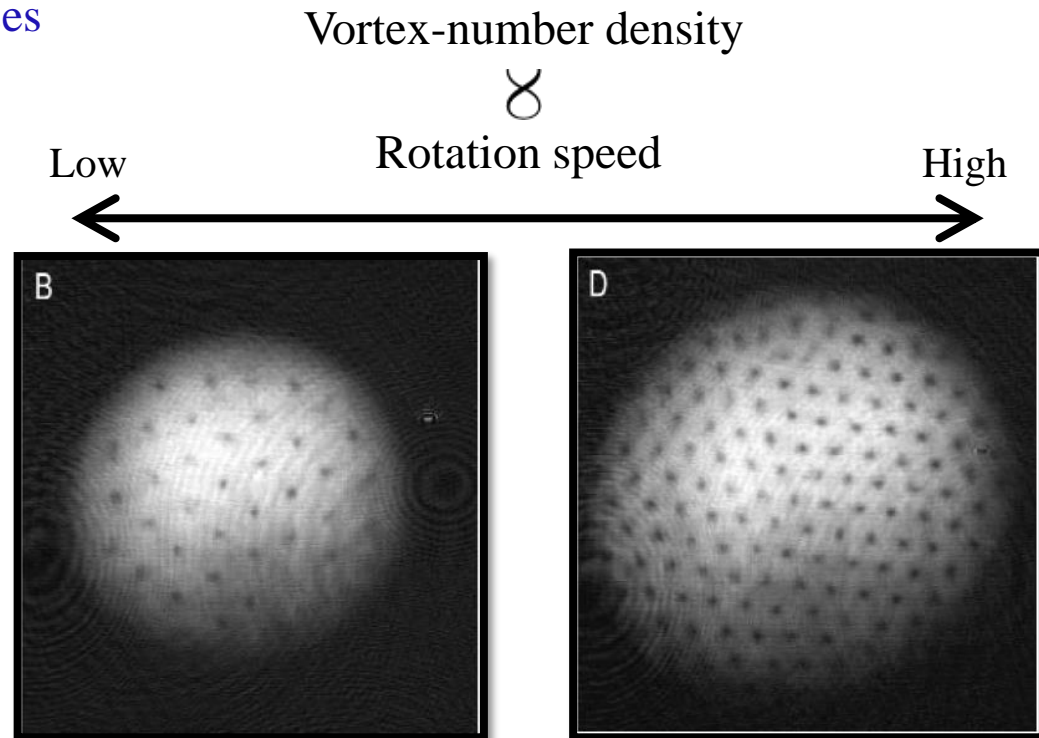
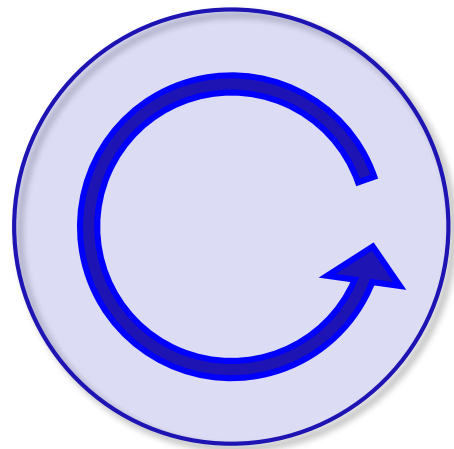
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- ✓ It offers extreme situations which can not be realized in terrestrial experiments



In rotating superfluid, an array of quantum vortices is generated

□ Observation in ultra-cold atomic gases

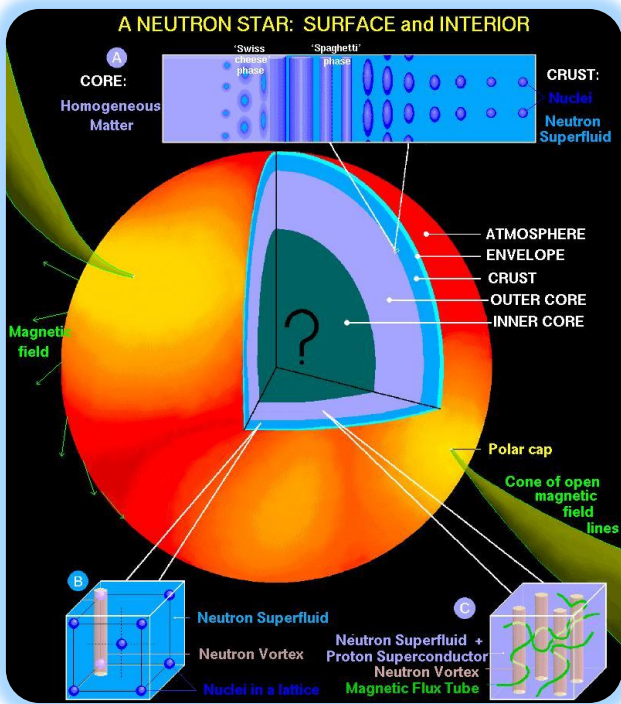


W. Ketterle, MIT Physics Annual. 2001

Quantum vortices in a neutron star

In rotating superfluid, an array of quantum vortices is generated

Observation in ultra-cold atomic gases



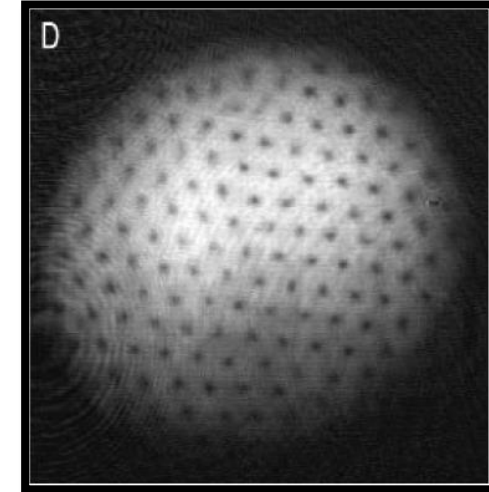
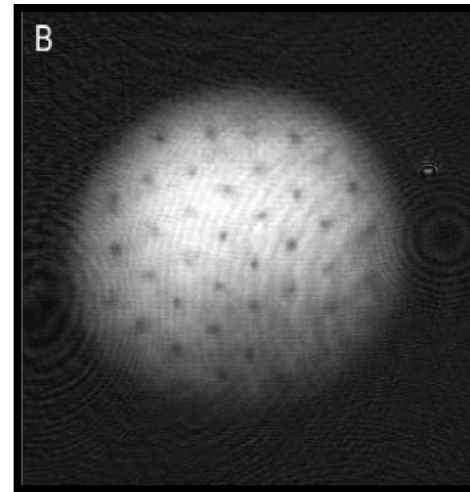
Vortex-number density

\propto

Rotation speed

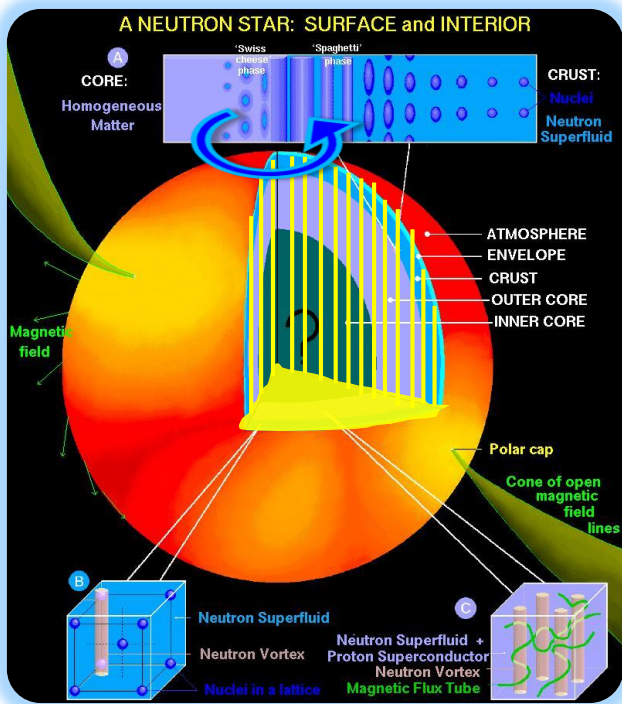
Low

High



W. Ketterle, MIT Physics Annual. 2001

There must be a huge number ($\sim 10^{18}$) of vortices inside a neutron star!!



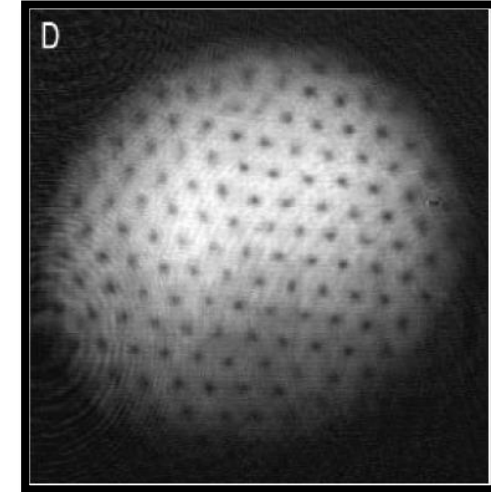
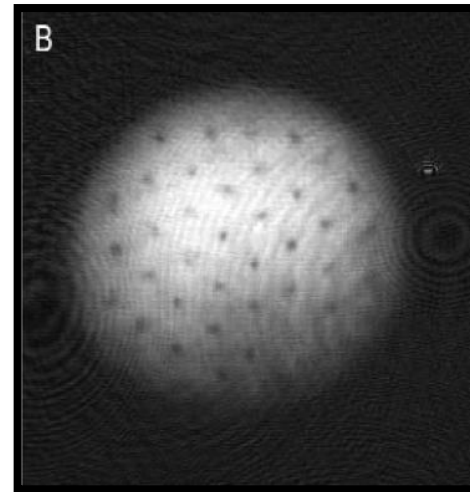
Vortex-number density

\propto

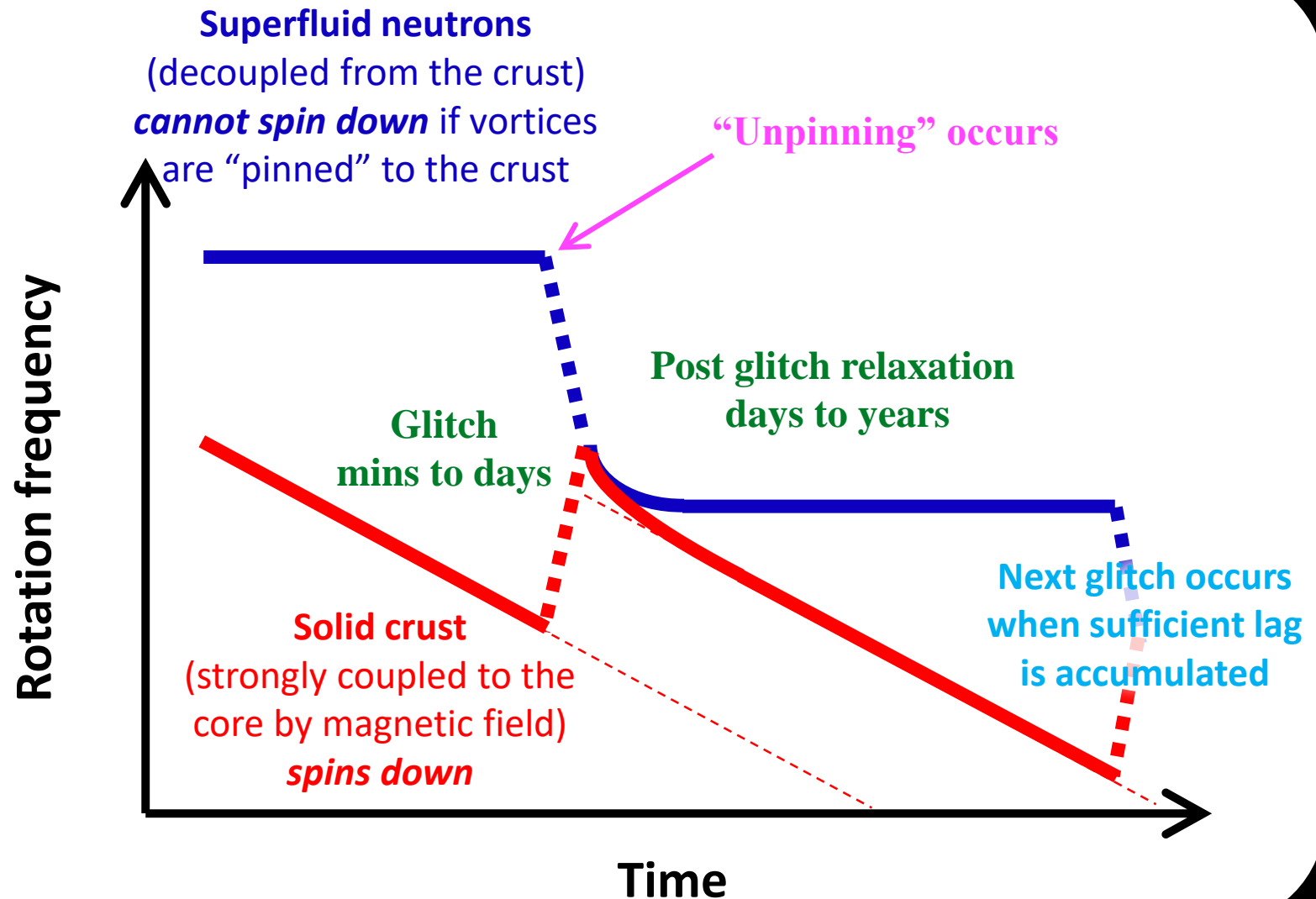
Rotation speed

Low

High



The vortex mediated glitch: Naive picture



To fully understand the glitches, we need to clarify:

Glitch dynamics

and, of course,
details of NS matter..

How do vortices move?

Pinning mechanism

How are vortices pinned?

Trigger mechanism

How are vortices unpinned?



We attacked this problem using
the state-of-the-art microscopic nuclear theory

We attack this problem with HPC on GPU supercomputers
with TDDFT for superfluid systems, TDSLDA!

TDSLDA: TDDFT with local treatment of pairing

Kohn-Sham scheme is extended for non-interacting quasiparticles

➤ TDSLDA equations (formally equivalent to TDHFB or TD-BdG equations)

$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} u_{k,\uparrow}(\mathbf{r}, t) \\ u_{k,\downarrow}(\mathbf{r}, t) \\ v_{k,\uparrow}(\mathbf{r}, t) \\ v_{k,\downarrow}(\mathbf{r}, t) \end{pmatrix} = \begin{pmatrix} h_{\uparrow\uparrow}(\mathbf{r}, t) & h_{\uparrow\downarrow}(\mathbf{r}, t) & 0 & \Delta(\mathbf{r}, t) \\ h_{\downarrow\uparrow}(\mathbf{r}, t) & h_{\downarrow\downarrow}(\mathbf{r}, t) & -\Delta(\mathbf{r}, t) & 0 \\ 0 & -\Delta^*(\mathbf{r}, t) & -h_{\uparrow\uparrow}^*(\mathbf{r}, t) & -h_{\uparrow\downarrow}^*(\mathbf{r}, t) \\ \Delta^*(\mathbf{r}, t) & 0 & -h_{\downarrow\uparrow}^*(\mathbf{r}, t) & -h_{\downarrow\downarrow}^*(\mathbf{r}, t) \end{pmatrix} \begin{pmatrix} u_{k,\uparrow}(\mathbf{r}, t) \\ u_{k,\downarrow}(\mathbf{r}, t) \\ v_{k,\uparrow}(\mathbf{r}, t) \\ v_{k,\downarrow}(\mathbf{r}, t) \end{pmatrix}$$

$$h_{\sigma} = \frac{\delta E}{\delta n_{\sigma}} \quad : \text{ s.p. Hamiltonian}$$

$$\Delta = -\frac{\delta E}{\delta \nu^*} \quad : \text{ pairing field}$$

$$n_{\sigma}(\mathbf{r}, t) = \sum_{E_k < E_c} |v_{k,\sigma}(\mathbf{r}, t)|^2 \quad : \text{ number density}$$

$$\nu(\mathbf{r}, t) = \sum_{E_k < E_c} u_{k,\uparrow}(\mathbf{r}, t) v_{k,\downarrow}^*(\mathbf{r}, t) \quad : \text{ anomalous density}$$

$$\mathbf{j}_{\sigma}(\mathbf{r}, t) = \hbar \sum_{E_k < E_c} \text{Im}[v_{k,\sigma}^*(\mathbf{r}, t) \nabla v_{k,\sigma}(\mathbf{r}, t)] \quad : \text{ current}$$

A large number (10⁴-10⁶) of 3D coupled non-linear PDEs have to be solved!!

of qp orbitals ~ # of grid points

TDSLDA: TDDFT with local treatment of pairing

Kohn-Sham scheme is extended for non-interacting quasiparticles

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Supercomputing!!

$$h_{\sigma} = \frac{\delta E}{\delta n_{\sigma}} : \text{s.p. Hamiltonian}$$

$$n_{\sigma}(\mathbf{r}, t) = \sum_{E_k < E_c} |v_{k,\sigma}(\mathbf{r}, t)|^2 : \text{number density}$$

$$\nu(\mathbf{r}, t) = \sum_{E_k < E_c} u_{k,\uparrow}(\mathbf{r}, t) v_{k,\downarrow}^*(\mathbf{r}, t) : \text{anomalous density}$$

$$\Delta = -\frac{\delta E}{\delta \nu^*} : \text{pairing field}$$

$$\mathbf{j}_{\sigma}(\mathbf{r}, t) = \hbar \sum_{E_k < E_c} \text{Im}[v_{k,\sigma}^*(\mathbf{r}, t) \nabla v_{k,\sigma}(\mathbf{r}, t)] : \text{current}$$

A large number (10^4 - 10^6) of 3D coupled non-linear PDEs have to be solved!!

of qp orbitals ~ # of grid points

*The number indicates the rank according to the [TOP500 list \(June 2022\)](#)

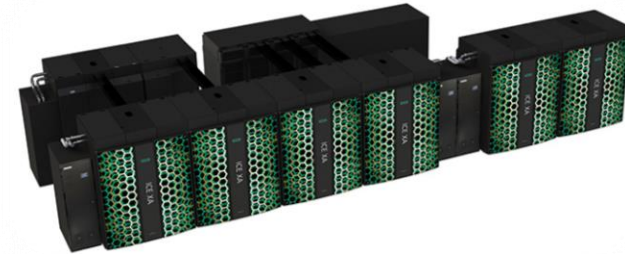
Piz Daint, CSCS, Switzerland (No. 23)



TITAN, ORNL, USA



TSUBAME3.0, Japan (No. 64)



Summit, ORNL, USA (No. 4)
GPU, 200 PFlops/s

No.4: Summit, ORNL, USA

No.5: Sierra, LLNL, USA

No.7: Perlmutter, NERSC, USA

No.8: Selene, NVIDIA Co., USA

No.11: JUWELS Booster Module, FZJ, Germany

No.12: HPC5, Eni S. p. A., Italy

No.13: Voyager-EUS2, Azure East US 2, USA

No.14: Polaris, ANL, USA

No.15: SSC-21, Samsung Electronics, South Korea

No.18: Damman-7, Saudi Aramco, Saudi Arabia

No.19: ABCI 2.0, AIST, Japan

**GPU machines
within Nos. 1-20**

**Certainly, GPU is competing
with CPU machines!!**

TDSLDA has been successfully applied for both UFG and nuclear systems

Unitary Fermi Gas (UFG)

$$\left(k_{\text{F}} a \rightarrow \infty, k_{\text{F}} r_{\text{eff}} \rightarrow 0 \quad \begin{array}{l} a : \text{s-wave scattering length} \\ r_{\text{eff}} : \text{effective range} \end{array} \right)$$

A. Bulgac and S. Yoon, PRL**102**(2009)085302.

A. Bulgac *et al.*, Science **332**(2011)1288.

A. Bulgac *et al.*, PRL**108**(2012)150401.

A. Bulgac *et al.*, PRL**112**(2014)025301.

G. Wlazłowski *et al.*, PRA**91**(2015)031602(R).

G. Wlazłowski *et al.*, PRL**120**(2018)253002.

P. Magierski *et al.*, PRA**100**(2019)033613.

Large-amplitude pairing field dynamics

Dynamics of quantum vortices

Quantum shock waves and domain walls

Dynamics of vortex rings

Dynamics of vortices and quantum turbulence

Solitonic cascades in spin-polarized UFG

Spin-polarized droplet in UFG

Nuclear systems

I. Stetcu *et al.*, PRC**84**(2011)051309(R).

I. Stetcu *et al.*, PRL**114**(2015)012701.

A. Bulgac *et al.*, PRL**116**(2016)122504;

PRC**100**(2019)034615.

G. Wlazłowski *et al.*, PRL**117**(2016)232701.

Isovector giant dipole resonance (IVGDR)

Relativistic coulomb excitation

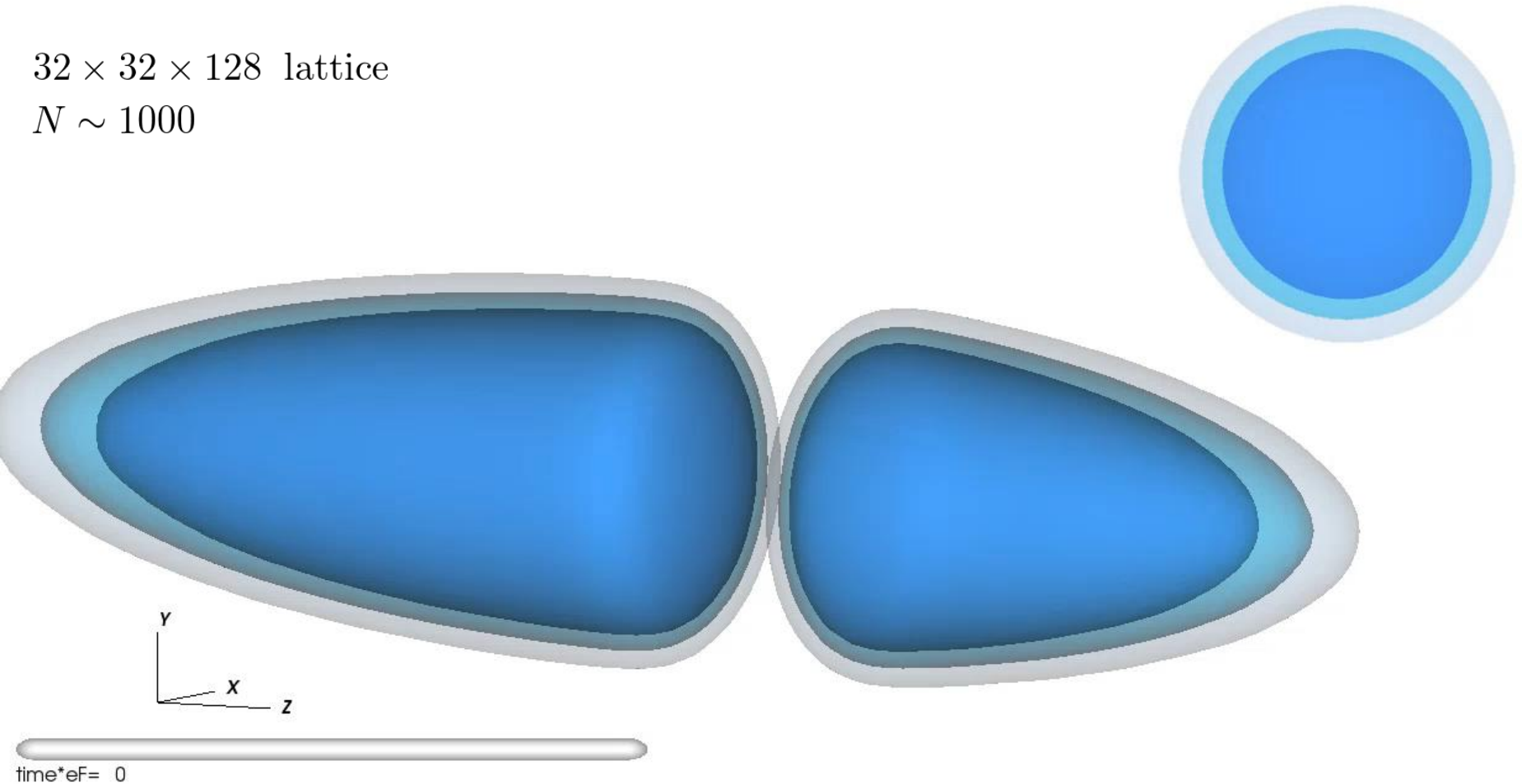
Induced fission of ^{240}Pu

Vortex-nucleus interaction

Result of TDSLDA simulation:

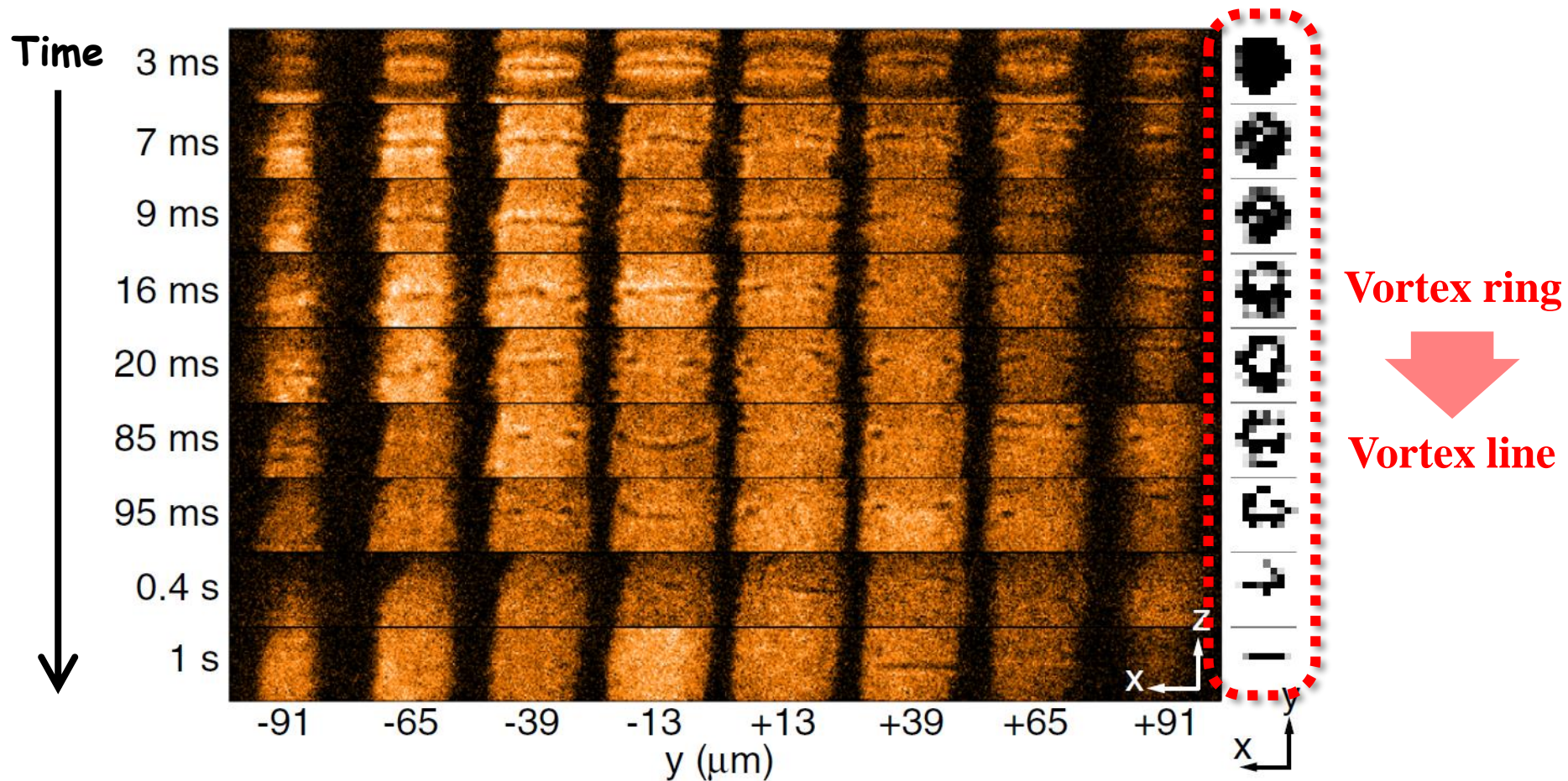
Phase discontinuity creates a vortex ring which decays into a vortex line

$32 \times 32 \times 128$ lattice
 $N \sim 1000$



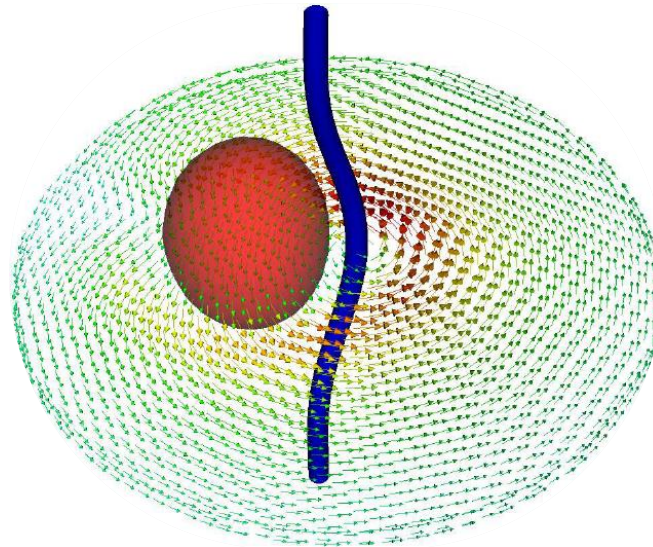
G. Wlazłowski, A. Bulgac, M.M. Forbes, and K.J. Roche, Phys. Rev. A **91**, 031602(R) (2015)

The cascades of solitonic excitations have been identified experimentally



M.J.H. Ku, B. Mukherjee, T. Yefsah, and M.W. Zwierlein, Phys. Rev. Lett. **116**, 045304 (2016)

Vortex-nucleus dynamics within TDSLDA



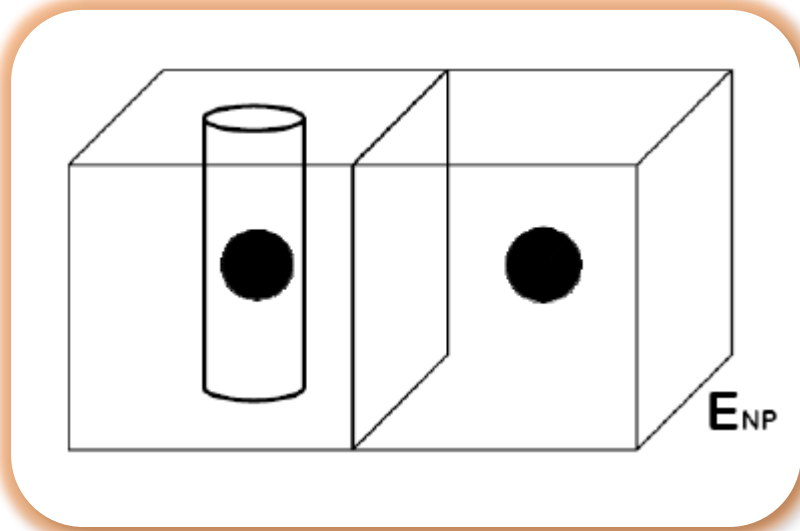
A key to understand the glitches is:
Vortex pinning mechanism in the inner crust of neutron stars

Q. Is the vortex-nucleus interaction

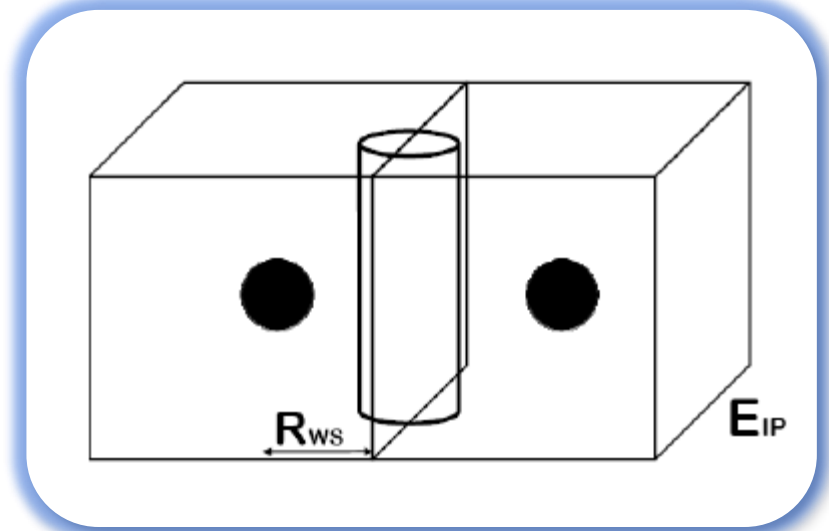
Attractive?

or

Repulsive?



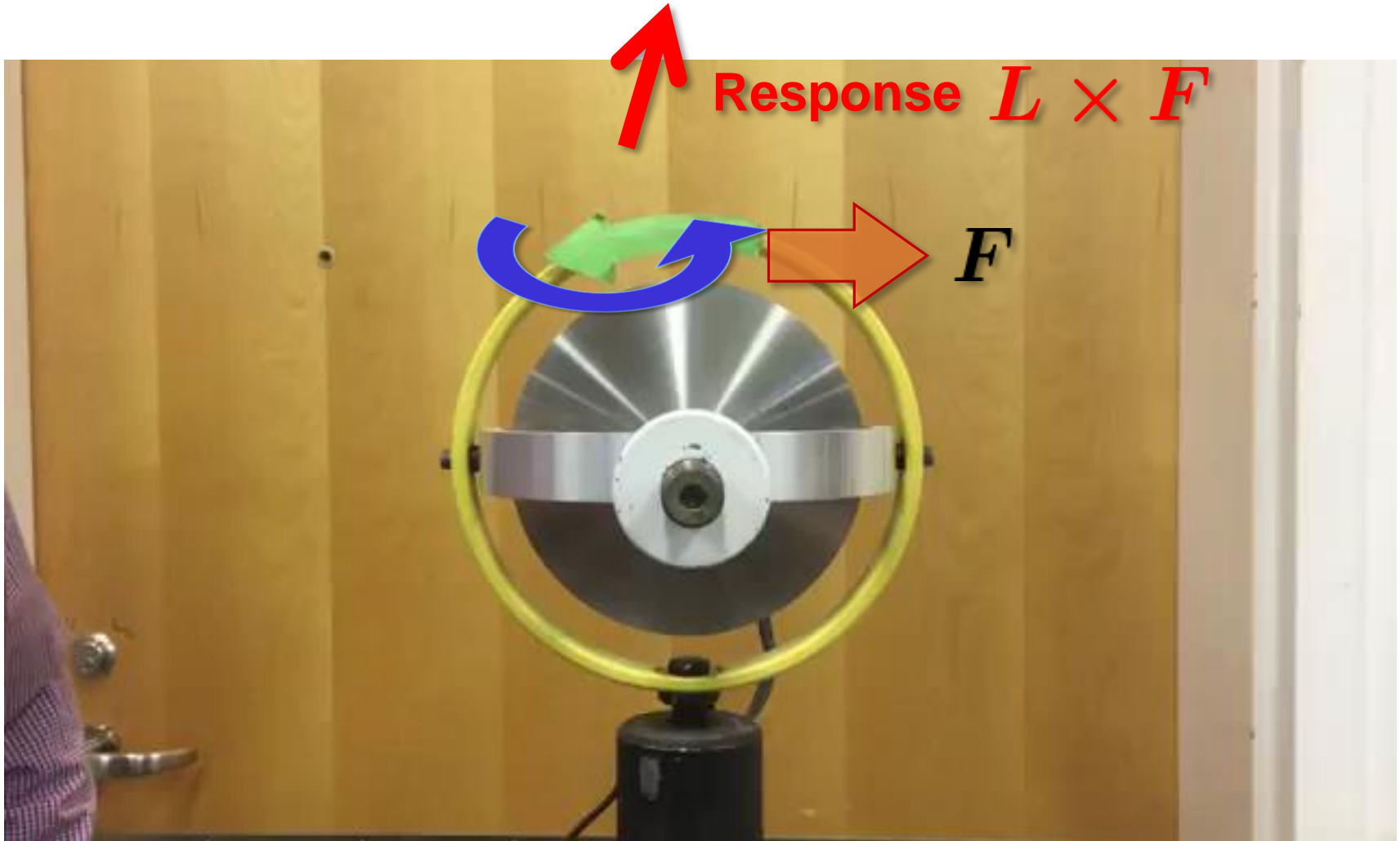
“Nuclear pinning”



“Interstitial pinning”

What we investigated - Vortex-nucleus dynamics

Response of a spinning gyroscope when pushed



We performed 3D, dynamical simulations by TDDFT with superfluidity

□ TDSLDA equations (or TDHFB, TD-BdG)

$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} u_i(\mathbf{r}) \\ v_i(\mathbf{r}) \end{pmatrix} = \begin{pmatrix} h(\mathbf{r}) & \Delta(\mathbf{r}) \\ \Delta^*(\mathbf{r}) & -h(\mathbf{r}) \end{pmatrix} \begin{pmatrix} u_i(\mathbf{r}) \\ v_i(\mathbf{r}) \end{pmatrix}$$

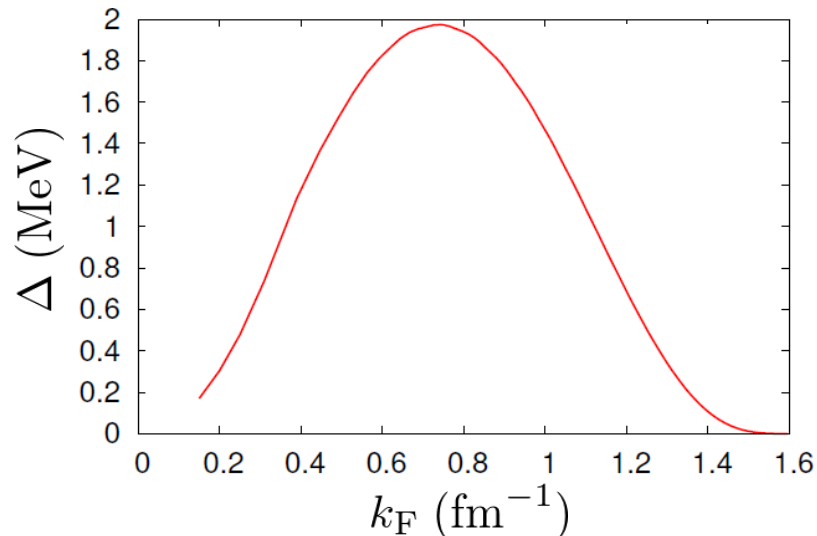
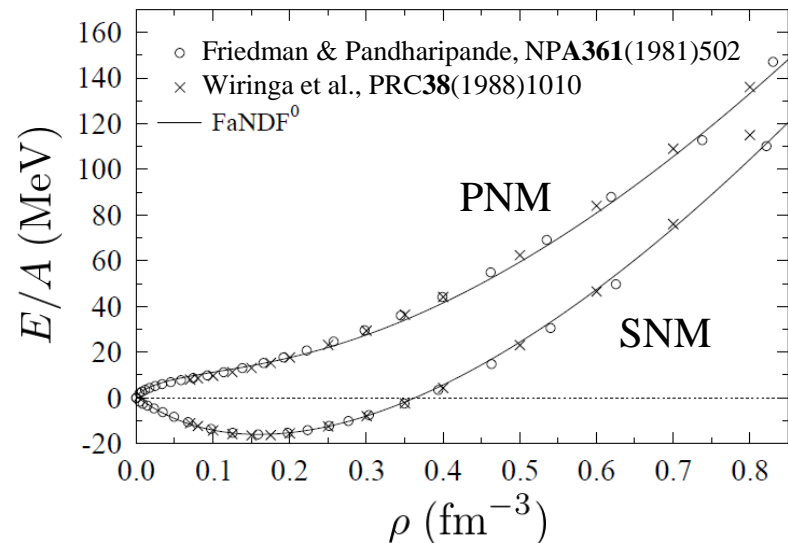
□ Energy density functional (EDF)

$$\mathcal{E}(\mathbf{r}) = \mathcal{E}_0(\mathbf{r}) + \mathcal{E}_{\text{pair}}(\mathbf{r})$$

$\mathcal{E}_0(\mathbf{r})$: Fayans EDF (FaNDF⁰) w/o LS

$$\mathcal{E}(\mathbf{r}) = \sum_{q=n,p} g[\rho_q(\mathbf{r})] |\nu_q(\mathbf{r})|^2$$

S.A. Fayans, JETP Lett. **68**, 169 (1998)



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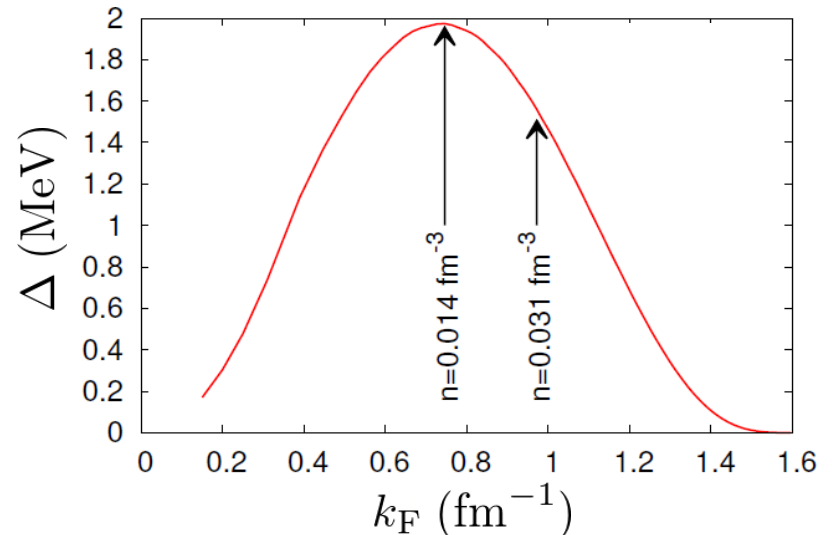
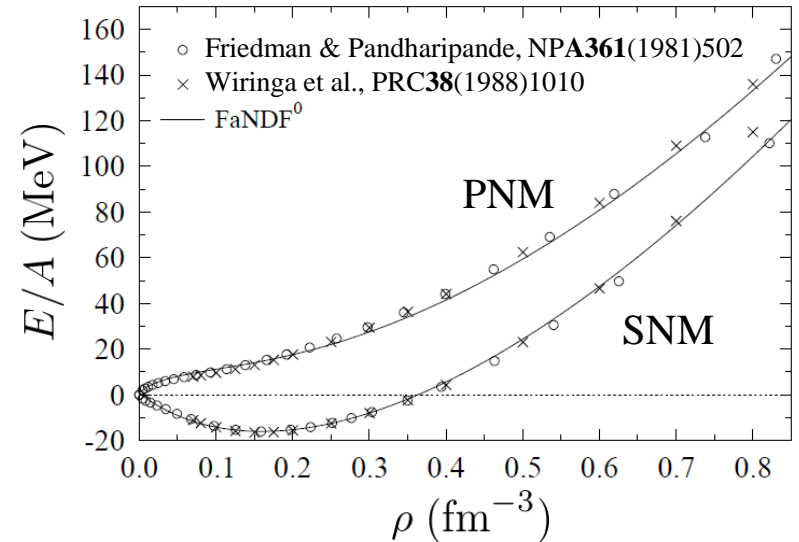
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□ Computational details

75 fm × 75 fm × 60 fm

(50 × 50 × 40, Δx = 1.5 fm)

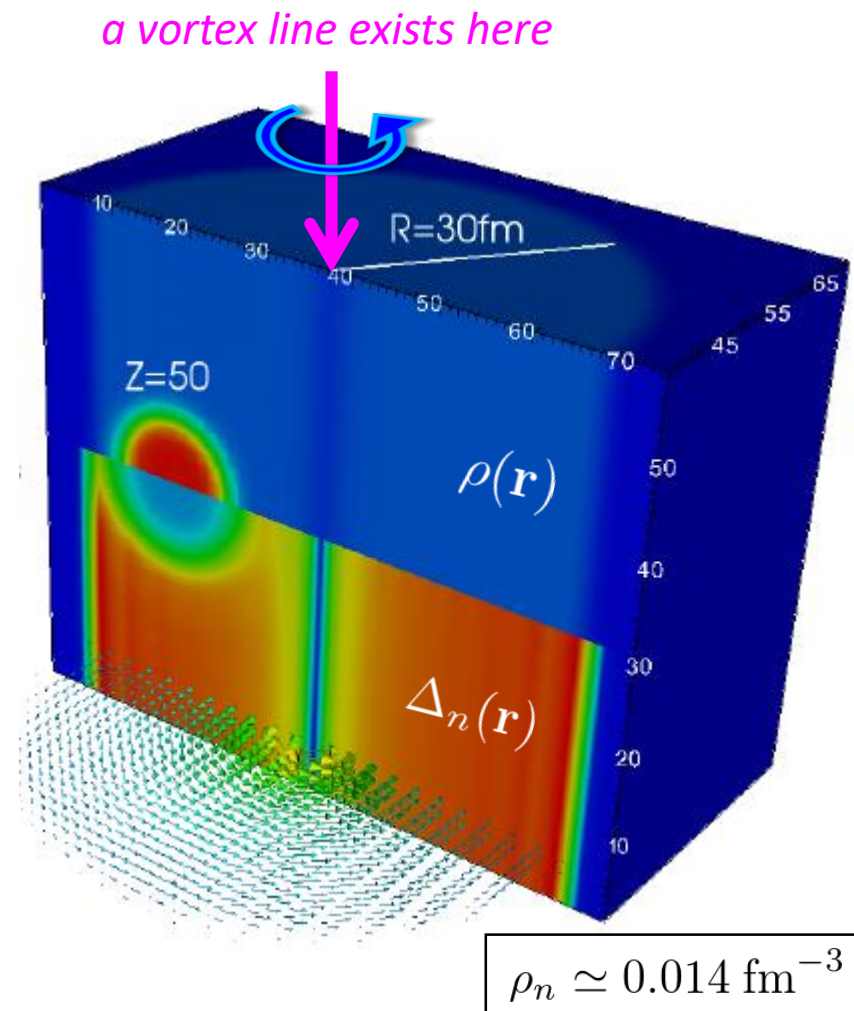
$k_c = \pi/\Delta x > k_F$ $k_F = (3\pi^2\rho_n)^{1/3}$

Nuclear impurity: Z = 50

$\rho_n \simeq 0.014 \text{ fm}^{-3}$ (N ≈ 2,530)

$\rho_n \simeq 0.031 \text{ fm}^{-3}$ (N ≈ 5,714)

of quasi-particle w.f. ≈ 100,000



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$\rho_n \simeq 0.031 \text{ fm}^{-3}$ ($N \simeq 5,714$)

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TITAN, Oak Ridge



NERSC Edison, Berkeley

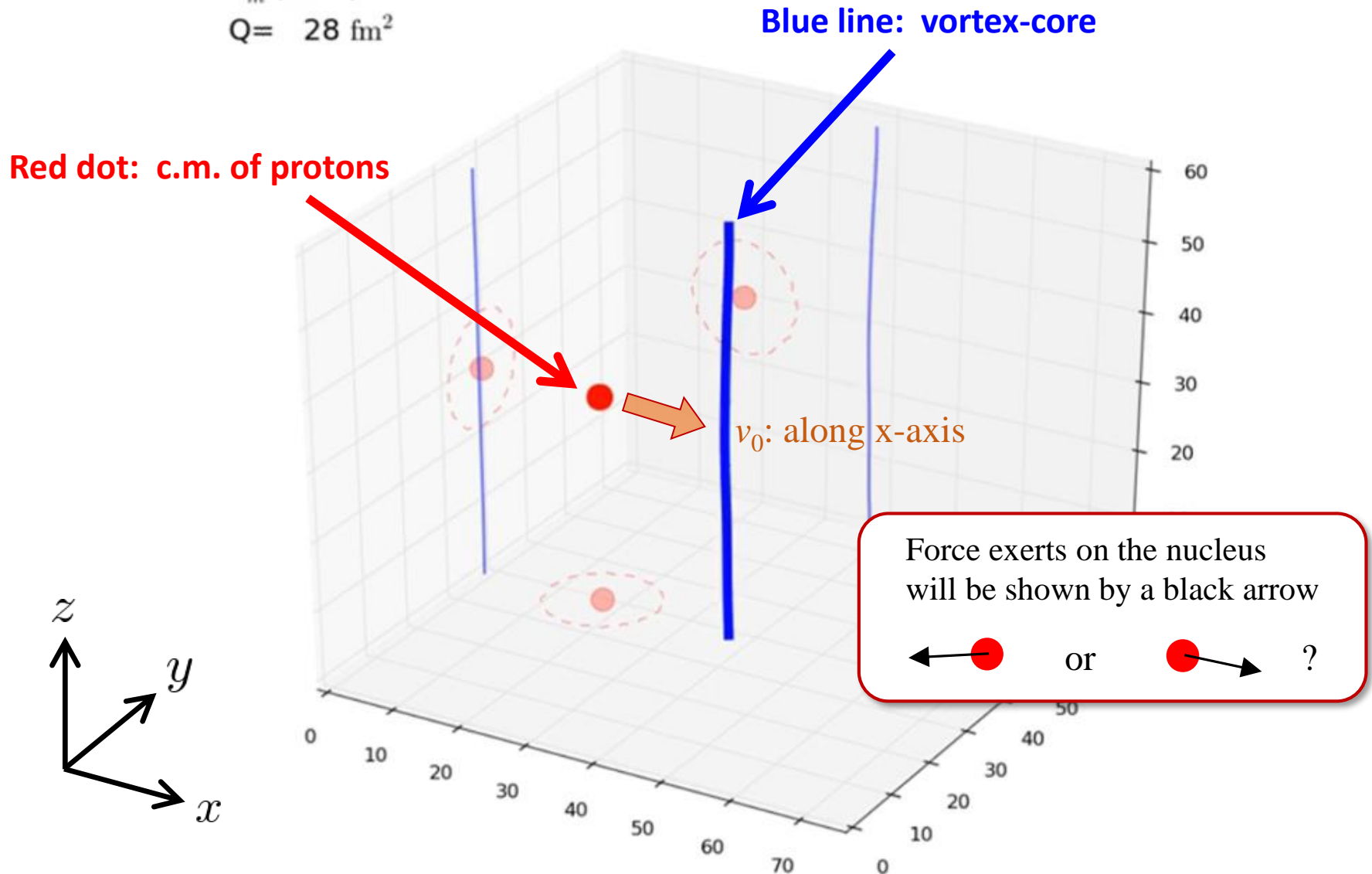


HA-PACS, Tsukuba

MPI+GPU
→ 48h w/ 200GPUs
for 10,000 fm/c

Results of TDSLDA calculation: $\rho_n \simeq 0.014 \text{ fm}^{-3}$

time= 0 fm/c
 $F_m(19.1)$ = unknown
 $Q = 28 \text{ fm}^2$



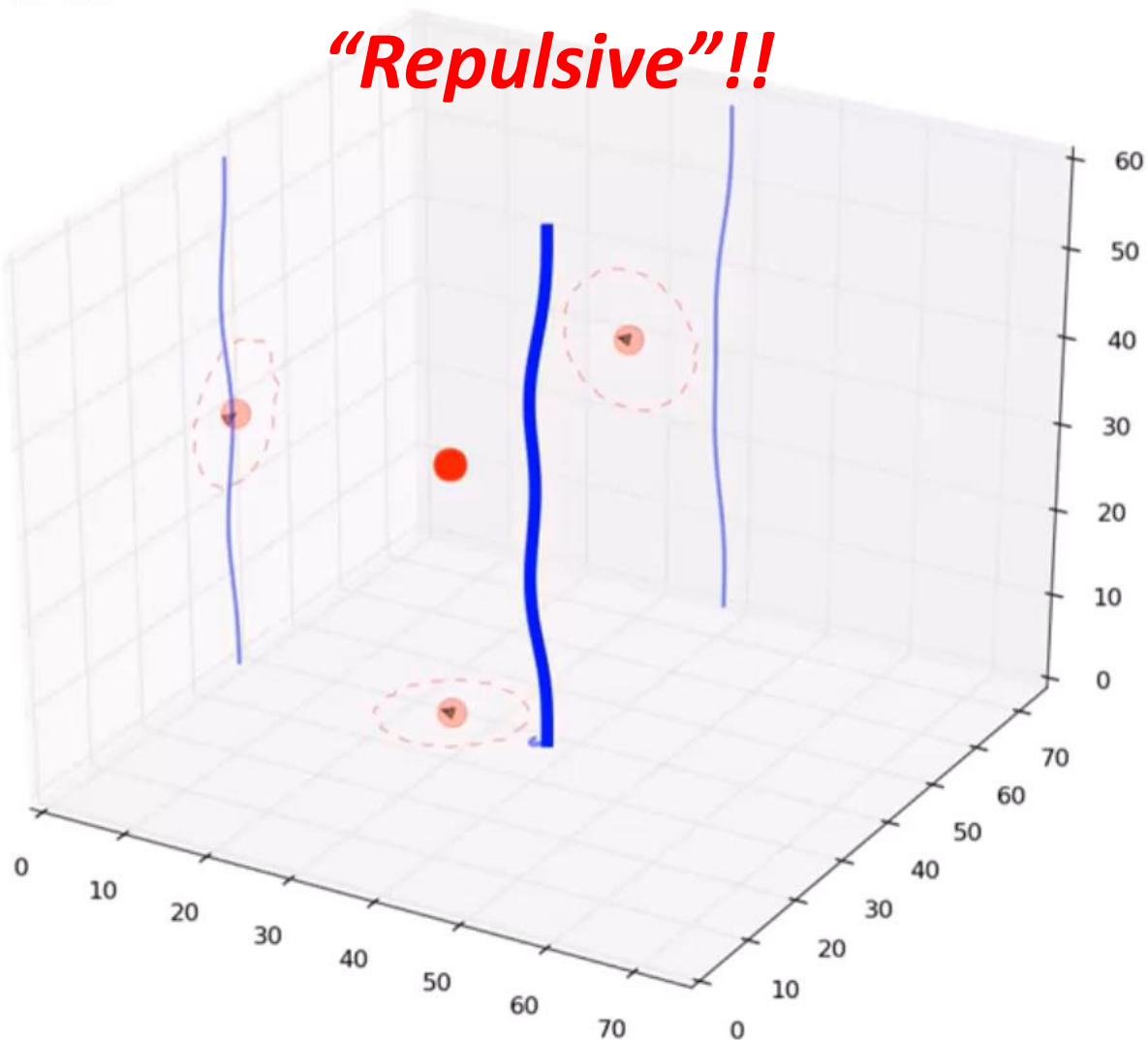
Results of TDSLDA calculation: $\rho_n \simeq 0.014 \text{ fm}^{-3}$

time= 8032 fm/c

$F_m(10.6) = 0.17 \text{ MeV/fm}$

$Q = 13 \text{ fm}^2$

“Repulsive”!!

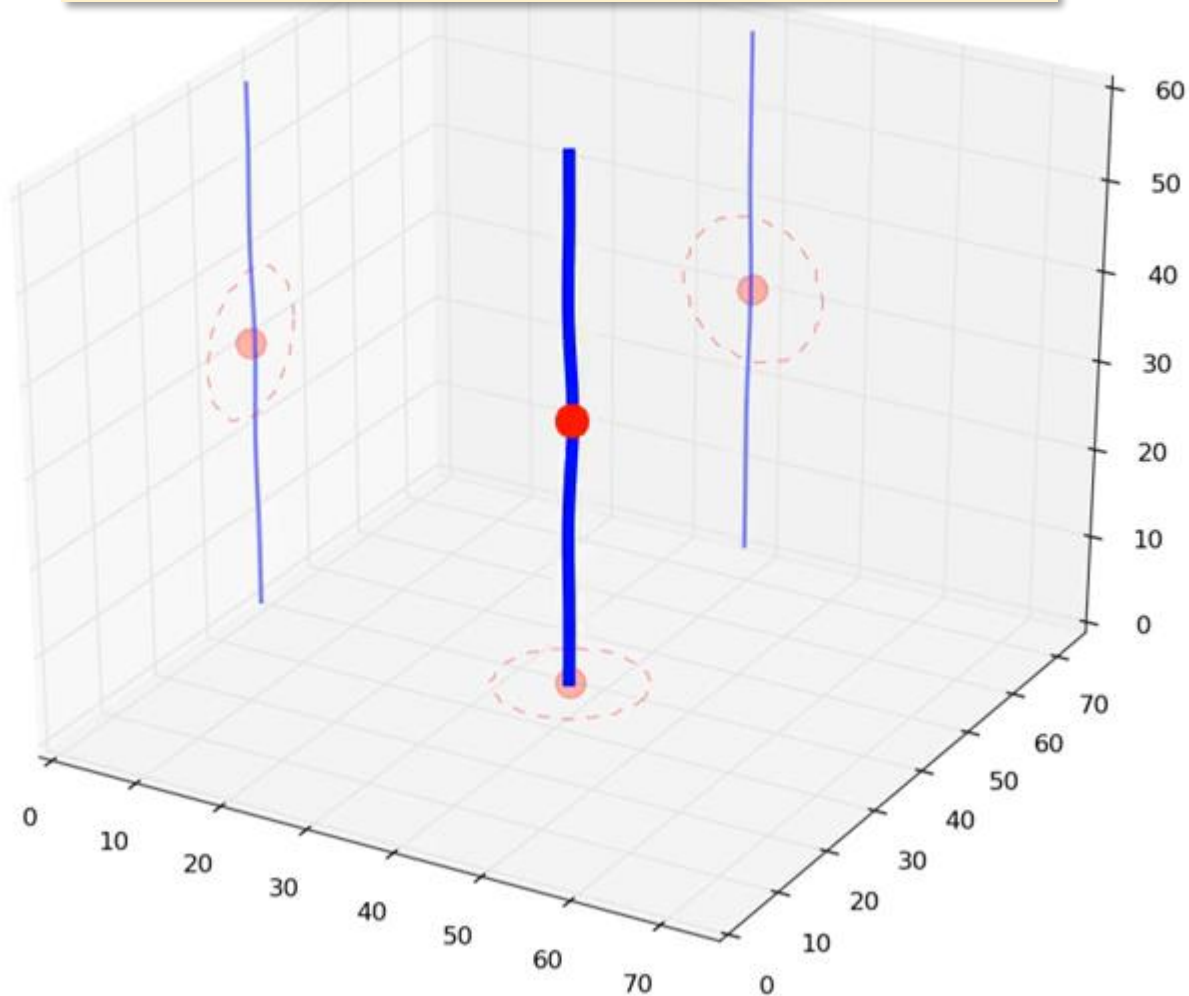


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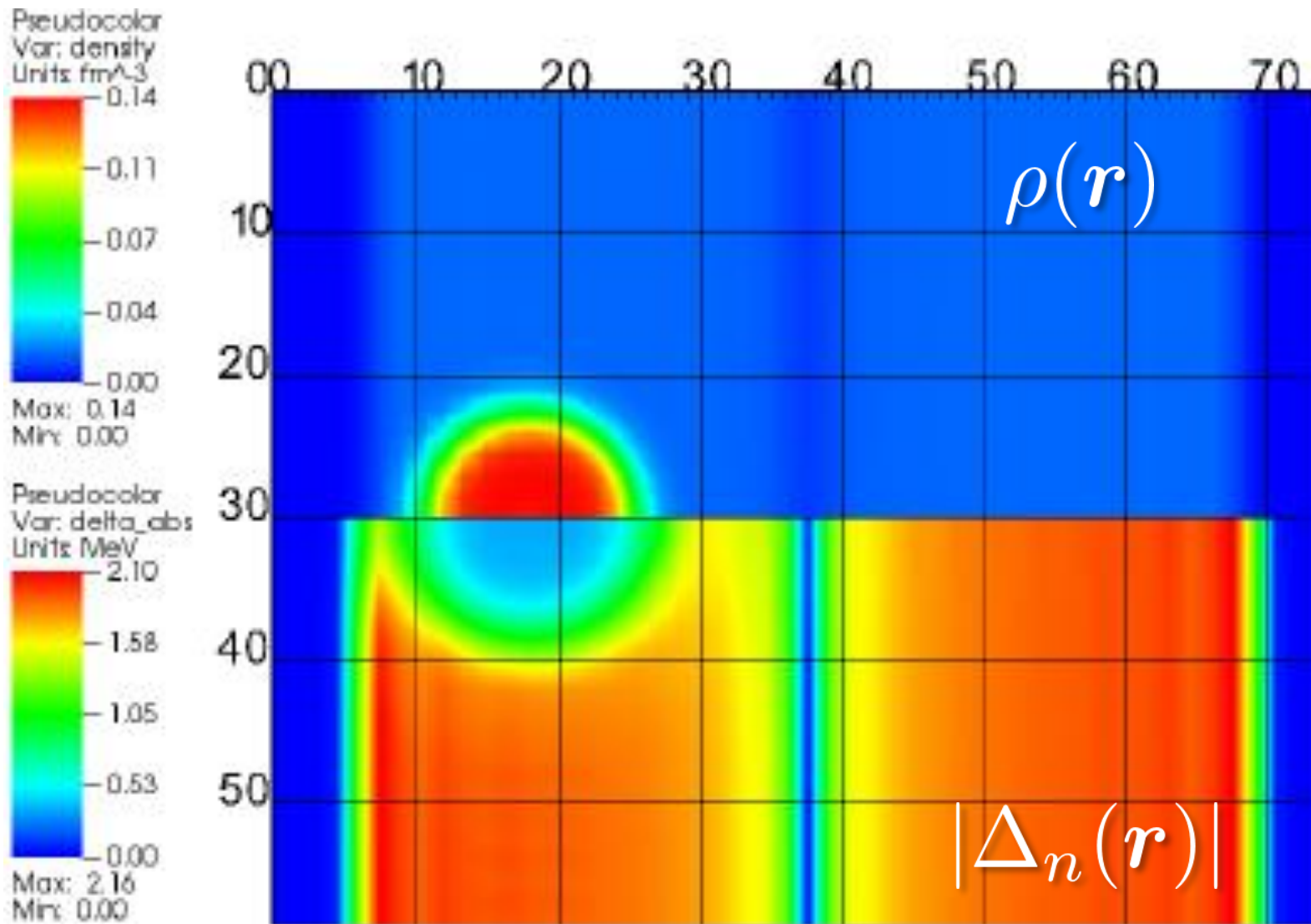
time= 0 fm/c

Q= -11 fm²

Pinned configuration is dynamically unstable

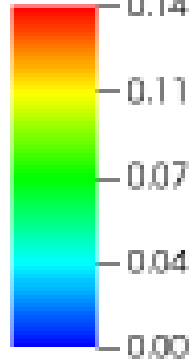


“Unpinned configuration”



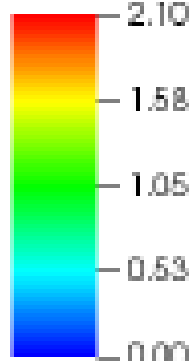
“Pinned configuration”

Pseudocolor
Var: density
Units: fm⁻³

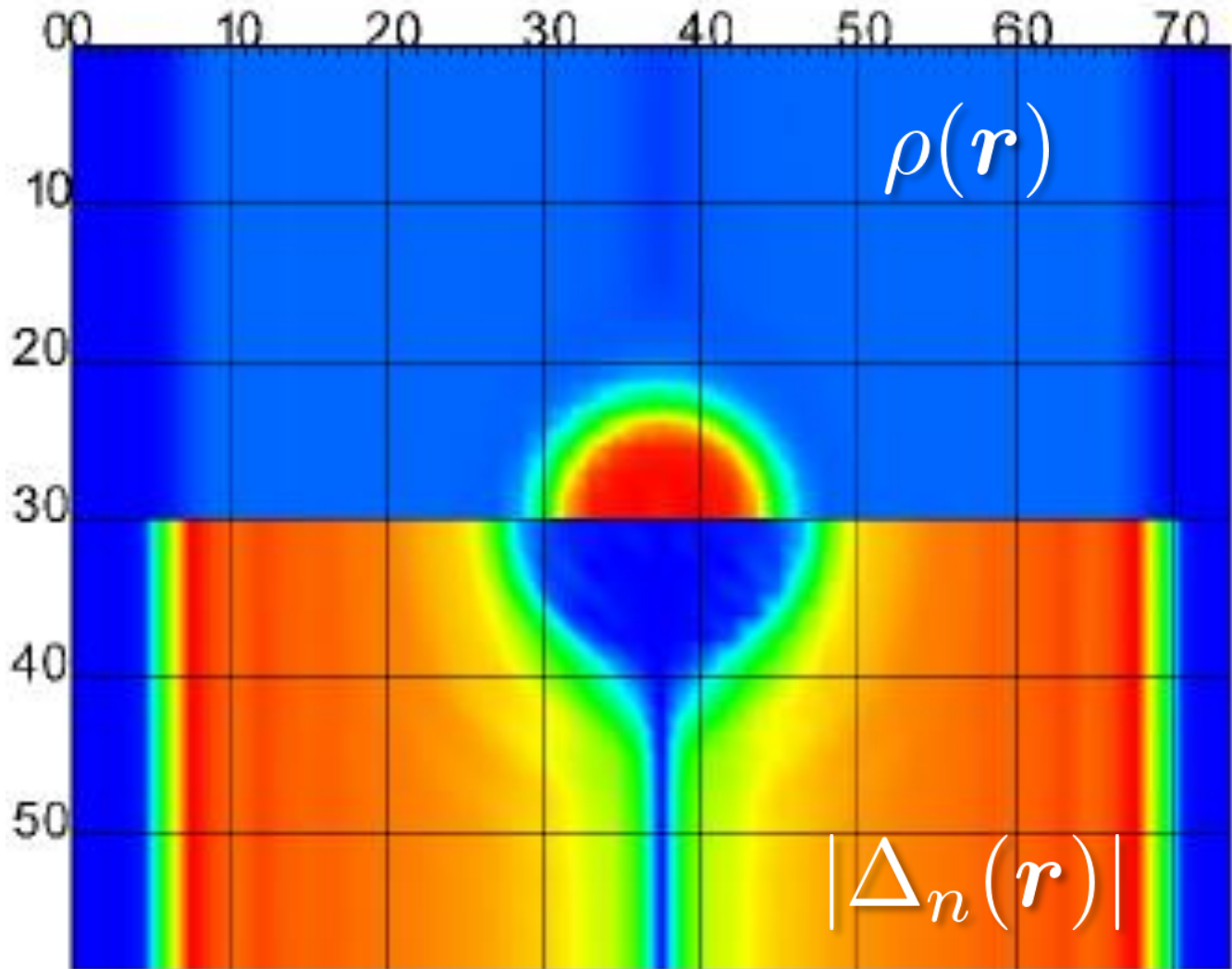


Max: 0.14
Min: 0.00

Pseudocolor
Var: delta_at
Units: MeV



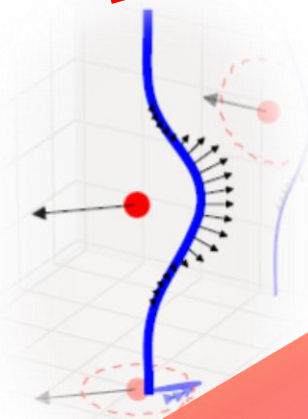
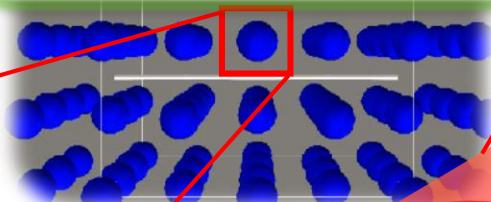
Max: 2.16
Min: 0.00



Our goal and strategy

Goal: Unveil the mechanism of glitches

New collaboration started:
Nicolaus Copernicus Astronomical Centre
B. Haskell et al.



10^{-15} - 10^{-13} m

$\sim 10^{-10}$ m

Mesoscopic

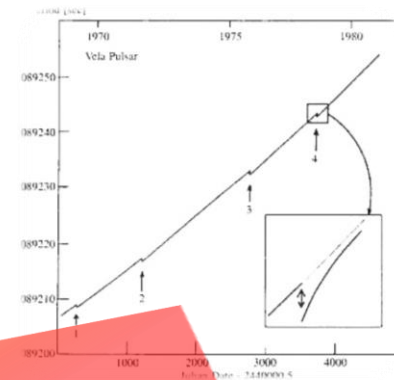
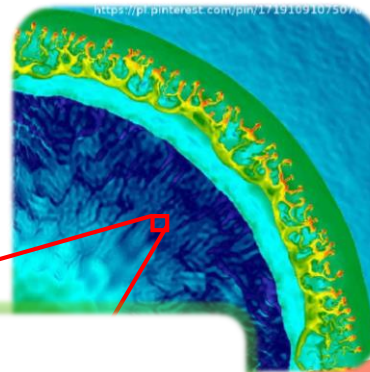
- dynamics of *vortices* in a lattice of *nuclei* (e.g. filament model)

Provide model ingredients

Microscopic

Nuclear Physics!!

- vortex-nucleus dynamics from *neutrons and protons*



10^4 m

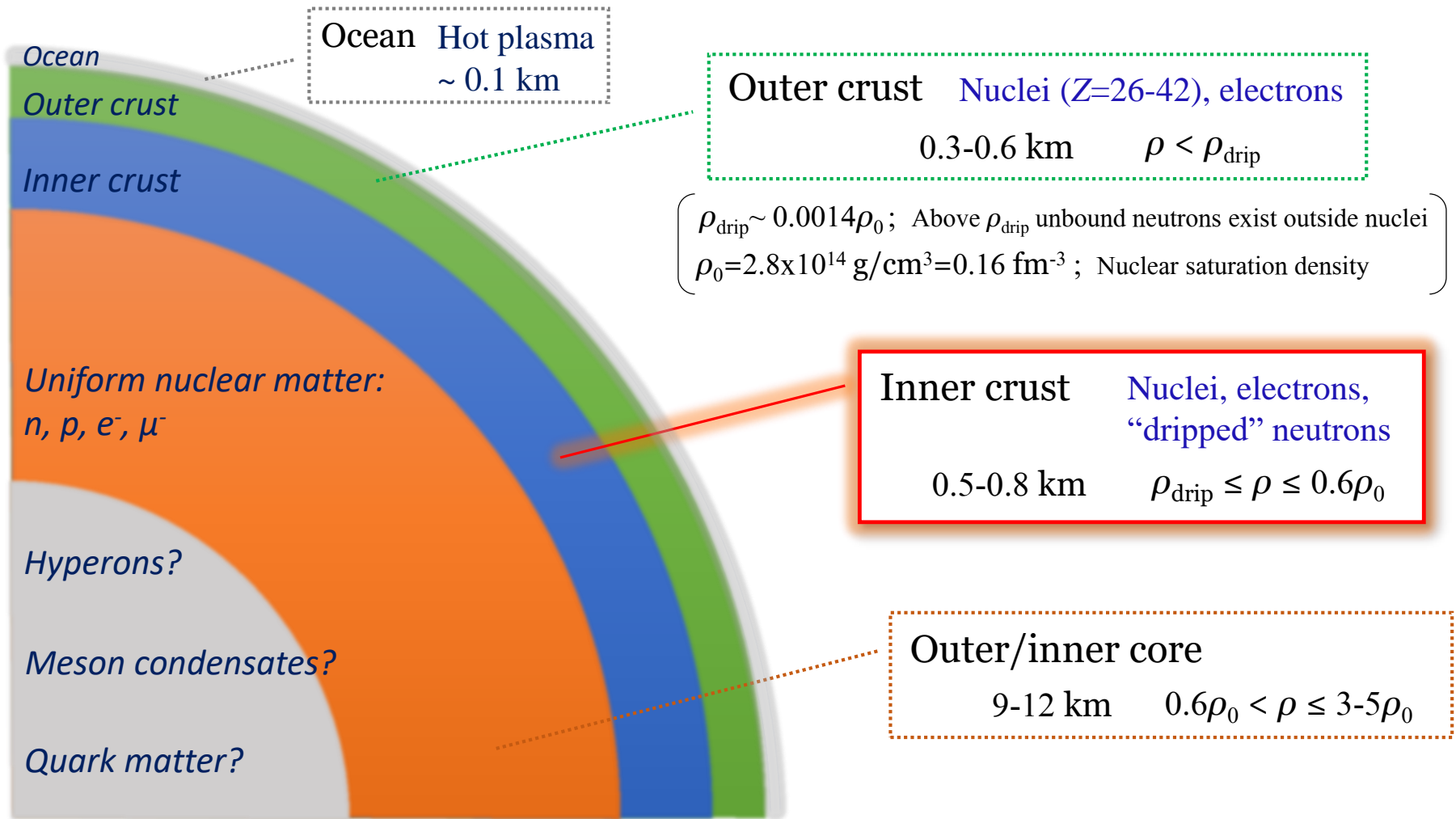
Macroscopic

- observations
- hydrodynamics

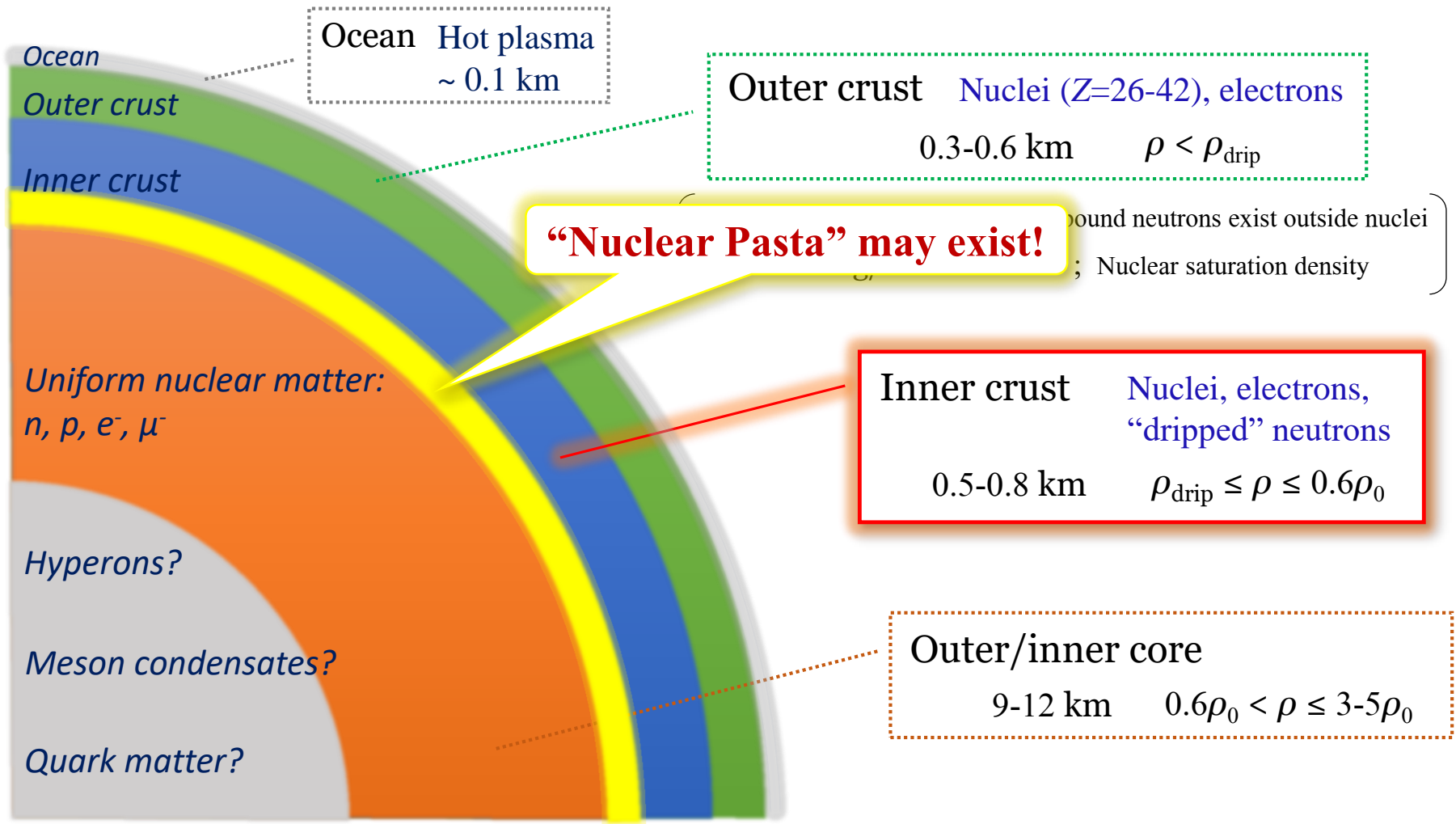
Time-
Dependent
Band Theory
for the
Inner Crust of
Neutron Stars



Neutron star is a great playground for nuclear physicists



Neutron star is a great playground for nuclear physicists



What is Nuclear Pasta?



Gnocchi

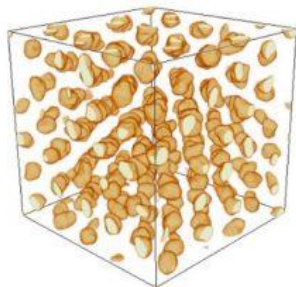


Lasagna

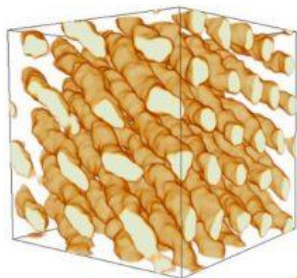


Spaghetti

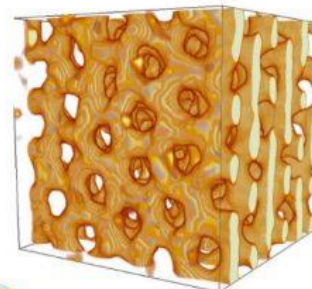
(a) *Gnocchi*



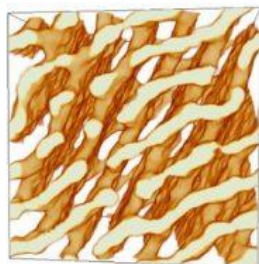
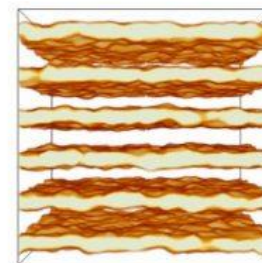
(b) *Spaghetti*



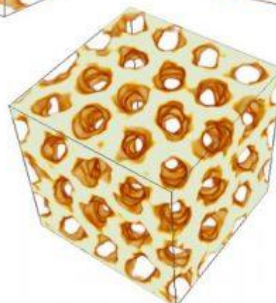
(c) *Waffles*



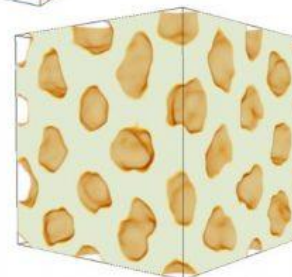
(d) *Lasagna*



(e) *Defects*



(f) *Antispaghetti*

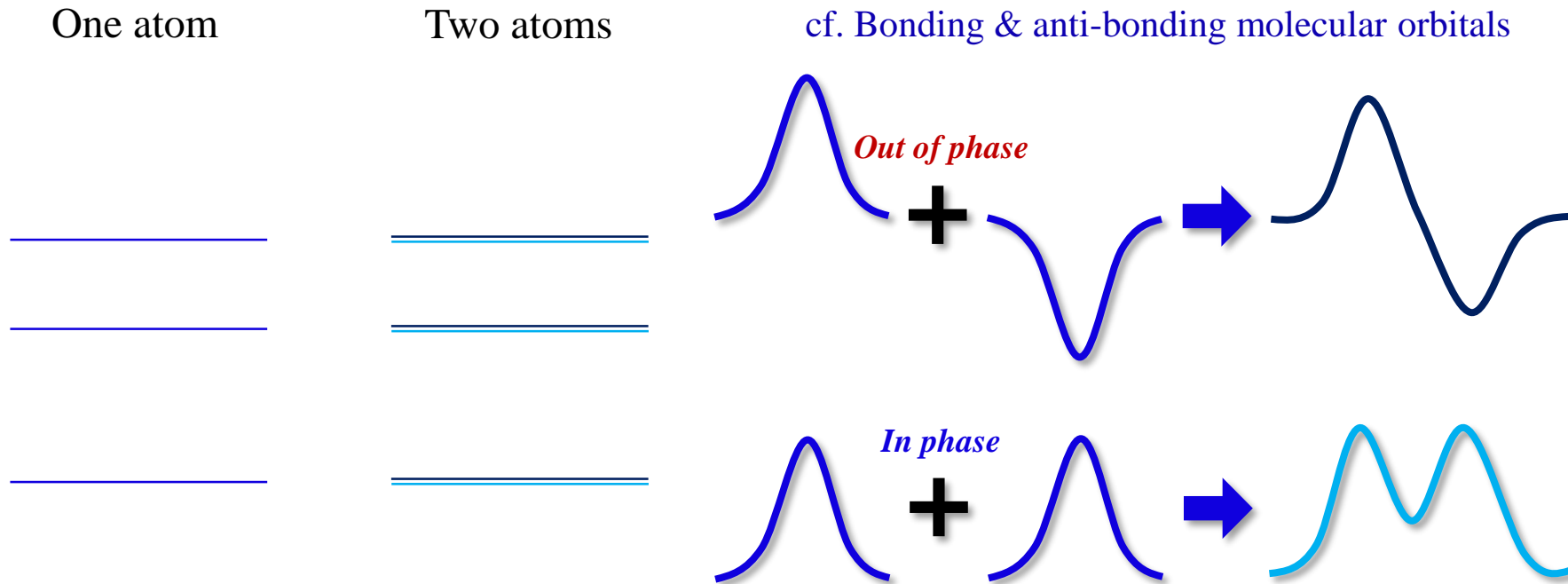


(g) *Antignocchi*

What is the “band structure” in solids?

An energy band: a bunch of shifted energy levels of atomic orbitals

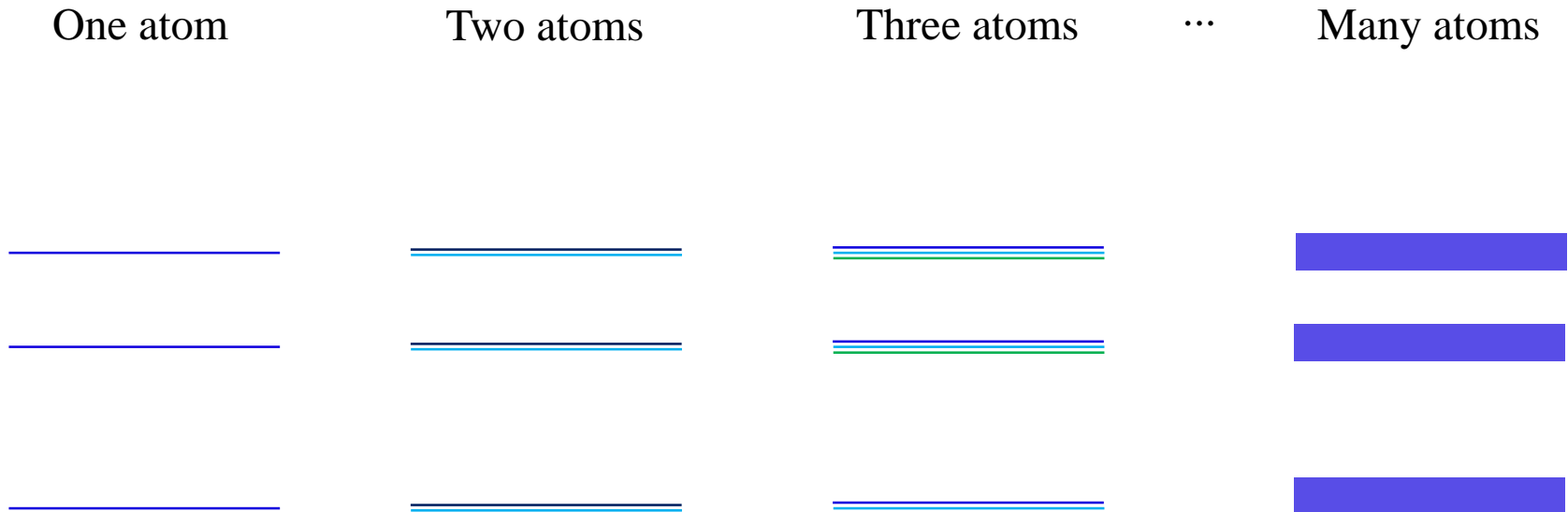
- ✓ Energy levels must be the same for each atom when those atoms are *completely separated*
- ✓ As they come closer, energy levels are spread out (cf. bonding/anti-bonding molecular orbital)
- ✓ The energy splitting depend on the inter-atomic distance (an optimum value is realized)



What is the “band structure” in solids?

An energy band: a bunch of shifted energy levels of atomic orbitals

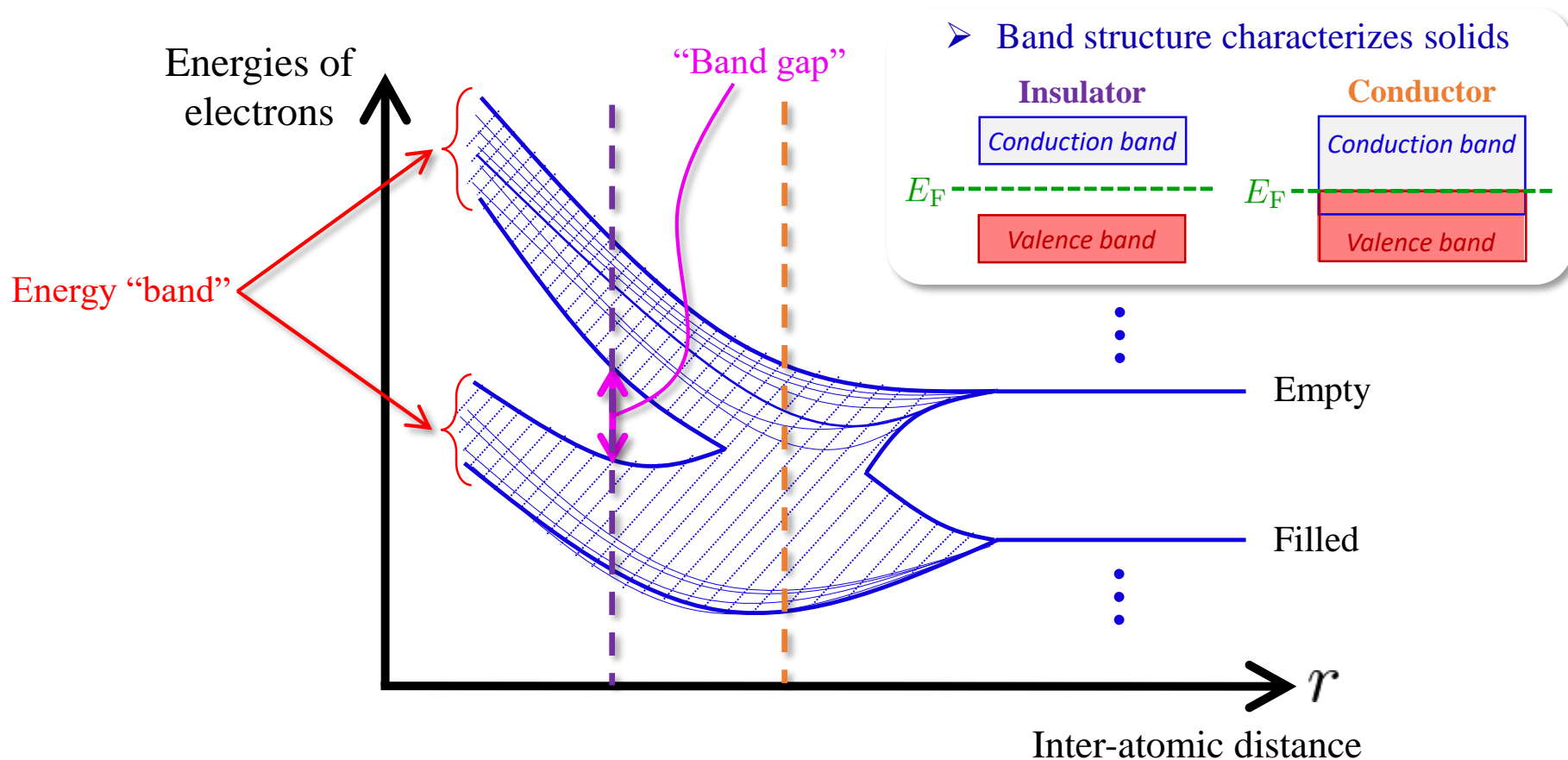
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The rest of the talk is based on one of my most recent publications:

PHYSICAL REVIEW C **105**, 045807 (2022)

Time-dependent extension of the self-consistent band theory for neutron star matter: Anti-entrainment effects in the slab phase

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²Nuclear Physics Division, Center for Computational Sciences, University of Tsukuba, Ibaraki 305-8577, Japan

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(Received 28 December 2021; accepted 4 April 2022; published 25 April 2022)

in collaboration with



Sorataka Kobayashi

(Finished MSc in Mar. 2019)



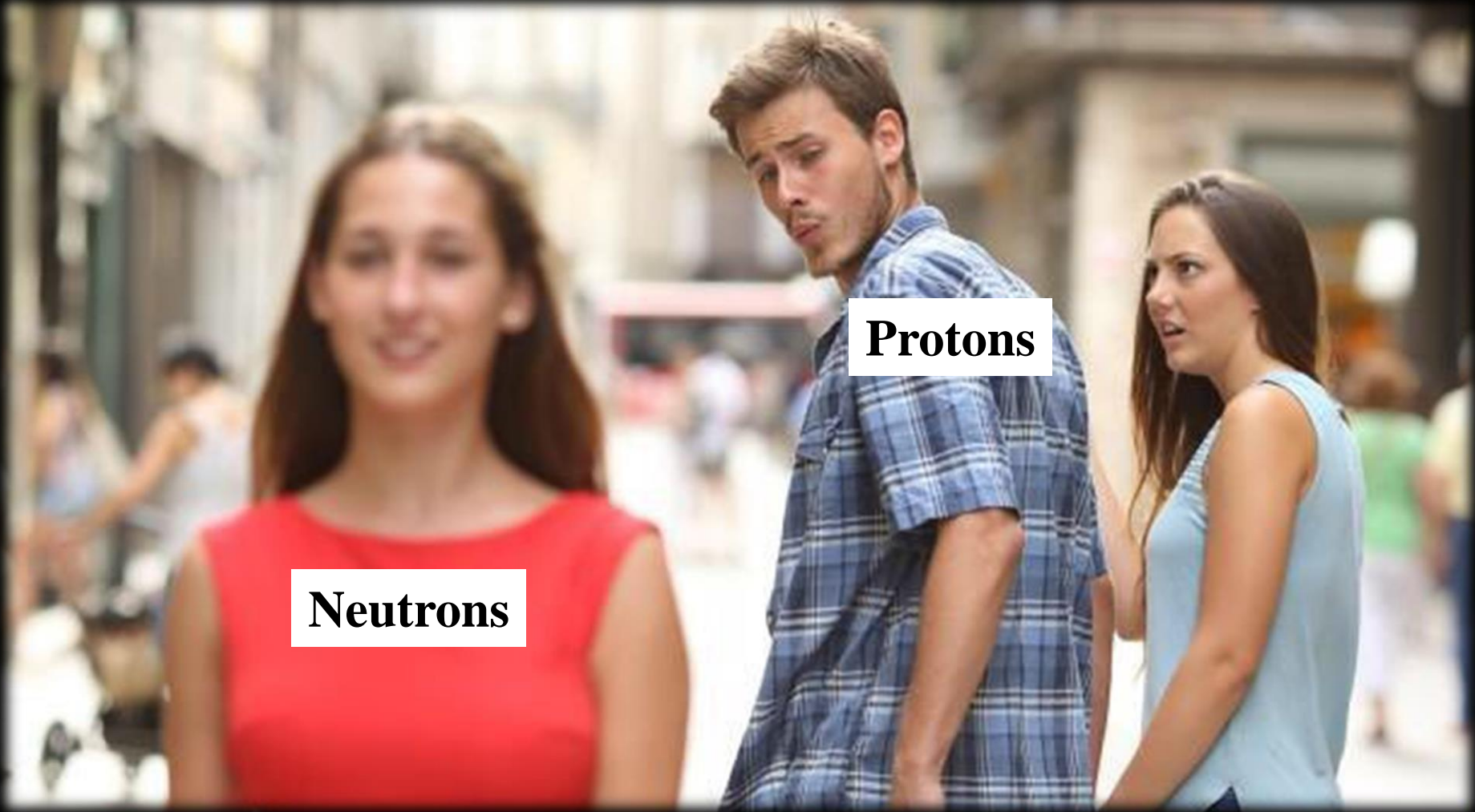
Masayuki Matsuo



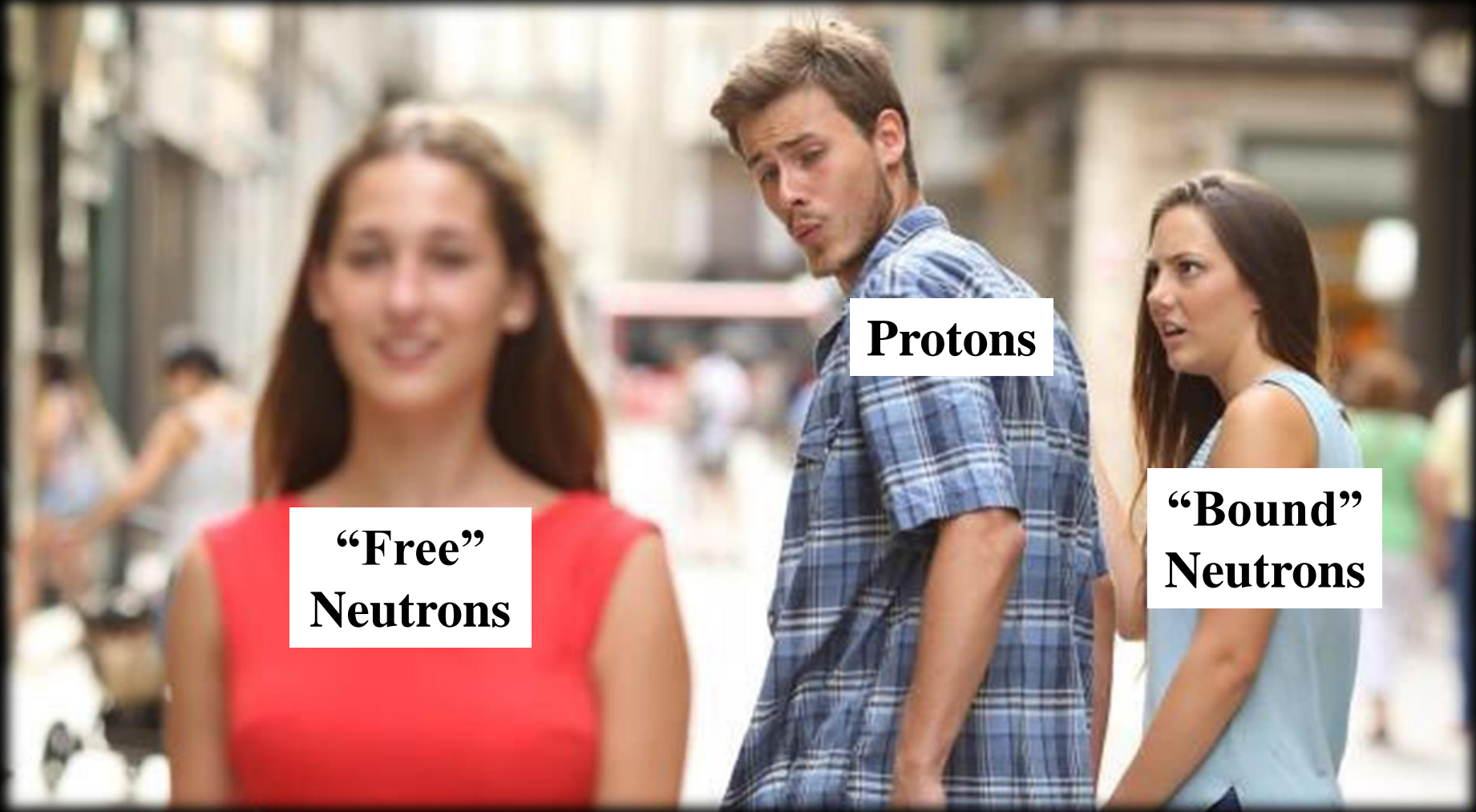
Kenta Yoshimura (M1)



“Entrainment” is a phenomenon between two species (particles, gases, fluids, etc.), where a motion of one component attracts the other.



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**“Free”
Neutrons**

Protons

**“Bound”
Neutrons**

“Entrainment” in the inner crust

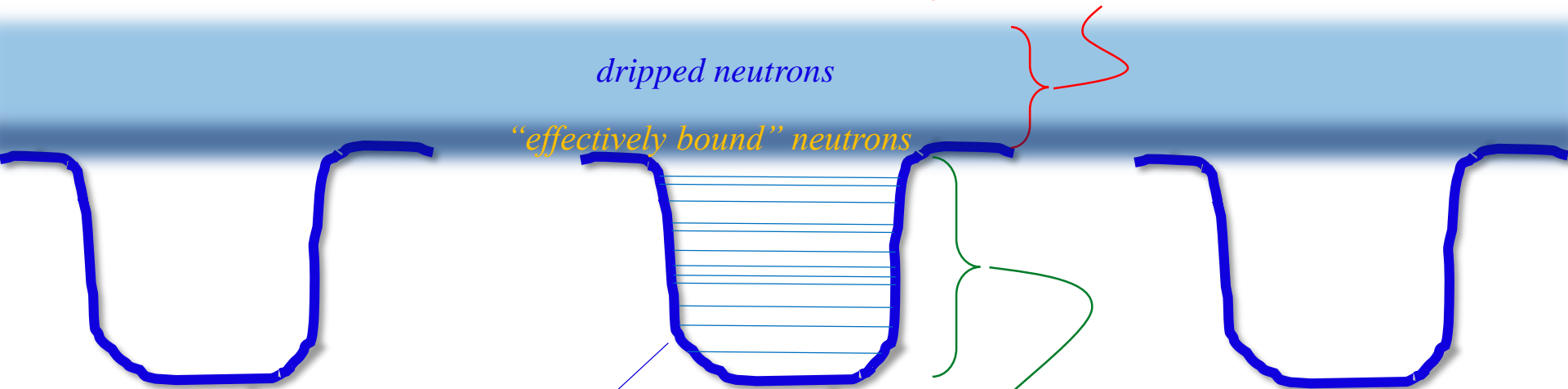
- Part of dripped neutrons are “effectively bound” (immobilized) by the periodic structure (due to Bragg scatterings), resulting in a larger effective mass

$$m_n n_n^f = m_n^* n_n^c$$

n_n^c : Conduction neutron number density
(neutrons that can actually flow)

m_n^* : (Macroscopic) Effective mass

Dripped neutrons extend spatially
→ Affected by the lattice, and a band structure is formed!



Entrainment leads:

- reduction of n_c
- enhancement of m^*

Potential for neutrons

Bound orbitals are well **localized**
→ Not affected by the lattice

The “entrainment effect” is still a debatable problem

- The first consideration for 1D, square-well potential

K. Oyamatsu and Y. Yamada, NPA**578**(1994)184

- Band calculations for slab (1D) and rod (2D) phases

B. Carter, N. Chamel, and P. Haensel, NPA**748**(2005)675

➡ Entrainment effects are **weak** for the slab & rod phases:

$$\frac{m^*}{m} \sim \begin{cases} 1.02 - 1.03 & \text{for the slab phase} \\ 1.11 - 1.40 & \text{for the rod phase} \end{cases}$$

- Band calculations for cubic-lattice (3D) phases

N. Chamel, NPA**747**(2005)109 (2005); NPA**773**(2006)263; PRC**85**(2012)035801; J. Low Temp. Phys. **189**, 328 (2017)

➡ **Significant** entrainment effects were found in a low-density region:

$$\frac{m^*}{m} \gtrsim 10 \text{ or more! for the cubic lattice}$$

- The first *self-consistent* band calculation for the slab (1D) phase (based on DFT with a BCPM EDF)

➡ “**Reduction**” of the effective mass was observed for the slab phase:

$$\frac{m^*}{m} \sim 0.65 - 0.75 \text{ for the slab phase}$$

Yu Kashiwaba and T. Nakatsukasa, PRC**100**(2019)035804

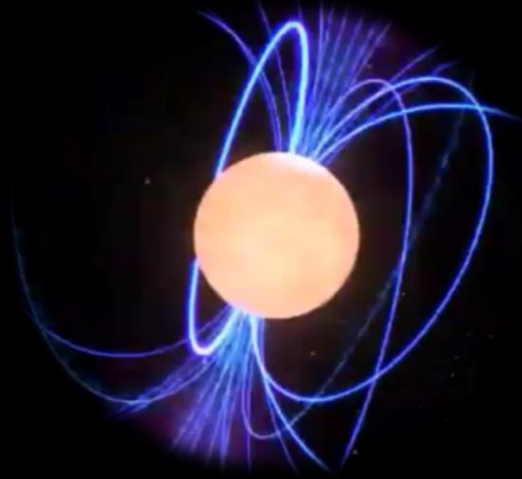
- **Time-dependent extension of the self-consistent band theory (based on TDDFT with a Skyrme EDF)**

➡ “**Reduction**” was observed, consistent with the Tsukuba group.

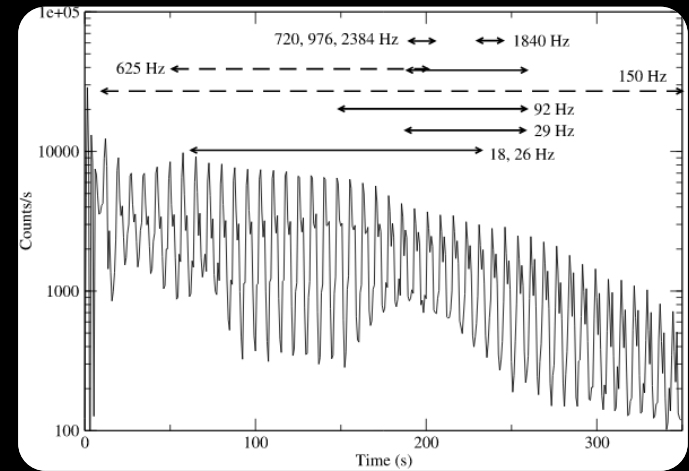
K. Sekizawa, S. Kobayashi, and M. Matsuo, PRC**105**(2022)045807

It may affect interpretation of various phenomena, e.g.:

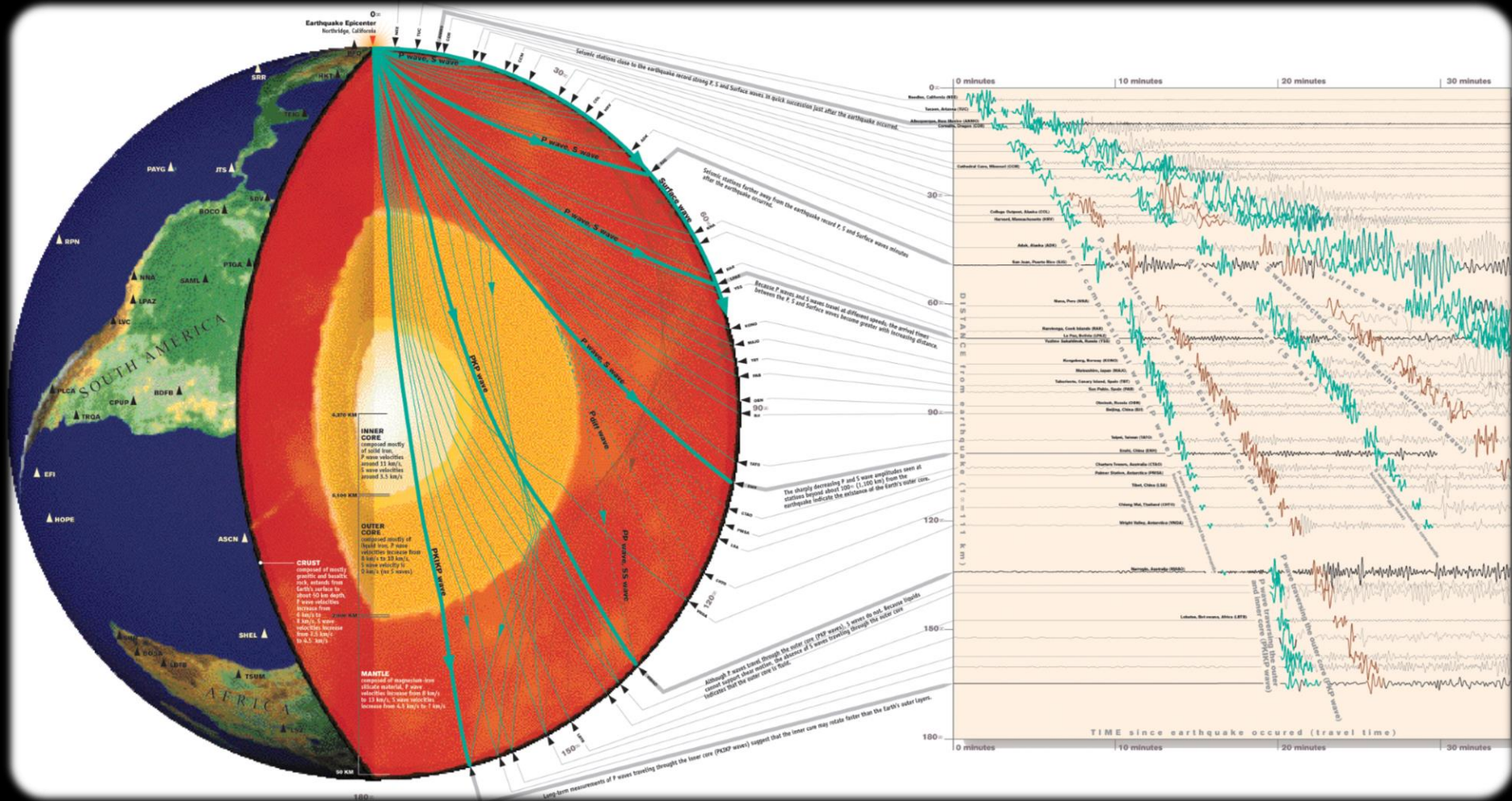
Neutron-star glitch



Quasi-periodic oscillation



Seismology (地震学): Studying inside of the Earth from earthquakes and their propagation





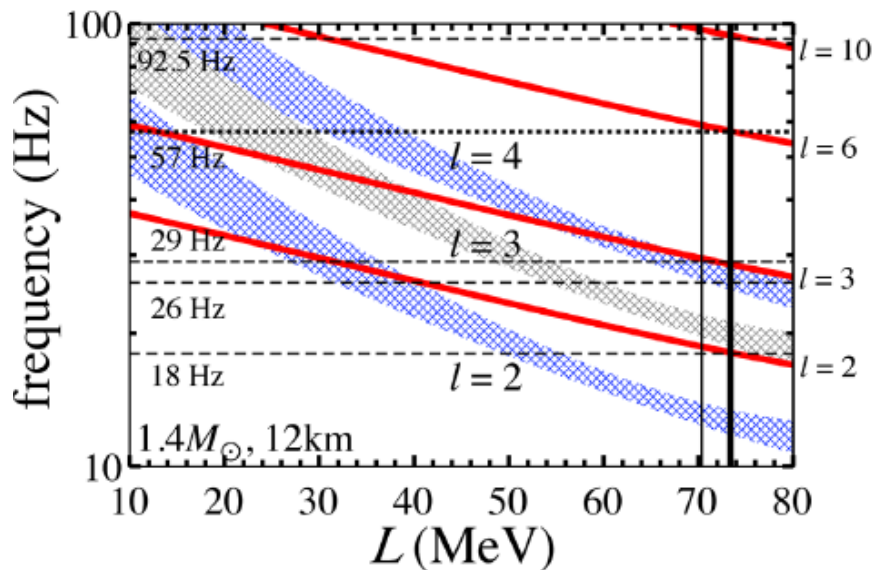
Astrophysical implications of double-layer torsional oscillations in a neutron star crust as a lasagna sandwich

Hajime Sotani¹,[★] Kei Iida² and Kazuhiro Oyamatsu³

¹Division of Science, National Astronomical Observatory of Japan, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan

²Department of Mathematics and Physics, Kochi University, 2-5-1 Akebono-cho, Kochi 780-8520, Japan

³Department of Human Informatics, Aichi Shukutoku University, 2-9 Katahira, Nagakute, Aichi 480-1197, Japan



➤ Many (~30) observed QPO frequencies, and prediction by a Bayesian analysis, have been nicely explained by torsional oscillations of tube–bubble or sphere–cylinder layer

Astrophysical implications of double-layer torsional oscillations in a neutron star crust as a lasagna sandwich

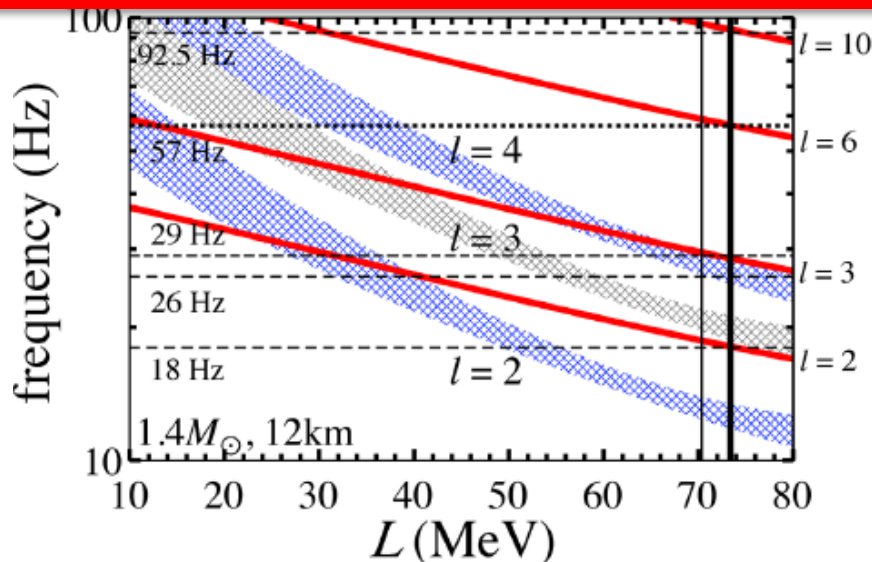
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The interpretation could be affected by the entrainment effects!



- Many (~30) observed QPO frequencies, and prediction by a Bayesian analysis, have been nicely explained by torsional oscillations of tube–bubble or sphere–cylinder layer

We employ the Skyrme-Kohn-Sham DFT with the Bloch boundary condition

✓ The Bloch boundary condition for single-particle orbitals

$$\psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r}) = \frac{1}{\sqrt{V}} u_{\alpha\mathbf{k}}^{(q)}(z) e^{i\mathbf{k}\cdot\mathbf{r}} \quad u_{\alpha\mathbf{k}}^{(q)}(z + na) = u_{\alpha\mathbf{k}}^{(q)}(z)$$

Periodicity of the slabs

α : Band index \mathbf{k} : Bloch wave vector q : Isospin (n or p) a : Period of the slabs

✓ Skyrme EDF

$$\frac{E}{A} = \frac{1}{N_b} \int_0^a \left(\frac{\hbar^2}{2m} \tau(z) + \sum_{t=0,1} \left[C_t^p [n] n_t^2(z) + C_t^{\Delta\rho} n_t(z) \partial_z^2 n_t(z) + C_t^{\tau} (n_t(z) \tau_t(z) - \mathbf{j}_t^2(z)) \right] + \mathcal{E}_{\text{Coul}}^{(p)}(z) \right) dz$$

Number density:

$$n_q(z) = 2 \sum_{\alpha,\mathbf{k}}^{\text{occ.}} |\psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r})|^2$$

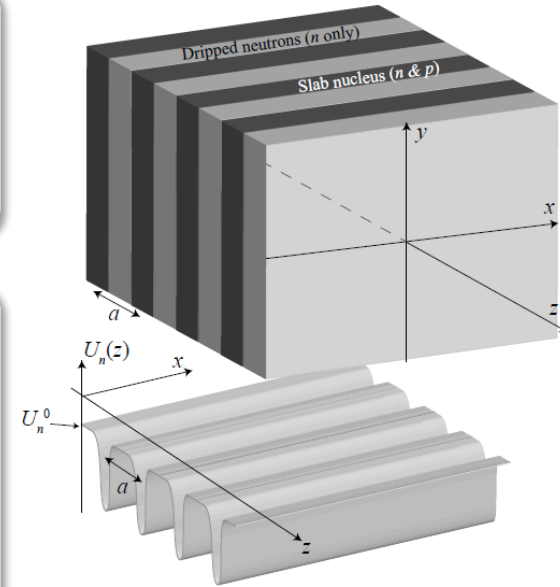
Kinetic density:

$$\tau_q(z) = 2 \sum_{\alpha,\mathbf{k}}^{\text{occ.}} |\nabla \psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r})|^2$$

Current (momentum) density:

$$\mathbf{j}_q(z) = 2 \sum_{\alpha,\mathbf{k}}^{\text{occ.}} \text{Im} [\psi_{\alpha\mathbf{k}}^{(q)*}(\mathbf{r}) \nabla \psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r})]$$

*Uniform background electrons are assumed for the charge neutrality condition: $n_e = \bar{n}_p$



Picture from PRC100(2019)035804

✓ Skyrme-Kohn-Sham equations

Note: While we deal with 3D slabs, the equations to be solved are 1D!

$$\hat{h}^{(q)}(z) \psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r}) = \varepsilon_{\alpha\mathbf{k}}^{(q)} \psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r}) \quad \rightarrow \quad \left(\hat{h}^{(q)}(z) + \hat{h}_{\mathbf{k}}^{(q)}(z) \right) u_{\alpha\mathbf{k}}^{(q)}(z) = \varepsilon_{\alpha\mathbf{k}}^{(q)} u_{\alpha\mathbf{k}}^{(q)}(z)$$

Ordinary single-particle Hamiltonian:

$$\hat{h}^{(q)}(z) = -\nabla \cdot \frac{\hbar^2}{2m_q^\oplus(z)} \nabla + U^{(q)}(z) + \frac{1}{2i} [\nabla \cdot \mathbf{I}^{(q)}(z) + \mathbf{I}^{(q)}(z) \cdot \nabla]$$

Additional (k -dependent) term:

$$\hat{h}_{\mathbf{k}}^{(q)}(z) = \frac{\hbar^2 \mathbf{k}^2}{2m_q^\oplus(z)} + \hbar \mathbf{k} \cdot \hat{\mathbf{v}}^{(q)}(z)$$

Velocity operator:

$$\hat{\mathbf{v}}^{(q)}(z) \equiv \frac{1}{i\hbar} [\mathbf{r}, \hat{h}^{(q)}(z)]$$

Proton fraction:

$$Y_p = \frac{\bar{n}_p}{\bar{n}_n + \bar{n}_p}$$

Average nucleon density:

$$\bar{n}_q = \frac{1}{a} \int_0^a n_q(z) dz$$

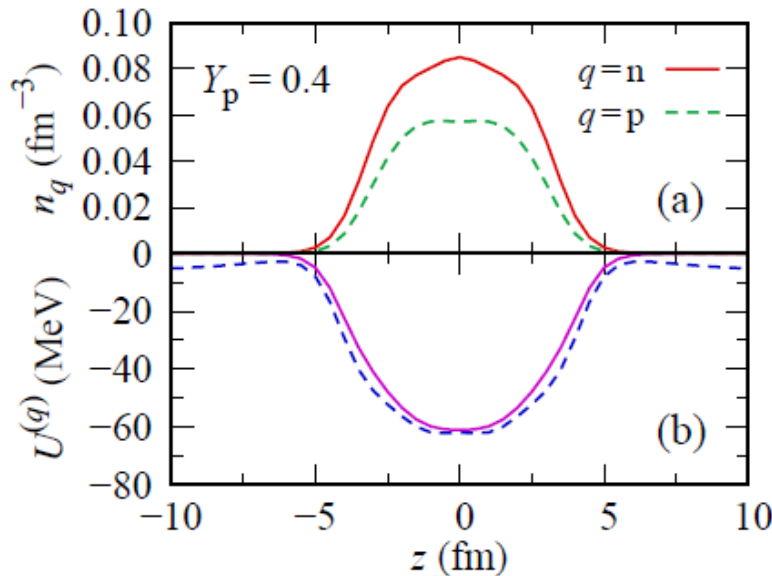
Single-particle energy:

$$\varepsilon_{\alpha\mathbf{k}}^{(q)} = \underbrace{e_{\alpha\mathbf{k}}^{(q)}}_{z\text{-component}} + \underbrace{\varepsilon_{\text{kin-xy},\alpha\mathbf{k}}^{(q)}}_{\approx \frac{\hbar^2 k_{\parallel}^2}{2m}} \quad k_{\parallel} = \sqrt{k_x^2 + k_y^2}$$

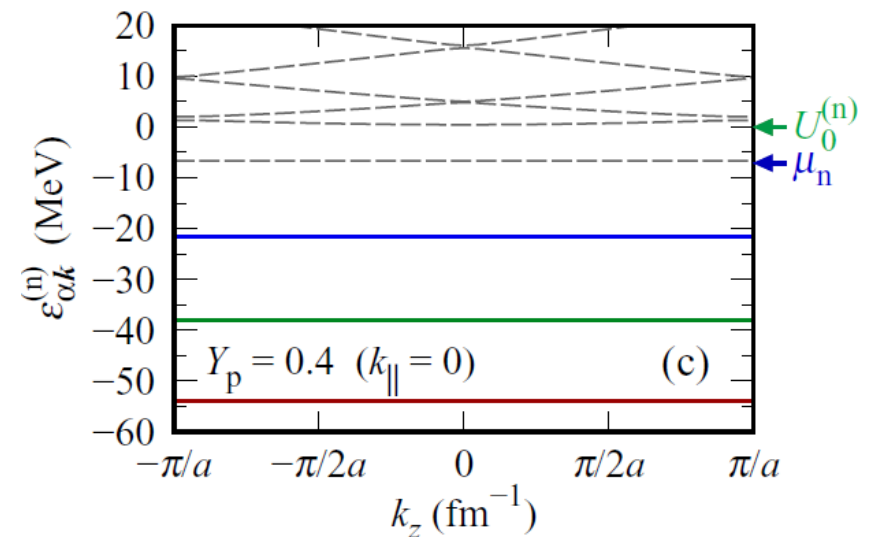
✓ Bound orbitals do not show band structure (k_z dependence)

$Y_p = 0.4, n_B = 0.4 \text{ fm}^{-3}$: Isolated slab (no neutron drip)

Density and potential



Neutron single-particle energies



Proton fraction:

$$Y_p = \frac{\bar{n}_p}{\bar{n}_n + \bar{n}_p}$$

Average nucleon density:

$$\bar{n}_q = \frac{1}{a} \int_0^a n_q(z) dz$$

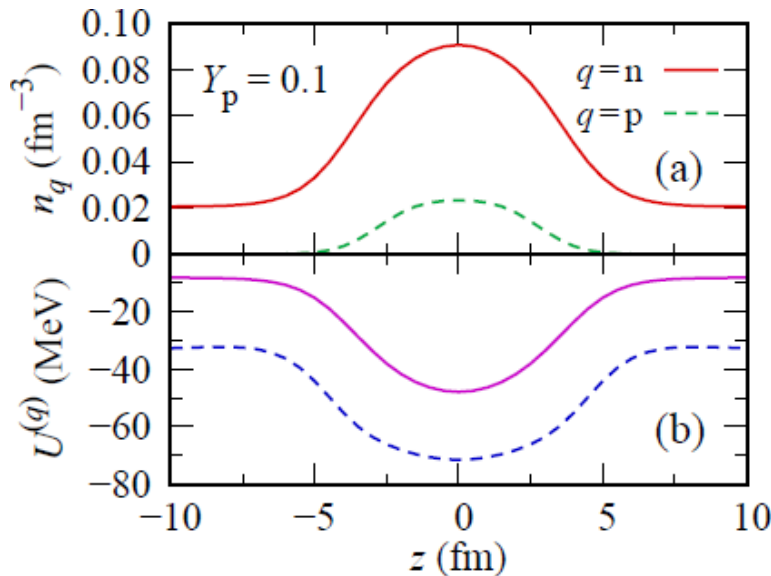
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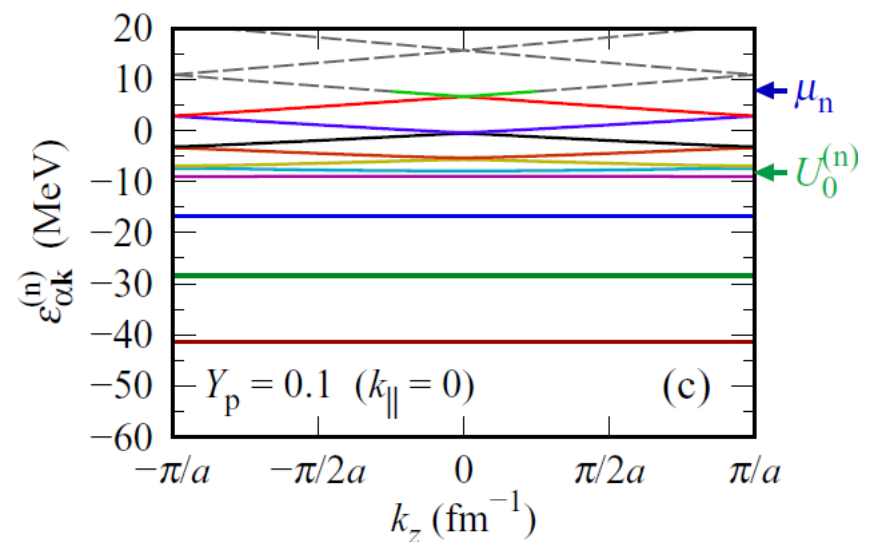
✓ Dripped neutrons show band structure (k_z dependence)

$Y_p = 0.1, n_B = 0.4 \text{ fm}^{-3}$: Neutron-dripped slab

Density and potential



Neutron single-particle energies



Proton fraction:

$$Y_p = \frac{\bar{n}_p}{\bar{n}_n + \bar{n}_p}$$

Average nucleon density:

$$\bar{n}_q = \frac{1}{a} \int_0^a n_q(z) dz$$

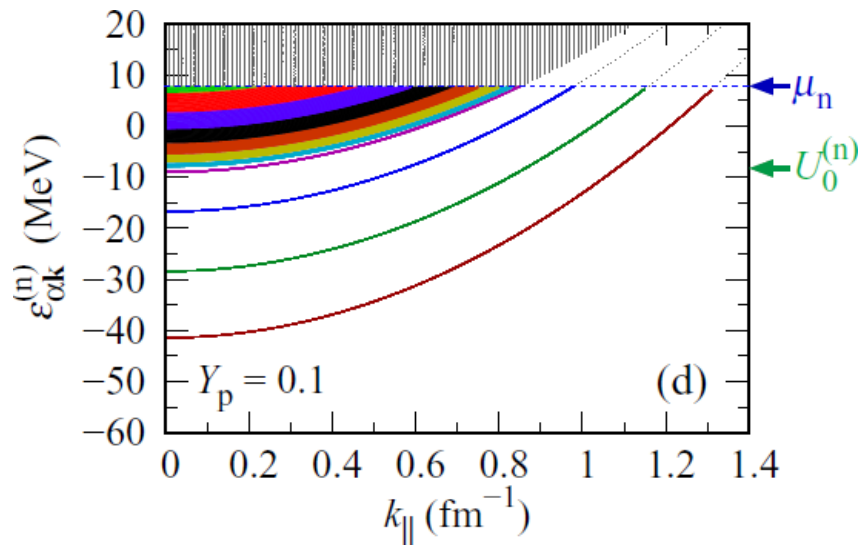
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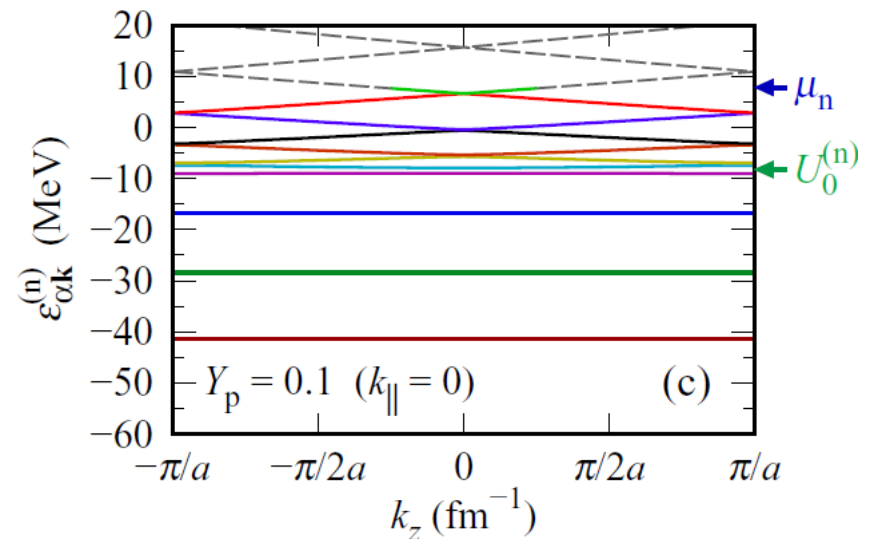
✓ Dripped neutrons show band structure (k_z dependence)

$Y_p = 0.1, n_B = 0.4 \text{ fm}^{-3}$: Neutron-dripped slab

k_{\parallel} dependence of s.p. energies:



Neutron single-particle energies



Static approach for conduction neutrons

- ✓ In the static approach, **conduction neutrons** are analyzed

In the **static** approach, the *conduction neutron number density* is defined by

$$n_n^c \equiv m_{n,\text{bg}}^\oplus \mathcal{K}_{zz}^{(n)}$$

where $\mathcal{K}_{zz}^{(n)}$ is the so-called *mobility coefficient*:

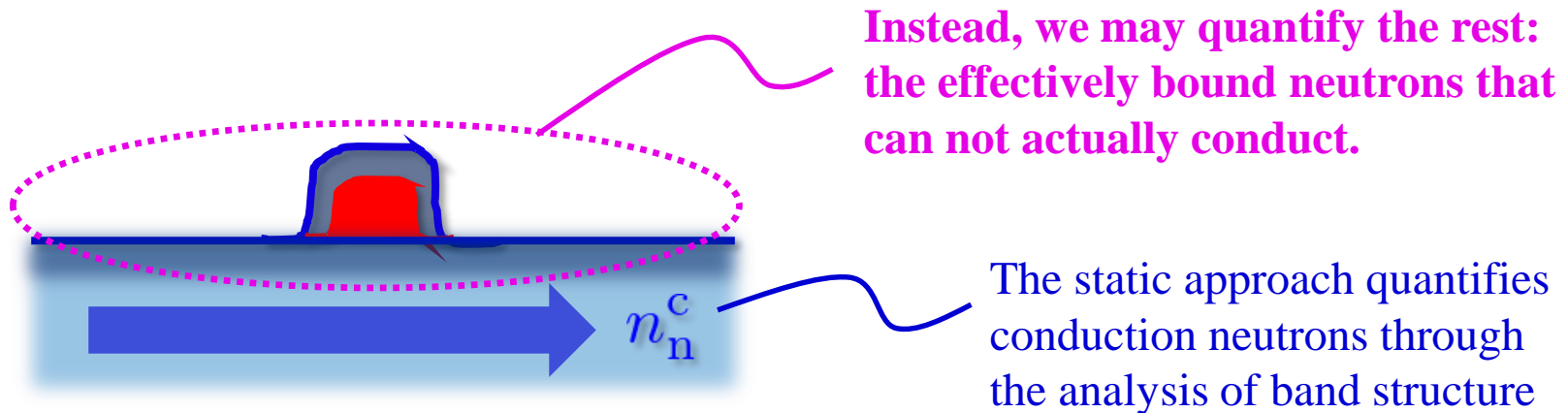
$$\mathcal{K}_{zz}^{(n)} = \frac{1}{\pi L} \sum_{\alpha, k_z} \int k_{\parallel} \left(m_{n,\alpha\mathbf{k}}^{\star-1} \right)_{zz} \theta(\mu_n - \varepsilon_{\alpha\mathbf{k}}^{(n)}) dk_{\parallel}$$

Inverse of the “macroscopic” effective mass tensor

$$\left(m_{n,\alpha\mathbf{k}}^{\star-1} \right)_{\mu\nu} = \frac{1}{\hbar^2} \frac{\partial^2 \varepsilon_{\alpha\mathbf{k}}^{(n)}}{\partial k_{\mu} \partial k_{\nu}}$$

For bound orbitals, there is no k_z dependence $\Rightarrow 1/m \rightarrow 0$, i.e., $m \rightarrow \infty$ (can not conduct).
 \Rightarrow The mobility coefficient quantifies dripped neutrons that can actually conduct.

Let's look at the same phenomenon
from a different side



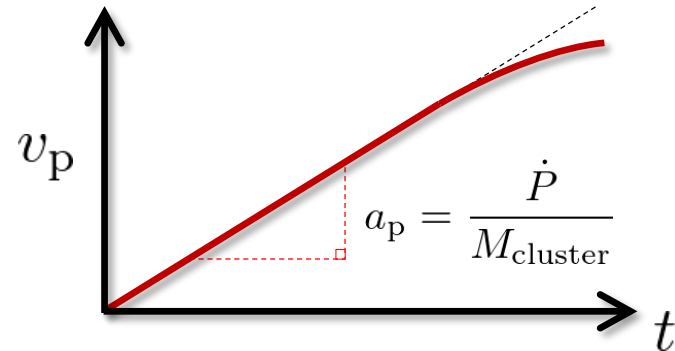
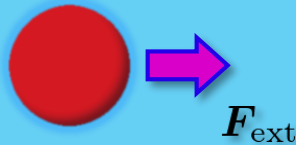




- ✓ The collective mass is extracted from **acceleration motion under constant force**

The real-time method: Idea

Dripped neutrons



How to introduce spatially-uniform electric field

- ✓ TDKS equation in a “velocity gauge”

$$i\hbar \frac{\partial \tilde{u}_{\alpha\mathbf{k}}^{(q)}(z, t)}{\partial t} = \left(\hat{h}^{(q)}(z, t) + \hat{h}_{\mathbf{k}(t)}^{(q)}(z, t) \right) \tilde{u}_{\alpha\mathbf{k}}^{(q)}(z, t)$$

Spatially-uniform
Vector potential

$$\mathbf{k}(t) = \mathbf{k} + \frac{e}{\hbar c} A_z(t) \hat{\mathbf{e}}_z$$

Gauge transformation for the Bloch orbitals:

$$\tilde{u}_{\alpha\mathbf{k}}^{(q)}(z, t) = \exp\left[-\frac{ie}{\hbar c} A_z(t) z\right] u_{\alpha\mathbf{k}}^{(q)}(z, t)$$

Electric field:

$$E_z(t) = -\frac{1}{c} \frac{dA_z}{dt}$$

k -dependent term:

$$\hat{h}_{\mathbf{k}}^{(q)}(z) = \frac{\hbar^2 \mathbf{k}^2}{2m_q^\oplus(z)} + \hbar \mathbf{k} \cdot \hat{\mathbf{v}}^{(q)}(z)$$

Velocity operator:

$$\hat{\mathbf{v}}^{(q)}(z) \equiv \frac{1}{i\hbar} [\mathbf{r}, \hat{h}^{(q)}(z)]$$

cf. K. Yabana and G.F. Bertsch, Phys. Rev. B **54**, 4484 (1996); G.F. Bertsch *et al.*, Phys. Rev. B **62**, 7998 (2000)

Results: The collective mass

Acceleration:

$$a_p = \frac{d^2 Z}{dt^2}$$

C.m. position of protons:

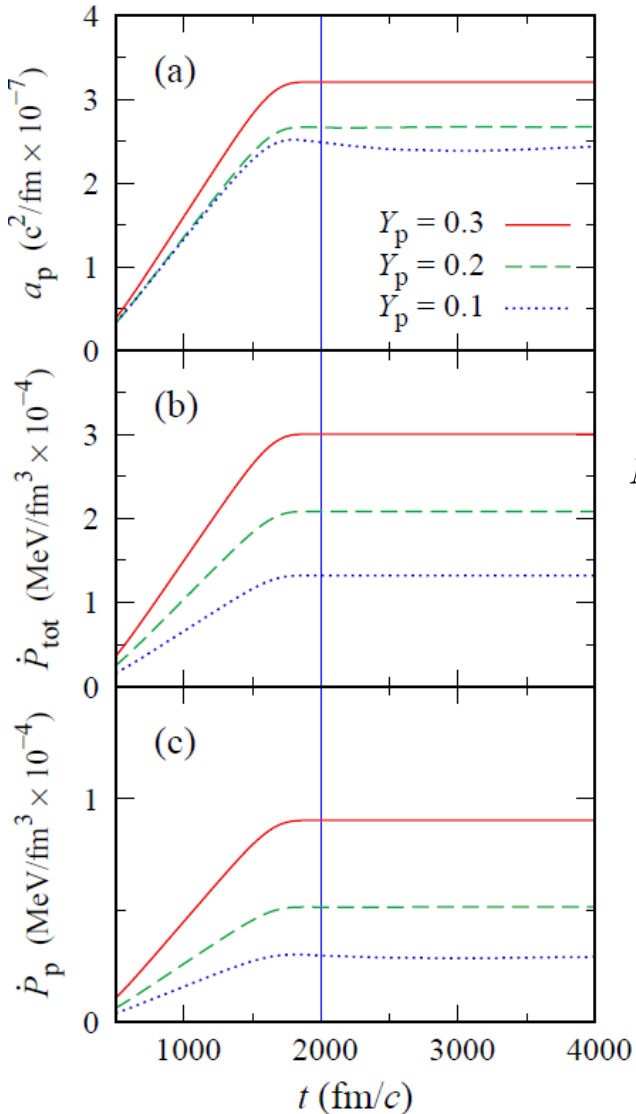
$$Z(t) = \frac{1}{a} \int_0^a z n_p(z, t) dz$$

Momentum of nucleons:

$$P_q(t) = \hbar \int_0^a j_q(z, t) dz$$

Total momentum:

$$P_{\text{tot}}(t) = P_n(t) + P_p(t)$$



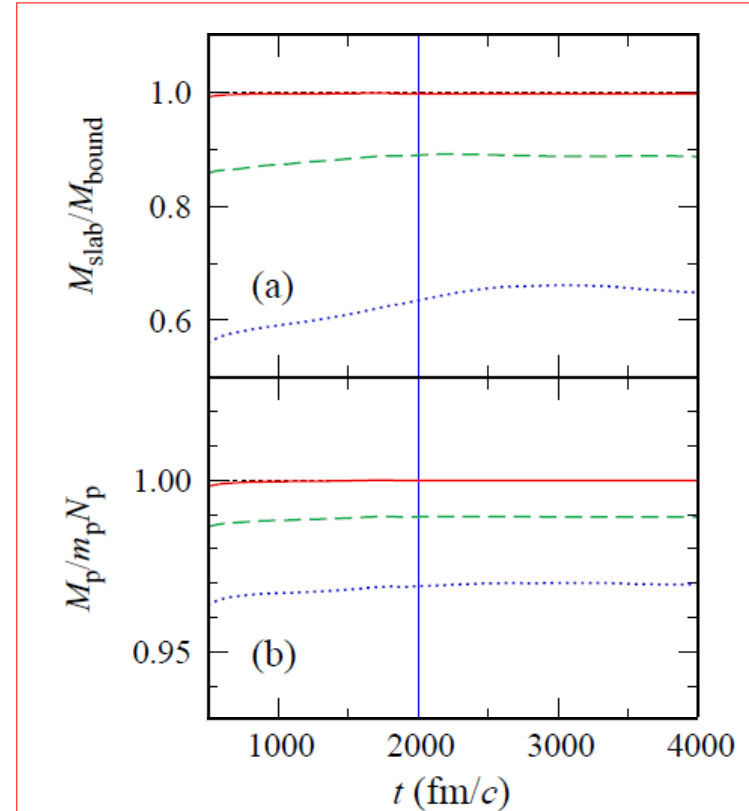
$$M_{\text{slab}} = \dot{P}_{\text{tot}}/a_p$$



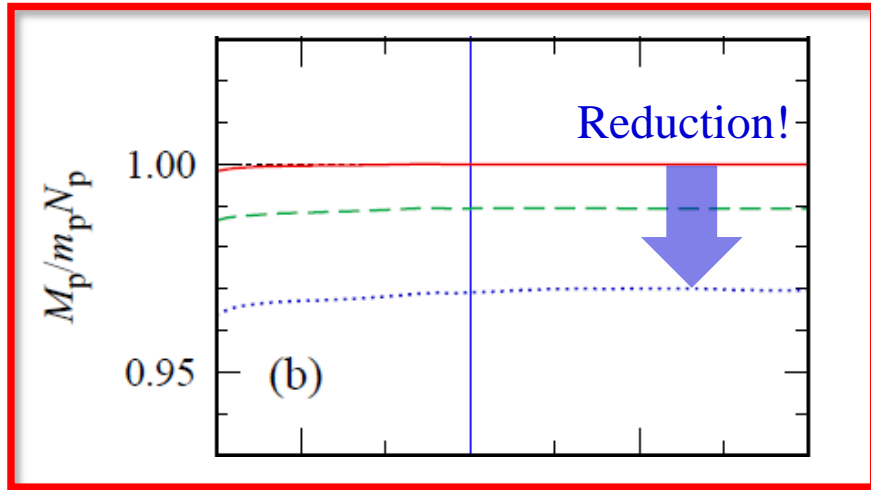
$$M_p = \dot{P}_p/a_p$$

✓ For neutron-dripped slabs, we find significant **reduction** of the collective mass!

➤ What is the origin of the reduction?



- ✓ Cause of the reduction of the collective mass of protons:
the density-dependent “microscopic” effective mass



Collective mass of protons

$$M_p \leq m_p N_p$$

$$\approx m_p^\oplus [n_n^{\text{b.g.}}] N_p$$

Protons and bound neutrons move together



There must be a velocity lag between protons and background neutrons!

The continuity equation within Skyrme TDDFT reads:

$$\frac{\partial \rho_q(\mathbf{r}, t)}{\partial t} + \hbar \nabla \cdot \mathbf{p}_q(\mathbf{r}, t) = 0$$

where

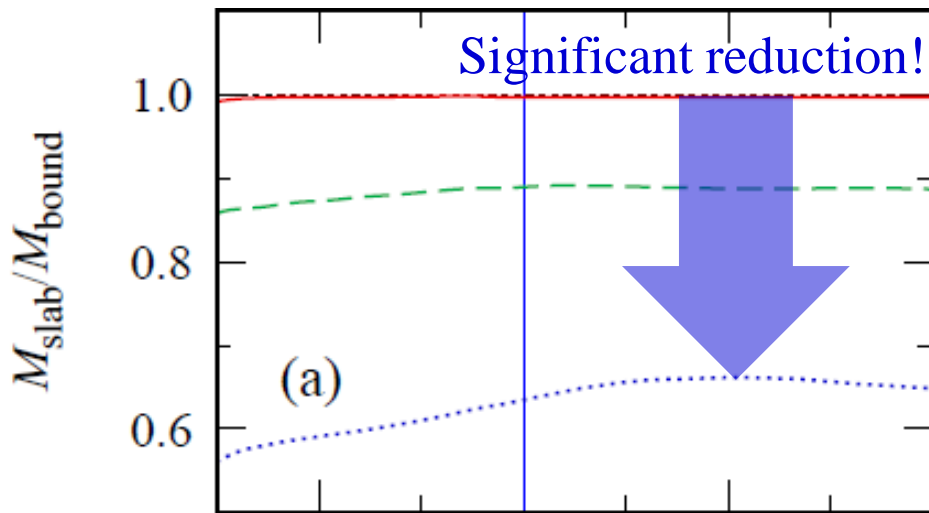
$$\mathbf{p}_q(\mathbf{r}, t) = \mathbf{j}_q(\mathbf{r}, t) + q \frac{2m_q}{\hbar^2} (C_0^\tau - C_1^\tau) n_n(\mathbf{r}, t) n_p(\mathbf{r}, t) \left(\frac{\mathbf{j}_p(\mathbf{r}, t)}{n_p(\mathbf{r}, t)} - \frac{\mathbf{j}_n(\mathbf{r}, t)}{n_n(\mathbf{r}, t)} \right)$$

+1 for protons
-1 for neutrons

velocity difference

Then, what is the cause of the reduction
of the collective mass of the slab?

→ an “anti-entrainment” effect!



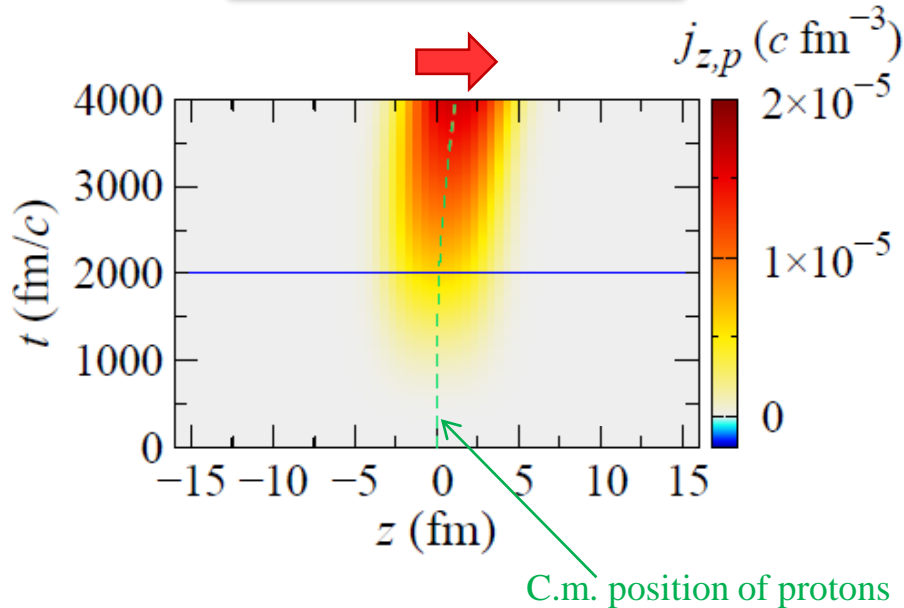
It can **not** be explained solely by
the microscopic effective mass.

Current density:

$$j_{z,q}(z,t) = \frac{\hbar}{m_q} \sum_{\alpha,\mathbf{k}}^{\text{occ.}} \text{Im}[\psi_{\alpha\mathbf{k}}^{(q)*}(\mathbf{r},t) \nabla \psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r},t)] = \frac{\hbar}{m_q} \frac{1}{aN_{k_z}} \sum_{\alpha,k_z} \int \frac{k_{\parallel}}{\pi} \text{Im}[u_{\alpha\mathbf{k}}^{(q)*}(z,t)(\partial_z + ik_z)u_{\alpha\mathbf{k}}^{(q)}(z,t)] \theta(\mu_q - \varepsilon_{\alpha\mathbf{k}}^{(q)}) dk_{\parallel}$$

- ✓ Protons inside the slab move toward the direction of the external force, as expected.

Proton current density



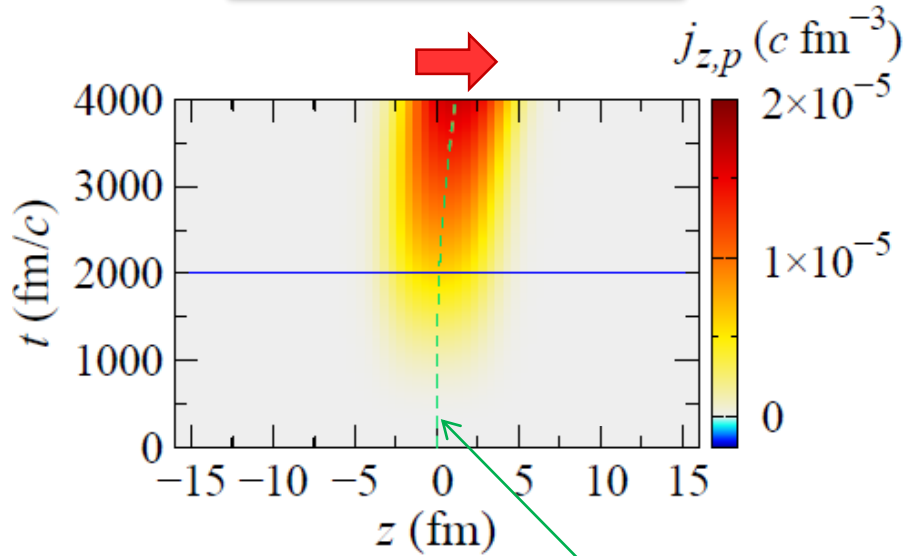
Current density:

$$j_{z,q}(z,t) = \frac{\hbar}{m_q} \sum_{\alpha, \mathbf{k}}^{\text{occ.}} \text{Im}[\psi_{\alpha\mathbf{k}}^{(q)*}(\mathbf{r}, t) \nabla \psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r}, t)] = \frac{\hbar}{m_q} \frac{1}{aN_{k_z}} \sum_{\alpha, k_z} \int \frac{k_{\parallel}}{\pi} \text{Im}[u_{\alpha\mathbf{k}}^{(q)*}(z,t) (\partial_z + ik_z) u_{\alpha\mathbf{k}}^{(q)}(z,t)] \theta(\mu_q - \varepsilon_{\alpha\mathbf{k}}^{(q)}) dk_{\parallel}$$

✓ Dripped neutrons outside the slab move toward the opposite direction!

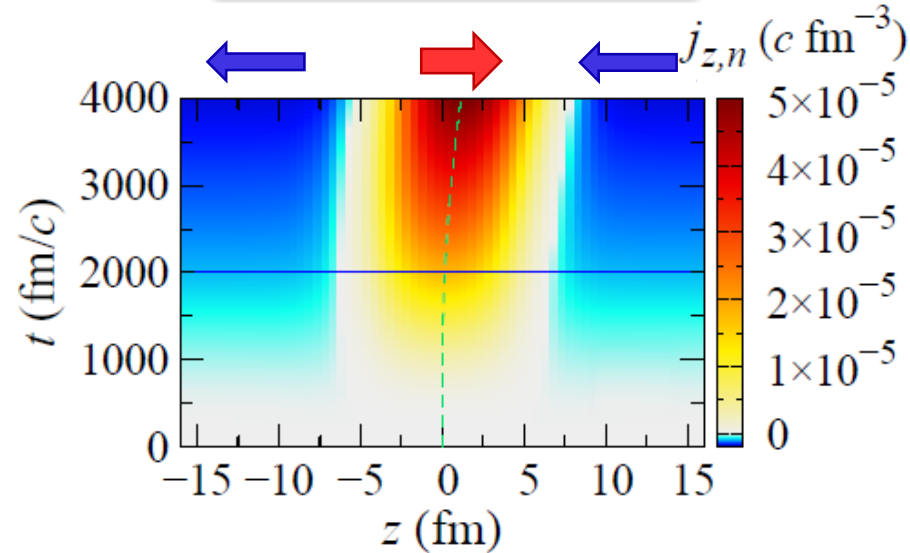
Since it reduces P_{tot} and \dot{P}_{tot} , $M_{\text{slab}} = \dot{P}_{\text{tot}}/a_p$ is reduced

Proton current density



C.m. position of protons

Neutron current density



$$(m_{n,\alpha\mathbf{k}}^{*-1})_{\mu\nu} = \frac{1}{\hbar^2} \frac{\partial^2 \varepsilon_{\alpha\mathbf{k}}^{(n)}}{\partial k_{\mu} \partial k_{\nu}}$$

Current density:

$$j_{z,q}(z,t) = \frac{\hbar}{m_q} \sum_{\alpha, \mathbf{k}}^{\text{occ.}} \text{Im}[\psi_{\alpha\mathbf{k}}^{(q)*}(\mathbf{r}, t) \nabla \psi_{\alpha\mathbf{k}}^{(q)}(\mathbf{r}, t)] = \frac{\hbar}{m_q} \frac{1}{aN_{k_z}} \sum_{\alpha, k_z} \int \frac{k_{\parallel}}{\pi} \text{Im}[u_{\alpha\mathbf{k}}^{(q)*}(z,t) (\partial_z + ik_z) u_{\alpha\mathbf{k}}^{(q)}(z,t)] \theta(\mu_q - \varepsilon_{\alpha\mathbf{k}}^{(q)}) dk_{\parallel}$$

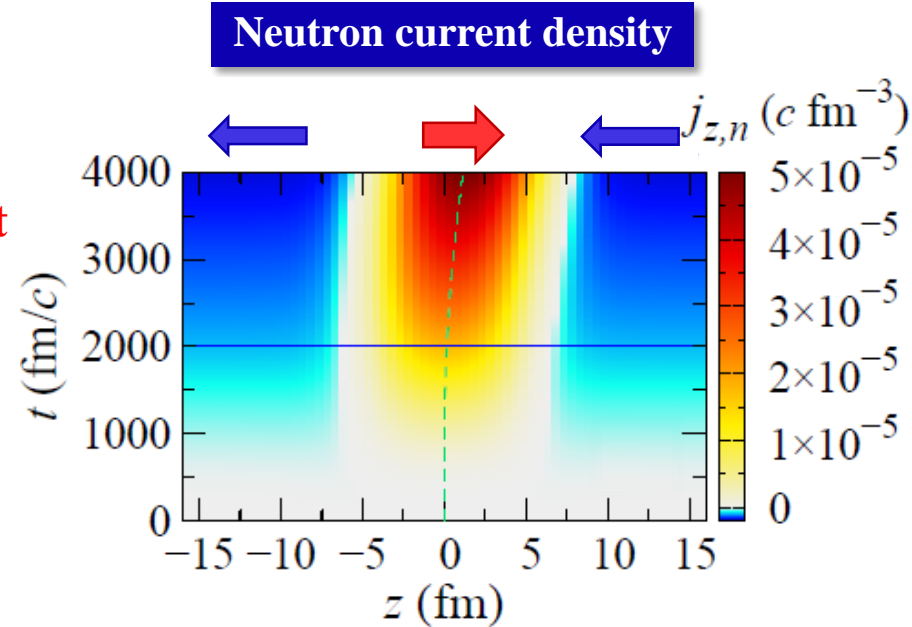
✓ Dripped neutrons outside the slab move toward the opposite direction!

Since it reduces P_{tot} and \dot{P}_{tot} , $M_{\text{slab}} = \dot{P}_{\text{tot}}/a_p$ is reduced

Reduction of M_{slab}
 → enhancement of n_c
 → reduction of m^*

We interpret it as an “anti-entrainment” effect

Y_p	n_n^f/\bar{n}_n	Static		Dynamic
		n_n^c/\bar{n}_n	m_n^*/m_n	n_n^c/\bar{n}_n
0.3	2.09×10^{-4}	0.005	0.040	0.005
0.2	0.127	0.256	0.496	0.229
0.1	0.362	0.630	0.574	0.586



$$(m_{n,\alpha\mathbf{k}}^{*-1})_{\mu\nu} = \frac{1}{\hbar^2} \frac{\partial^2 \varepsilon_{\alpha\mathbf{k}}^{(n)}}{\partial k_{\mu} \partial k_{\nu}}$$

At the frontiers in nuclear physics III:

Neutron-star merger, gravitational wave,
and nucleosynthesis

Fe is the most stable!



47Ag



79Au

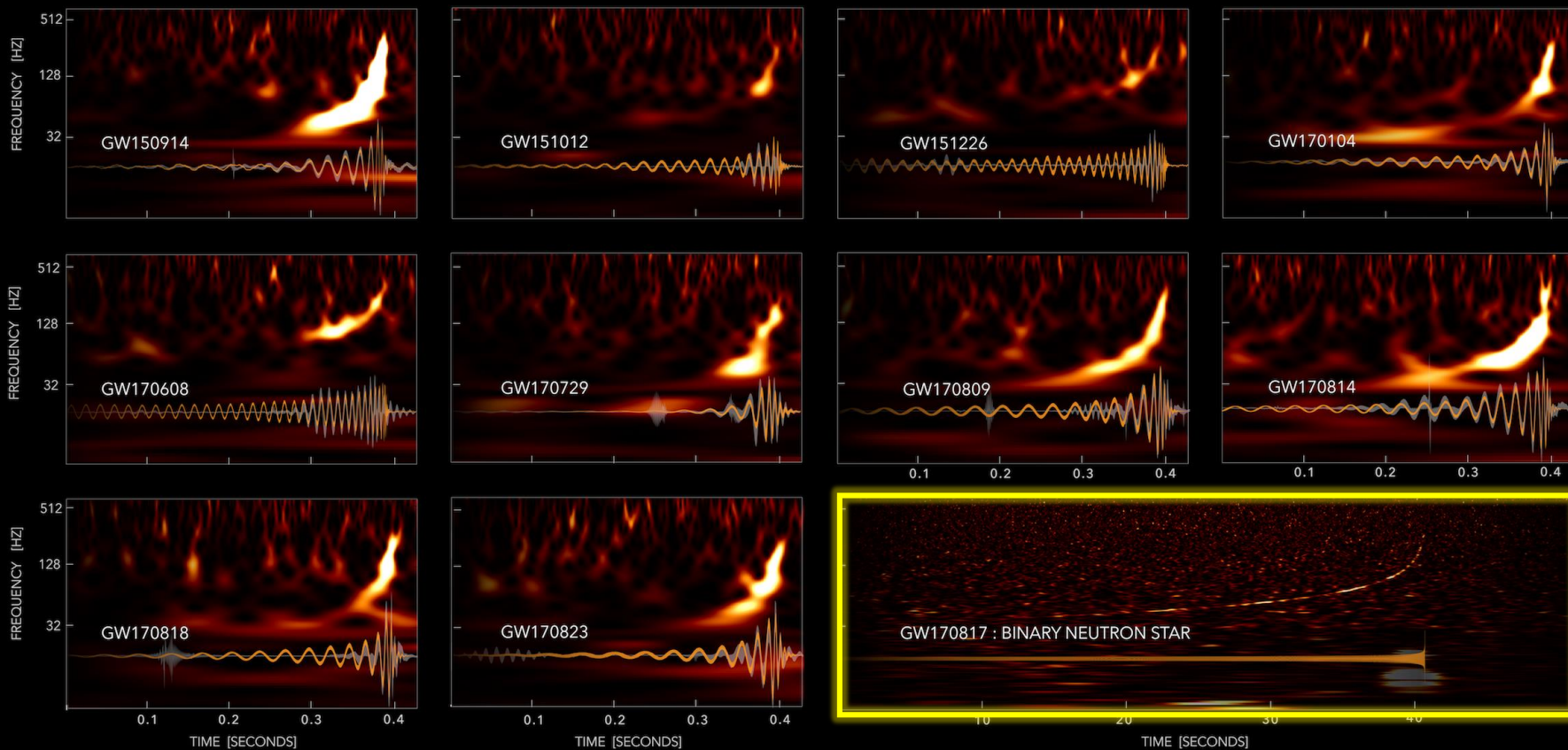
Well, then how were elements heavier than iron produced!?

One of the unsolved problems in Physics

but, we have learned a lot!

Dawn of a new era of the multi-messenger astronomy

GRAVITATIONAL-WAVE TRANSIENT CATALOG-1



Neutron Star Merger!!

Relevant to gravitational waves, nucleosynthesis, as well as neutron stars

Gravitational waves and gamma rays were detected on Aug. 17, 2017

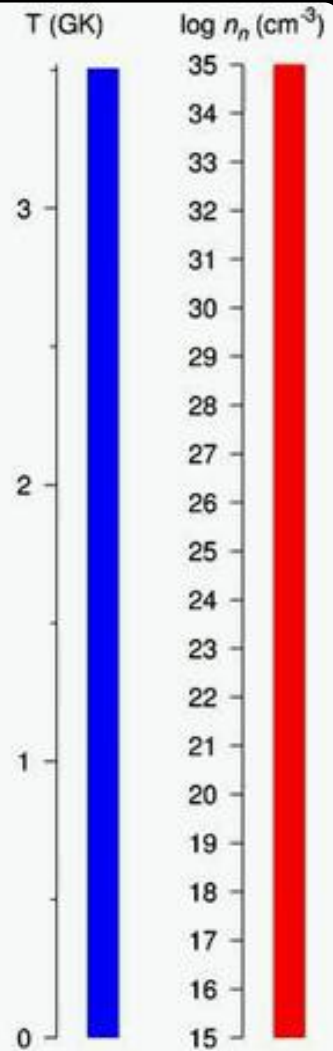
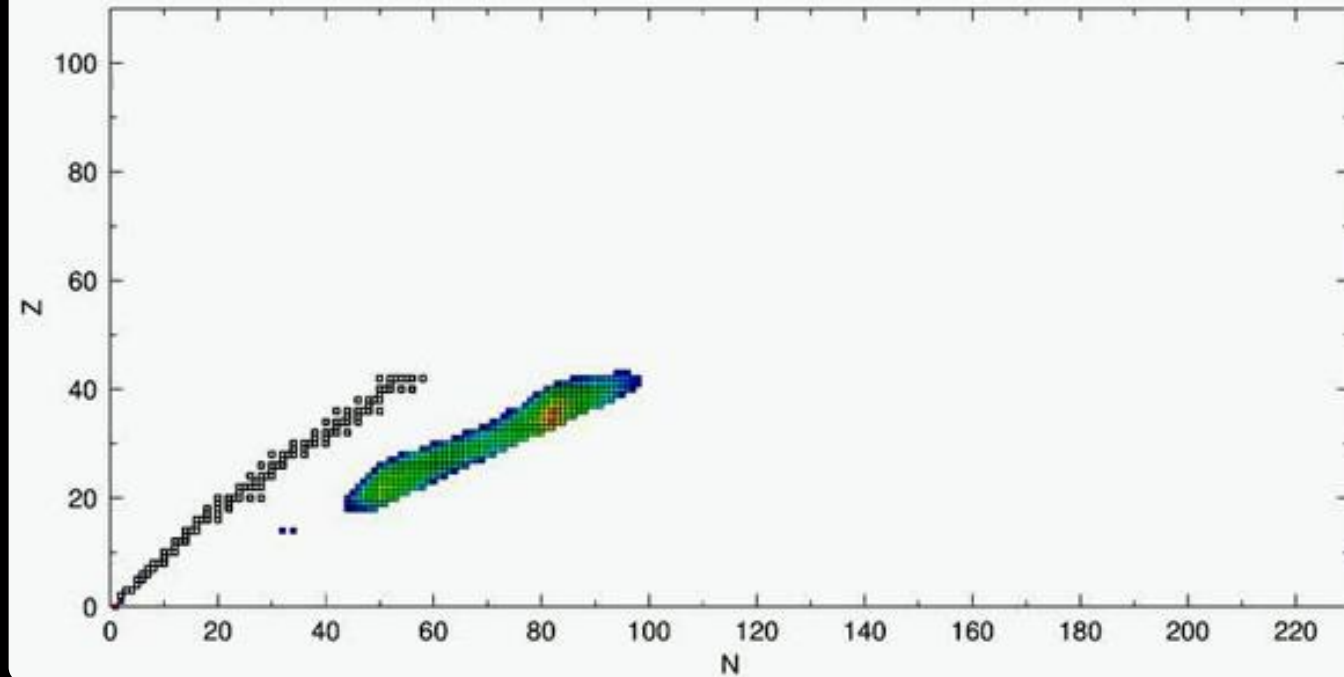
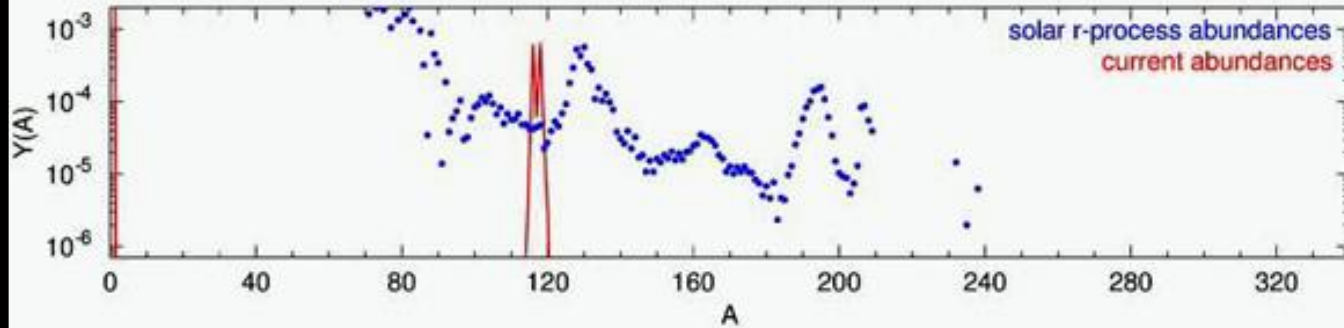
NASA's Goddard Space Flight Center

At the frontier in Physics!!

Georgia
Tech 

Simulation of r-process nucleosynthesis during a neutron-star merger

$T = 3.50$ GK, $n_n = 2.937e+35$ cm⁻³, $R_{n/s} = 623.3$, $s = 0.621$ k_B/nuc, $t = 0.0131$ s

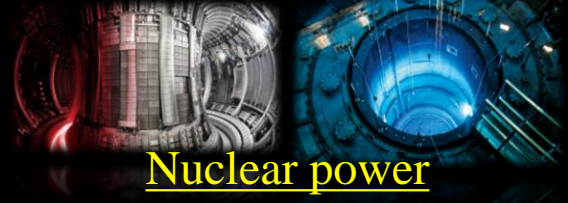




Summary

Nuclear physics - from fundamentals to applications to our life

JET, Oxfordshire, UK A reactor in Switzerland



Nuclear power

Strong + EM + Weak + Gravity

Accelerator science
Detector systems

Equation of state and structure
Superfluidity/superconductivity
Thermal evolution
Supernova explosions

Neutron stars

n, p, e⁻, μ⁻
hyperons
meson condensates
quark matter

Nuclear security 

Medical physics

MRI, Hadron therapy,
Nuclear medicine



CNAO, Pavia, Italy

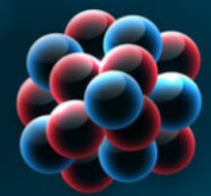
Radiometric dating

e.g. ¹⁴C (≲60,000 year)



Strong + EM + Weak

Structure
Reactions
Decays



Atomic Nuclei

neutrons and protons

Nuclear force

Strong force

QCD

quarks and gluons

Neutron star merger
Gravitational waves
Origin of the elements

Physics beyond the Standard Model
0νββ, CP violation (EDM, CN resonance),
X11, Axion, Dark matter, ...

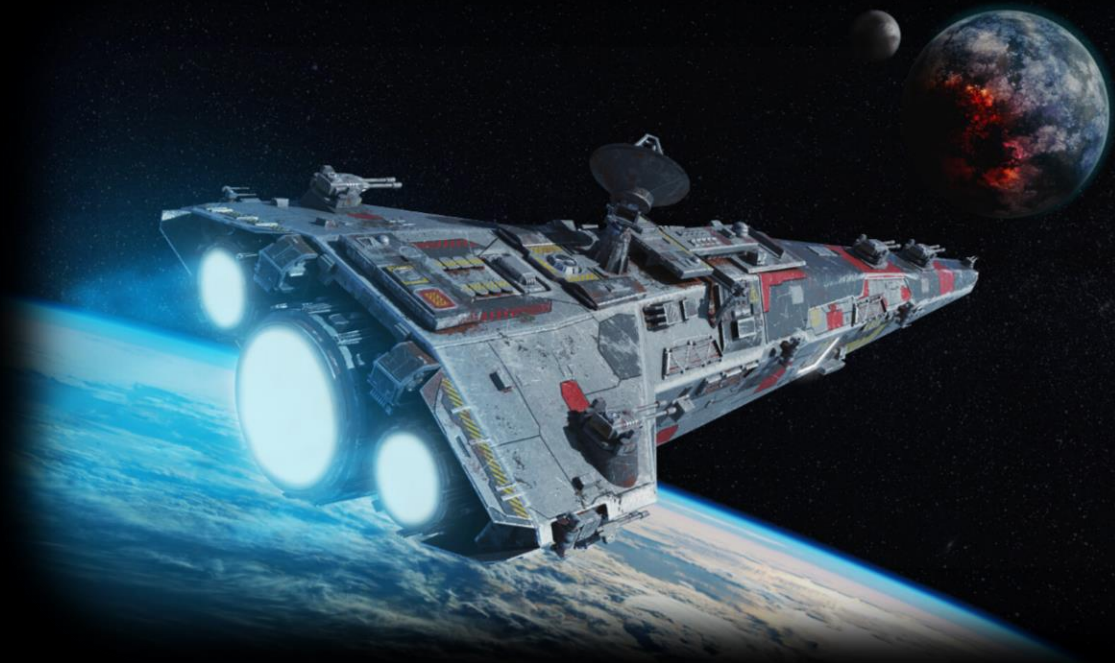


*Each of all pictures has been linked to its source URL.

I hope you enjoyed our exciting adventure!



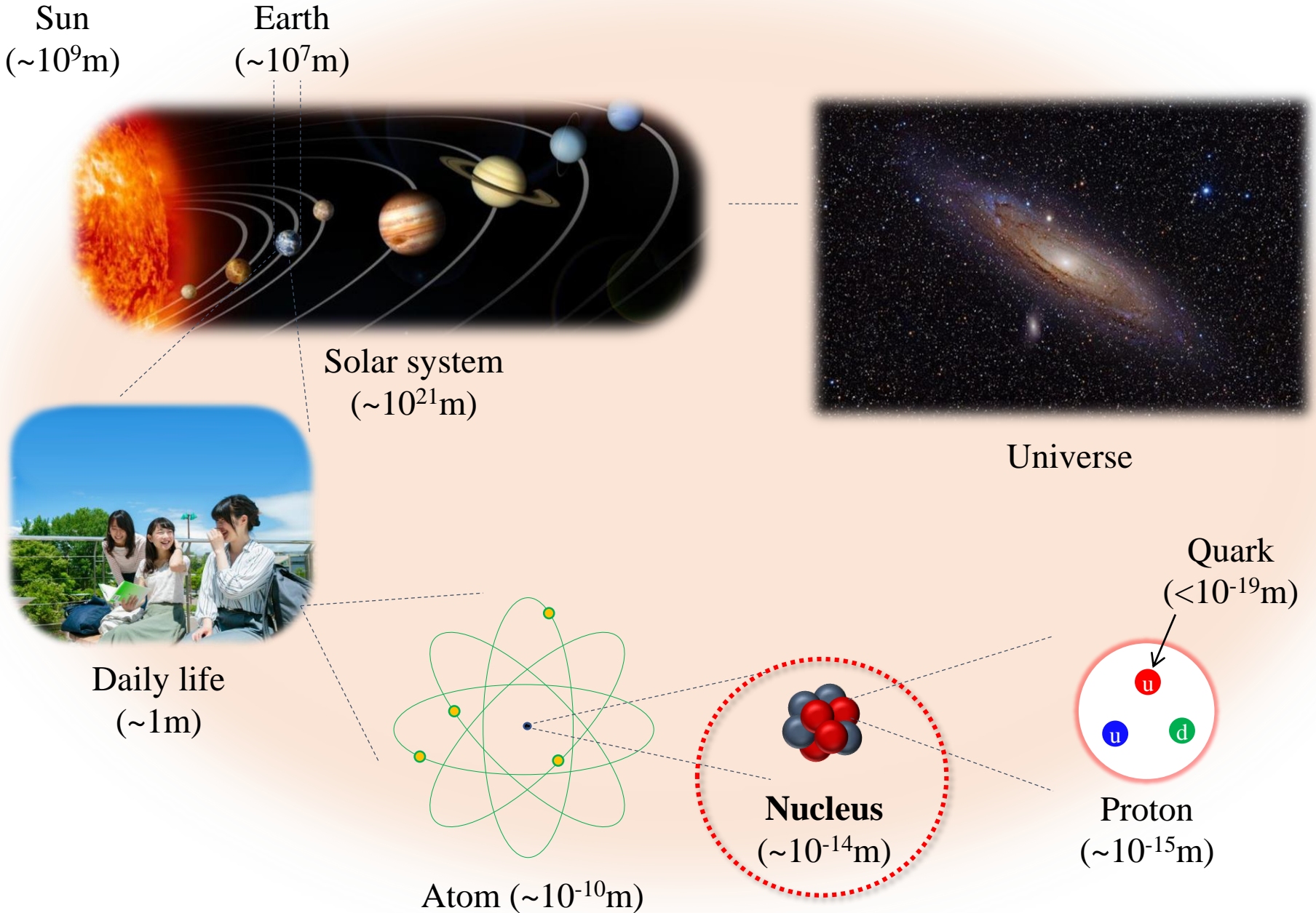
with a magic of DFT and TDDFT..



The transdisciplinary character is
one of the fascinating points of Nuclear Physics

What I showed today are only a tiny part of the huge field! :)

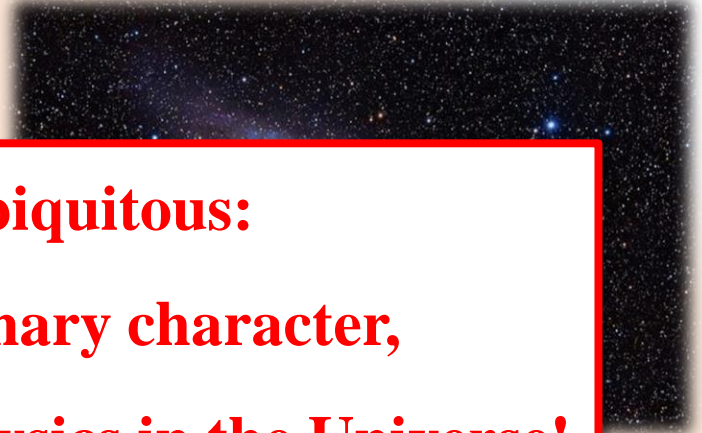
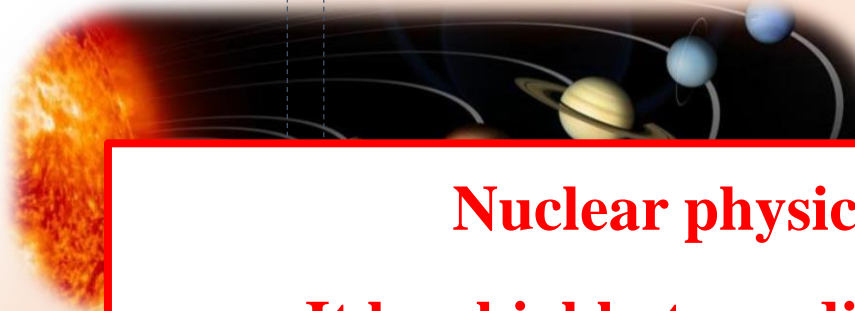
Takeaway message



Takeaway message

Sun
($\sim 10^9\text{m}$)

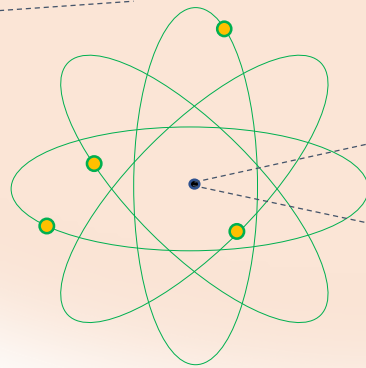
Earth
($\sim 10^7\text{m}$)



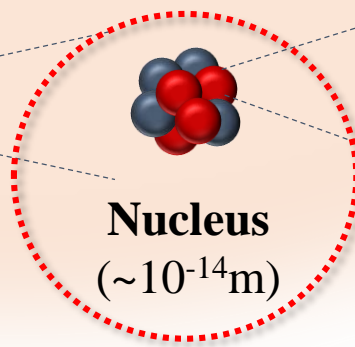
**Nuclear physics is ubiquitous:
It has highly transdisciplinary character,
connecting different scales of physics in the Universe!**



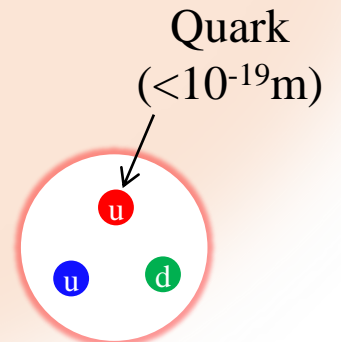
Daily life
($\sim 1\text{m}$)



Atom ($\sim 10^{-10}\text{m}$)



Nucleus
($\sim 10^{-14}\text{m}$)



Proton
($\sim 10^{-15}\text{m}$)

Quark
($< 10^{-19}\text{m}$)

Kazuyuki Sekizawa

Associate Professor

Department of Physics, School of Science

Tokyo Institute of Technology

2-12-1 O-Okayama, Meguro, Tokyo 152-8551, Japan

sekizawa @ phys.titech.ac.jp

About me: <http://sekizawa.fizyka.pw.edu.pl/english/>

About us: <https://nuclphystitech.wordpress.com/>

See also:



A person wearing a blue shirt, brown pants, a grey helmet, and climbing gear is rappelling down a light-colored rock face. The background is a clear blue sky with some distant mountains visible. The text is overlaid on the right side of the image.

Message for undergraduate course (BSc) students

All the subjects you learn (e.g., classical mechanics, electromagnetism, analytical mechanics, thermodynamics, statistical mechanics, quantum mechanics, etc.) **are indispensable to explore the wonderful world of physics in the universe.** It's like equipment for climbing. When completed, you'll see the breathtaking beauty of the nature that may change the rest of your life! :)

Study hard, be ambitious, and have fun!



**Message for
graduate course (MSc and PhD) students**

Study what interests you the most. Dig it deeper and deeper.

It doesn't matter what others say (of course it can be useful as an "opinion" though). Try to reach the deepest ever achieved and find a way to dig in further. This is your thesis work. **It will form your "roots."** Then, **a thick and high trunk, a wide variety of branches, and abundant fruits of your research "tree" will grow up.** The experience to grow up the tree allows you to plant other trees as well.

Be confident, anything can be interesting!